

Random walks in dynamic random environment: LLN and LDP

Luca Avena

Leiden University

joint work with Frank den Hollander and Frank Redig

WORKSHOP, Dynamic Random Environments, EURANDOM Eindhoven,
14-18 December 2009

Outline of the talk

- Introduction and Model.
- Law of large numbers.
- Large deviations and slow-down phenomenon.

RWRE in a Nutshell

A Random Walk in a Random Environment (RWRE) on \mathbb{Z}^d is a Markov chain/process with **random transition probabilities / rates** , namely

$$p(x, y | \xi) = \text{probability to go from } x \text{ to } y$$

depends on a random field or random process ξ called **Random Environment (RE)**.

Goal: to extend the theory of "classical" random walk, i.e., LLN, CLT, LDP, etc...

Static RE: when ξ is chosen at random at time 0 and kept fixed throughout the time evolution.

In $d = 1$ completely solved, in $d \geq 2$ many open problems.

Dynamic RE: when ξ changes in time according to some stochastic dynamics. Very little known, even in $d = 1$.

Model: dynamic random environment

Let $\xi = \left\{ \xi(x, t) : x \in \mathbb{Z}, t \geq 0 \right\}$

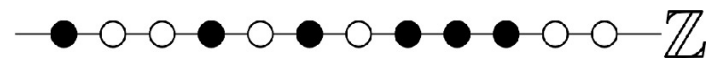
be a one-dimensional **Interacting Particle System (IPS)** on the integer lattice, namely, a Markov process with state space $\Omega = \{0, 1\}^{\mathbb{Z}}$.

$\xi(x, t) = 1 \leftrightarrow$ *Particle at site x at time t ,*

$\xi(x, t) = 0 \leftrightarrow$ *Vacancy at site x at time t .*

Examples: Independent spin-flips , stochastic Ising model, contact process, exclusion process, voter model, majority vote model.

For $\eta \in \Omega = \{0, 1\}^{\mathbb{Z}}$, P^η represents the law of ξ starting from η .



Given a measure μ on $\Omega = \{0, 1\}^{\mathbb{Z}}$, we denote by

$$P^\mu(\cdot) = \int_{\Omega} P^\eta(\cdot) \mu(d\eta).$$

the law of ξ starting from μ .

Assumption: ξ has an equilibrium measure μ that is shift-invariant and shift-ergodic.

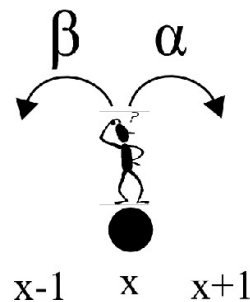
Model: random walk

Given a realization of ξ , let $X = \{X_t\}_{t \geq 0}$, $X_0 = 0$, be the continuous-time nearest-neighbor random walk with **transition rates**

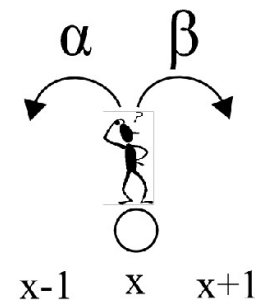
$$x \rightarrow x + 1 \quad \text{at rate} \quad \alpha \xi(x, t) + \beta [1 - \xi(x, t)]$$

$$x \rightarrow x - 1 \quad \text{at rate} \quad \beta \xi(x, t) + \alpha [1 - \xi(x, t)],$$

where $\alpha, \beta > 0$. Without loss of generality **we may assume $\alpha > \beta$** .



local **drift to the right on particles**



local **drift to the left on holes.**

Quenched Law

$P_0^\xi(X_t \in \cdot)$ = law of X_t starting from 0 **conditional on the environment** ξ .

Annealed Law

$$\mathbb{P}_{\mu,0}(X_t \in \cdot) = \int_{D_\Omega[0,\infty)} P^\mu(d\xi) P_0^\xi(X_t \in \cdot)$$

= law of X_t **averaged over the environment** ξ ,
where $D_\Omega[0,\infty)$ is the set of càdlàg paths in $\Omega = \{0, 1\}^{\mathbb{Z}}$.

