Mean-field behaviour of long- and finite range Ising model, percolation and self-avoiding walk

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The Ising model on \mathbb{Z}^d

Let $\Lambda \subset \subset \mathbb{Z}^d$. Energy of a configuration $\varphi = \{\varphi_x | x \in \Lambda\} \in \{-1, +1\}^{\Lambda}$ defined by

$$\mathcal{H}_{\Lambda}(\varphi) := -\sum_{\{x,y\} \in \Lambda \times \Lambda} J_{\{x,y\}} \, \varphi_x \, \varphi_y.$$

Assume J symmetric and $J \ge 0$ (ferromagnetic interaction)

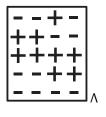
Example: $J_{\{x,y\}} = \mathbb{1}_{|x-y|=1}$ (nearest-neighbor model).

Probability of configuration $\varphi \in \{-1, +1\}^{\Lambda}$ is $e^{-\beta \mathcal{H}_{\Lambda}(\varphi)}/Z_{\Lambda}(\beta)$. Two-point correlation function:

$$\langle \varphi_0 \varphi_{\mathsf{x}} \rangle_{\mathsf{\Lambda}} := \frac{\sum_{\boldsymbol{\varphi} \in \{-1,1\}^{\mathsf{\Lambda}}} \varphi_0 \, \varphi_{\mathsf{x}} \exp(-\beta \mathcal{H}_{\mathsf{\Lambda}}(\boldsymbol{\varphi}))}{\sum_{\boldsymbol{\varphi} \in \{-1,1\}^{\mathsf{\Lambda}}} \exp(-\beta \mathcal{H}_{\mathsf{\Lambda}}(\boldsymbol{\varphi}))}.$$

Phase transition: $\exists m, M > 0$:

$$G_{\beta}(x) := \lim_{\Lambda \nearrow \mathbb{Z}^d} \langle \varphi_0 \varphi_x \rangle_{\Lambda} \begin{cases} \leq e^{-m|x|} & (\beta < \beta_c) \\ \geq M & (\beta > \beta_c). \end{cases}$$



Aim: Understanding of *critical* behaviour, i.e., $\beta = \beta_c$

Critical behaviour

Physics concept: Introduce critical exponents, e.g.

$$G_{\beta_c}(x) \stackrel{|x| \nearrow \infty}{\approx} |x|^{-(d-2+\eta)}, \qquad \sum_x G_{\beta}(x) \stackrel{\beta \nearrow \beta_c}{\approx} (\beta_c - \beta)^{-\gamma}.$$

Theorem (Aizenman, Barsky, Fernández, Graham '80s)

If
$$\sum_{x} G_{\beta_c}(x)^2 < \infty$$
 (bubble condition), then $\eta = 0$, $\gamma = 1$, $\delta = 3$, ...

Note: These values for critical exponents coincide with Curie-Weiss model (=Ising model on *complete* graph). Known as *mean-field values*.

Theorem (Fröhlich, Simon, Spencer '76)

For RP-version of J (reflection positivity) an *infrared bound* holds:

$$0 \le \hat{G}_{\beta}(k) \le \frac{\mathrm{const}}{|k|^2} \qquad (d \ge 2, \beta < \beta_c),$$

i.e., mean-field behaviour above 4 dimensions.

The spin-spin coupling J

Recall:

$$egin{aligned} G_eta(x) &:= \lim_{\Lambda
earrow \mathbb{Z}^d} rac{\sum_{oldsymbol{arphi} \in \{-1,1\}^\Lambda} arphi_0 \, arphi_x \exp(-eta \mathcal{H}_\Lambda(oldsymbol{arphi}))}{\sum_{oldsymbol{arphi} \in \{-1,1\}^\Lambda} \exp(-eta \mathcal{H}_\Lambda(oldsymbol{arphi}))}, \ \mathcal{H}_\Lambda(oldsymbol{arphi}) &:= -\sum_{\{x,y\} \in \Lambda imes \Lambda} J_{\{x,y\}} \, arphi_x \, arphi_y. \end{aligned}$$

Question: What is the role of J?