Law of Large Numbers (LLN)

Theorem 1 (LLN) Suppose that ξ is **cone-mixing**. Then there exists a deterministic **speed** $v \in \mathbb{R}$ such that

$$\lim_{t \rightarrow \infty} \frac{X_t}{t} = v \quad P^\mu - a.s.$$

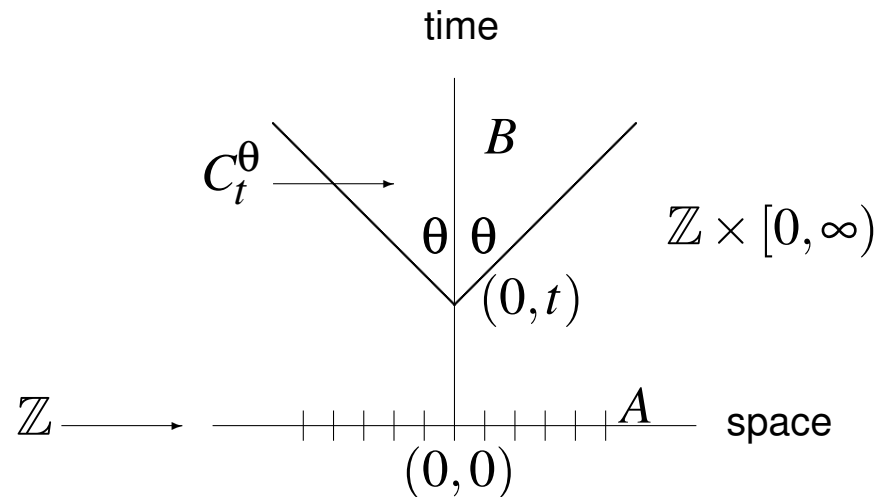
Remark 1: For **small local drift** $\alpha - \beta$, and **exponentially mixing** environments, we have an **explicit formula for the speed** (see **F. Redig's** talk).

Remark 2: Theorem 1 can be extended to higher dimension, provided the total jump rate of the walker is the same on top of particles and holes.

Definition: A random environment ξ is called **cone-mixing** if, for all $\theta \in (0, \frac{1}{2}\pi)$,

$$\lim_{t \rightarrow \infty} \sup_{\substack{A \in \mathcal{F}_0, B \in \mathcal{F}_t^\theta \\ \mu(A) > 0}} \left| P^\mu(B | A) - P^\mu(B) \right| = 0,$$

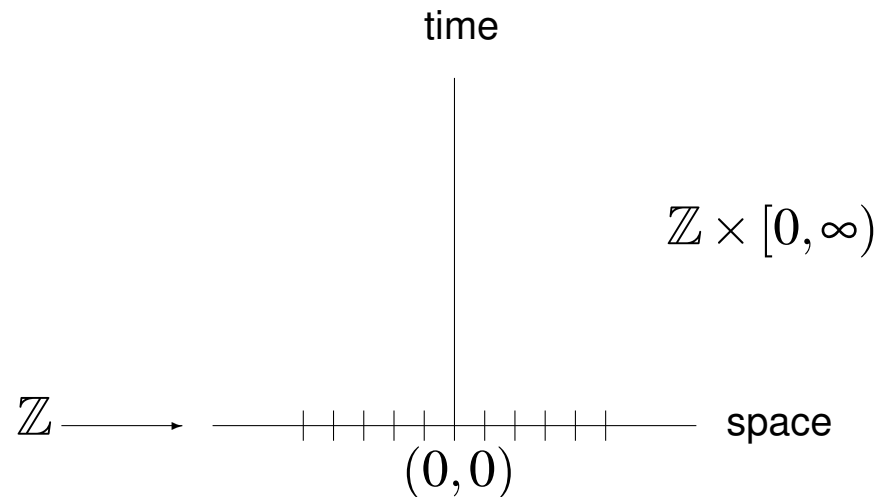
where $\mathcal{F}_0 = \sigma\{\xi_0(x) : x \in \mathbb{Z}\}$ and $\mathcal{F}_t^\theta = \sigma\{\xi_s(x) : (x, s) \in C_t^\theta\}$, with C_t^θ the cone in the picture:



Definition: An environment ξ is called **cone-mixing** if, for all $\theta \in (0, \frac{1}{2}\pi)$,

$$\lim_{t \rightarrow \infty} \sup_{\substack{A \in \mathcal{F}_0, B \in \mathcal{F}_t^\theta \\ \mu(A) > 0}} \left| P^\mu(B | A) - P^\mu(B) \right| = 0,$$

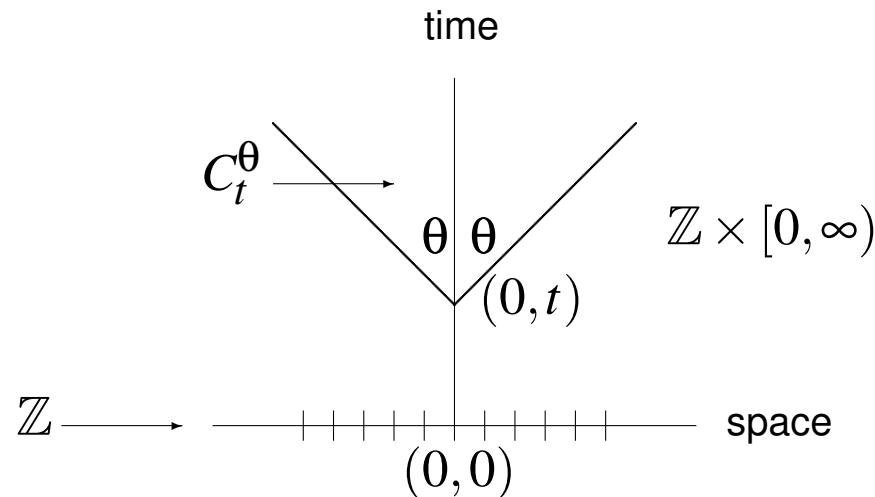
where $\mathcal{F}_0 = \sigma\{\xi_0(x) : x \in \mathbb{Z}\}$ and $\mathcal{F}_t^\theta = \sigma\{\xi_s(x) : (x, s) \in C_t^\theta\}$, with C_t^θ the cone in the picture:



Definition: An environment ξ is called **cone-mixing** if, for all $\theta \in (0, \frac{1}{2}\pi)$,

$$\lim_{t \rightarrow \infty} \sup_{\substack{A \in \mathcal{F}_0, B \in \mathcal{F}_t^\theta \\ \mu(A) > 0}} \left| P^\mu(B | A) - P^\mu(B) \right| = 0,$$

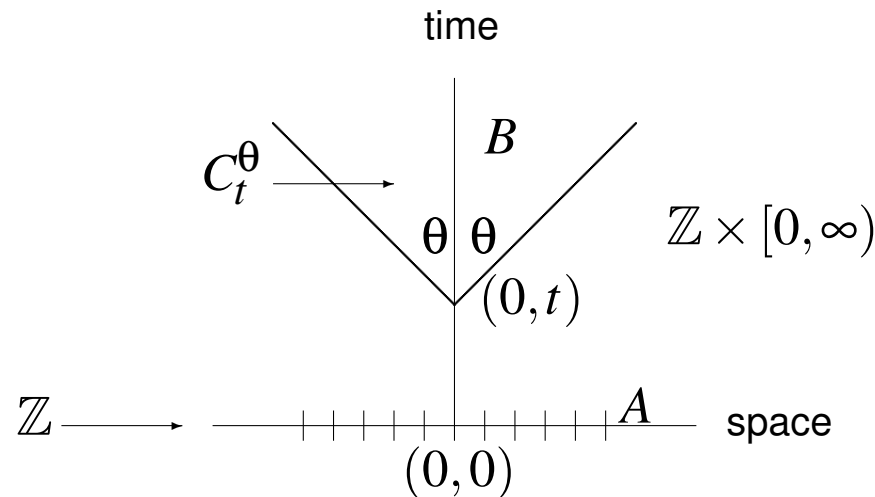
where $\mathcal{F}_0 = \sigma\{\xi_0(x) : x \in \mathbb{Z}\}$ and $\mathcal{F}_t^\theta = \sigma\{\xi_s(x) : (x, s) \in C_t^\theta\}$, with C_t^θ the cone in the picture:



Definition: An environment ξ is called **cone-mixing** if, for all $\theta \in (0, \frac{1}{2}\pi)$,

$$\lim_{t \rightarrow \infty} \sup_{\substack{A \in \mathcal{F}_0, B \in \mathcal{F}_t^\theta \\ \mu(A) > 0}} \left| P^\mu(B | A) - P^\mu(B) \right| = 0,$$

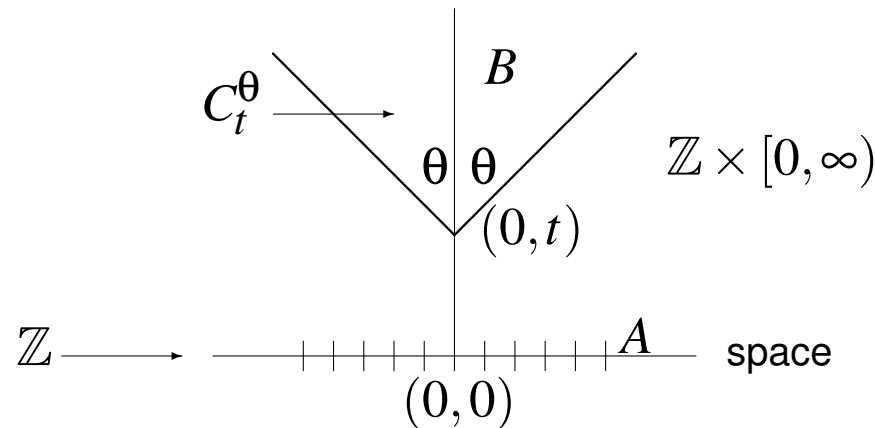
where $\mathcal{F}_0 = \sigma\{\xi_0(x) : x \in \mathbb{Z}\}$ and $\mathcal{F}_t^\theta = \sigma\{\xi_s(x) : (x, s) \in C_t^\theta\}$, with C_t^θ the cone in the picture:



Definition: An environment ξ is called **cone-mixing** if, for all $\theta \in (0, \frac{1}{2}\pi)$,

$$\lim_{t \rightarrow \infty} \sup_{\substack{A \in \mathcal{F}_0, B \in \mathcal{F}_t^\theta \\ \mu(A) > 0}} \left| P^\mu(B | A) - P^\mu(B) \right| = 0,$$

where $\mathcal{F}_0 = \sigma\{\xi_0(x) : x \in \mathbb{Z}\}$ and $\mathcal{F}_t^\theta = \sigma\{\xi_s(x) : (x, s) \in C_t^\theta\}$, with C_t^θ the cone in the picture:



Examples of cone-mixing environments: Independent spin-flips, uniquely ergodic attractive spin systems, IPS under the $M < \varepsilon$ -condition, two-dimensional Gibbs random field embedded in space-time.

Ingredients to prove the LLN

We adapt and simplify a regeneration-time argument originally developed by F. Comets and O. Zeitouni for static Random Environment (RE).

★ we consider a discrete-time version X_n of our walker on \mathbb{Z} in *dynamic* RE, and show that it is equivalent to a random walk $Y_n = (X_n, n)$ on \mathbb{Z}^2 in a *static* RE.

★ the **cone-mixing** property of the environment allows us to construct a sequence of **regeneration times**, i.e., a sequence of times at which the walker Y sees the environment almost as freshly sampled from the equilibrium distribution μ .

★ the increments of Y_n along this sequence of times can be treated as an i.i.d. sequence plus some small extra noise. This allows us to prove the LLN for Y_n , which in turn implies the LLN for X_n .

★ we pass from discrete to continuous time afterwards, via a limiting argument.