Introduce
$$D$$
 by $D(x) = \frac{\tanh(\beta J_{\{0,x\}})}{\sum_{y \in \mathbb{Z}^d} \tanh(\beta J_{\{0,y\}})}$.

Consider three cases:

- nearest-neighbor coupling: $D(x) = \frac{1}{2d} \mathbb{1}_{\{|x|=1\}}$
- finite-variance coupling: e.g., $D(x) \approx e^{-|x/L|}$ as $|x| \to \infty$, L suff. large
- long-range coupling: e.g., $D(x) \approx |x/L|^{-(d+\alpha)}$ as $|x| \to \infty$, $\alpha > 0$, L suff. large

Result

Fourier transform:

$$1-\hat{D}(k)=1-\sum_{x\in\mathbb{Z}^d}\mathrm{e}^{ik\cdot x}\ D(x)=\sum_{x\in\mathbb{Z}^d}[1-\cos(k\cdot x)]\ D(x)symp |k|^{2\wedge lpha}.$$

Theorem (HHS 2007)

Consider the Ising model in dimension

$$d \begin{cases} > 4 & \text{(finite-variance case)}, \\ > 2(2 \wedge \alpha) & \text{(power-law case)}, \end{cases}$$

and L sufficiently large. Then

$$\hat{G}_{z_c}(k) = \frac{1 + O(L^{-d})}{1 - \hat{D}(k)} \times (1 + O(L^{-d})) |k|^{-(2 \wedge \alpha)}.$$
 (1)

- Finite range models known by Sakai (CMP '07).
- Proof uses lace expansion.
- Result holds also for nearest-neighbor model with $O(L^{-d})$ replaced by O(1/d) and d suff. large.

Random Current representation

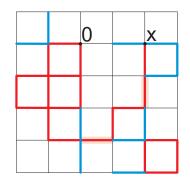
Starting point: RC representation by Griffiths, Hurst, Sherman Denote $\mathbb{B}_{\Lambda} = \big\{ \{x,y\} \mid x,y \in \Lambda, x \neq y \big\}$ set of bonds in Λ .

$$\begin{split} \langle \varphi_{0} \varphi_{x} \rangle_{\Lambda} &= \frac{2^{-|\Lambda|}}{Z_{\Lambda}} \sum_{\varphi \in \{\pm 1\}^{\Lambda}} \varphi_{0} \, \varphi_{x} \prod_{\{u,v\} \in \mathbb{B}_{\Lambda}} \exp(-\beta J_{\{u,v\}} \, \varphi_{u} \, \varphi_{v}) \\ &= \frac{2^{-|\Lambda|}}{Z_{\Lambda}} \sum_{\varphi \in \{\pm 1\}^{\Lambda}} \varphi_{0} \, \varphi_{x} \sum_{\mathbf{n} \in \mathbb{N}^{\mathbb{B}_{\Lambda}}} \left(\prod_{\underline{b} \in \mathbb{B}_{\Lambda}} \frac{(\beta J_{b})^{n_{b}}}{n_{b}!} \right) \prod_{v \in \Lambda} \varphi_{v}^{\sum_{\bar{v} \in \Lambda} n_{\{v,\bar{v}\}}} \\ &= \sum_{\mathbf{n} \in \mathbb{N}^{\mathbb{B}_{\Lambda}}} \frac{w_{\lambda}(\mathbf{n})}{Z_{\Lambda}} \prod_{v \in \Lambda} \left(\frac{1}{2} \sum_{\varphi_{v} = \pm 1} \varphi_{v}^{\mathbb{I}_{\{v \in \{0\} \Delta \{x\}\}} + \sum_{\bar{v} \in \Lambda} n_{\{v,\bar{v}\}}} \right) \\ &= \sum_{\partial \mathbf{n} = \{0\} \Delta \{x\}} \frac{w_{\lambda}(\mathbf{n})}{Z_{\Lambda}} \quad \left(\partial \mathbf{n} := \{v \in \Lambda : \sum_{\bar{v}} n_{\{v,\bar{v}\} \text{ is odd}} \} \right) \end{split}$$