Large Deviation Principle (LDP): attractive spin-flip systems

Definition: An IPS ξ is called a **spin-flip system** if only one coordinate changes at each transition, i.e., is a Markov process on $\{0, 1\}^{\mathbb{Z}}$ with generator

$$(Lf)(\eta) = \sum_{x \in \mathbb{Z}} c(x, \eta) [f(\eta^x) - f(\eta)],$$

where

$\eta \in \{0, 1\}^{\mathbb{Z}}$, f = local function, $\eta^x = \eta$ flipped at x ,
 $c(x, \eta)$ = flip rate at site x in the configuration η .

Large Deviation Principle (LDP): attractive spin-flip systems

Definition: An IPS ξ is called a **spin-flip system** if only one coordinate changes at each transition at rate, i.e., is a Markov process on $\{0, 1\}^{\mathbb{Z}}$ with generator

$$(Lf)(\eta) = \sum_{x \in \mathbb{Z}} c(x, \eta) [f(\eta^x) - f(\eta)].$$

Definition: A spin-flip system ξ is called **attractive** if, whenever $\eta \leq \zeta$,

$$\begin{aligned} c(x, \eta) &\leq c(x, \zeta) & \text{if } \eta(x) = \zeta(x) = 0, \\ c(x, \eta) &\geq c(x, \zeta) & \text{if } \eta(x) = \zeta(x) = 1, \quad \text{for all } x \in \mathbb{Z}. \end{aligned}$$

Attractive Idea: The evolution tends to make a coordinate agree with its neighborhood.

Large Deviation Principle (LDP): attractive spin-flip systems

Definition: An IPS ξ is called a **spin-flip system** if only one coordinate changes at each transition at rate, i.e., is a Markov process on $\{0, 1\}^{\mathbb{Z}}$ with generator

$$(Lf)(\eta) = \sum_{x \in \mathbb{Z}} c(x, \eta) [f(\eta^x) - f(\eta)].$$

Definition : A spin-flip system ξ is called **attractive** if, for all $x \in \mathbb{Z}$, whenever $\eta \leq \zeta$,

$$\begin{aligned} c(x, \eta) &\leq c(x, \zeta) & \text{if } \eta(x) = \zeta(x) = 0, \\ c(x, \eta) &\geq c(x, \zeta) & \text{if } \eta(x) = \zeta(x) = 1. \end{aligned}$$

Attractive Idea: The evolution tends to make a coordinate agree with its neighborhood.

Examples of attractive spin-flip systems: Independent spin-flips, ferromagnetic stochastic Ising model, contact process, voter model, majority vote model.

LDP for the Empirical Speed

Informally: We say that the family $\mathbb{P}(X_t/t \in \cdot)$, $t > 0$, satisfies a **Large Deviation Principle (LDP)** with rate t and with rate function $I: \mathbb{R} \rightarrow [0, \infty)$ if

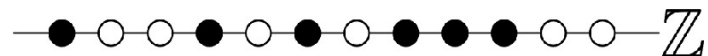
$$\exists \lim_{t \rightarrow \infty} -\frac{1}{t} \log \mathbb{P}\left(\frac{X_t}{t} \simeq \theta\right) = I(\theta) \quad \forall \theta \in \mathbb{R},$$

$$\text{i.e. , } \mathbb{P}\left(\frac{X_t}{t} \simeq \theta\right) \approx e^{-tI(\theta)} \text{ as } t \rightarrow \infty.$$

Remark: The rate function for “classical” random walks is zero at the typical speed v and positive elsewhere.

LDP for *static* Random Environment

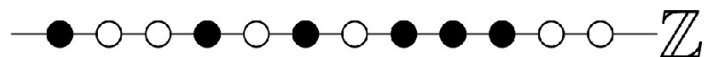
Static RE: $\xi \in \{0, 1\}^{\mathbb{Z}}$ distributed according to a **Bernoulli product measure** ν_ρ with parameter $\rho \in (0, 1)$.



F.Solomon '75 proved LLN. A.Greven, F.den Hollander '94, F.Comets, N.Gantert and O.Zeitouni '00 proved LDPs.