Lace expansion for the Ising model (Sakai '07)

This gives rise to a representation involving bonds:

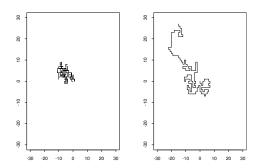
$$\langle \varphi_0 \varphi_x \rangle_{\Lambda} = \sum_{\substack{\mathbf{n} \in \mathbb{N}^{\mathbb{B}_{\Lambda}} \\ \partial \mathbf{n} = \{0\} \triangle \{x\}}} \frac{w_{\lambda}(\mathbf{n})}{Z_{\Lambda}}$$



Perform *lace expansion* using ideas from lace expansion for **percolation**. Use of BK-inequality replaced by use of RHS source switching lemma.

Self-avoiding walk

A *self-avoiding walk* is random walk conditioned to visit every site at most once.



Realization 200-step (nearest-neighbor) SRW and SAW for d=2.

Interesting: The asymptotic difference between SRW and SAW (in terms of critical exponents) disappears in dimension $d \ge 5$.

Recall versions of D:

- nearest-neighbor coupling: $D(x) = \frac{1}{2d} \mathbb{1}_{\{|x|=1\}}$
- finite-variance coupling: e.g., $D(x) \approx \mathrm{e}^{-|x/L|}$ as $|x| \to \infty$, L large
- long-range coupling: e.g., $D(x) \approx |x/L|^{-(d+\alpha)}$ as $|x| \to \infty$, L large

Let $W_n(x)$ be set of *n*-step walks from the origin 0 to x:

$$W_n(x) = \{(w_0, \dots, w_n) \mid w_0 = 0, w_n = x, w_i \in \mathbb{Z}^d, 1 \le i \le n-1\}$$

Call $w \in \mathcal{W}_n(x)$ self-avoiding if $w_i \neq w_j$ for $i \neq j$ with $i, j \in \{0, ..., n\}$. Define *SAW two-point function* by $c_0(x) = \delta_{0,x}$ and, for $n \geq 1$,

$$c_n(x) := \sum_{w \in \mathcal{W}_n(x)} \prod_{i=1}^n D(w_i - w_{i-1}) \, \mathbb{1}_{\{w \text{ is self-avoiding}\}}.$$

Self-avoiding walk Green's function:

$$G_z(x) := \sum_{n=0}^{\infty} c_n(x) z^n$$

Question: How does $G_z(x)$ behave?



Self-avoiding walk - result

Theorem (HHS)

Let d sufficiently large in the nearest-neighbor case, or d>4 and L sufficiently large in the finite-variance spread-out case, or $d>2(2\wedge\alpha)$ and L sufficiently large in the spread-out power-law case. Then

$$\hat{G}_z(k) = \frac{1 + O(L^{-d})}{\chi(z)^{-1} + z[1 - \hat{D}(k)]}$$
 (2)

uniformly for $z \in [0, z_c)$ and $k \in [-\pi, \pi)^d$.

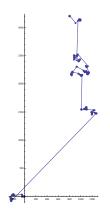
Here, $\chi(z)=\sum_x G_z(x)=\hat{G}_z(0)$ is the *susceptibility*. The critical value z_c is the radius of convergence for the power series defining G_z ; it is characterized by $z_c=\sup\{z\mid \chi(z)<\infty\}$. In the nearest-neighbor case, $O(L^{-d})$ replaced with O(1/d). As for the Ising model,

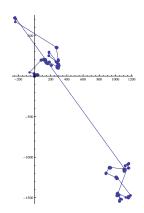
$$\hat{G}_{z_c}(k) = (1 + O(L^{-d})) [1 - \hat{D}(k)]^{-1} \times (1 + O(L^{-d})) |k|^{-(2 \wedge \alpha)}$$

and this implies mean-field behaviour.



Long-range walks





Realizations 500-step long-range random walk with $\alpha = 0.7$ for d = 2.