LDP for *static* Random Environment

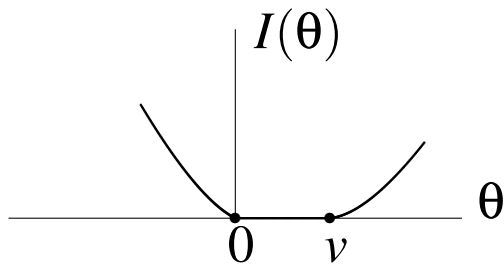
Static RE: $\xi \in \{0, 1\}^{\mathbb{Z}}$ distributed according to a **Bernoulli product measure** ν_ρ with parameter $\rho \in (0, 1)$.



F.Solomon '75 proved LLN. A.Greven, F.den Hollander '94, F.Comets, N.Gantert and O.Zeitouni '00 proved LDPs.

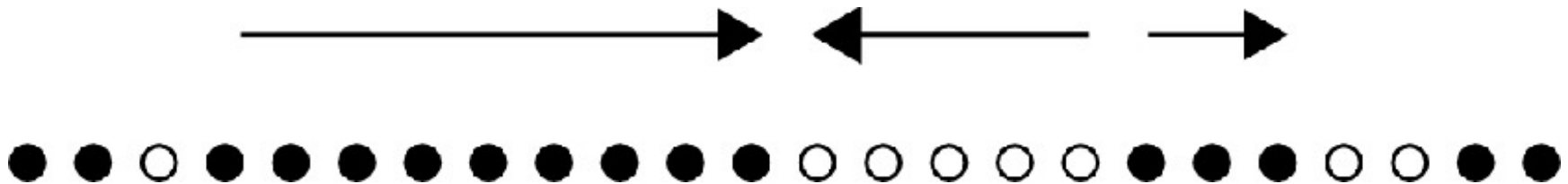
Slow-down phenomenon

When $\nu > 0$, the rate function is zero on $[0, \nu]$, namely, $\mathbb{P}\left(\frac{X_t}{t} \simeq \theta\right) \approx e^{-o(t)}$.



Traps

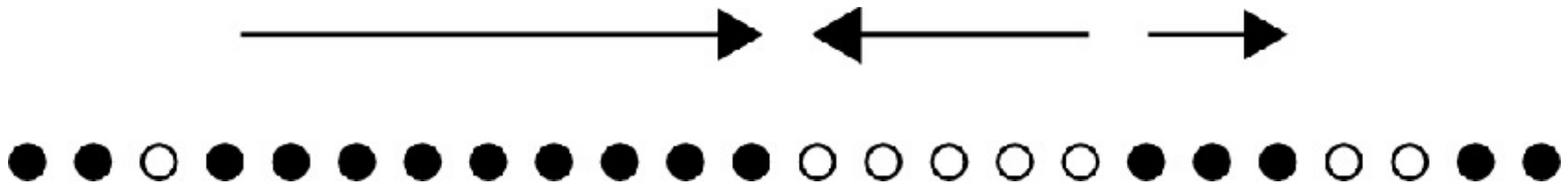
The slow-down phenomenon is due to the presence of traps.



Even though globally more particles than holes, presence of long strings of holes significantly slows down the walker.

Traps

The slow-down phenomenon is due to the presence of traps.



Even though globally more particles than holes, presence of long strings of holes significantly slows down the walker.

Does this slow-down survive in dynamic random environments?

LDP for the Empirical Speed: attractive spin-flip systems

Theorem 2 (Annealed LDP) The family $\mathbb{P}_{\mu,0}(X_t/t \in \cdot)$, $t > 0$, satisfies the LDP with rate t and with convex rate function $I^{\text{ann}}: \mathbb{R} \rightarrow [0, \infty)$ satisfying

$$I^{\text{ann}}(\theta) \begin{cases} = 0, & \text{if } \theta \in [v^-, v^+], \\ > 0, & \text{if } \theta \in \mathbb{R} \setminus [v^-, v^+], \end{cases}$$

for some $-(\alpha - \beta) \leq v^- \leq v \leq v^+ \leq \alpha - \beta$.

Moreover, for **small drift** $\alpha - \beta$ and **exponentially mixing** environments ,

$$v^- = v = v^+ \quad (\text{no slow-down phenomenon}).$$

Main ingredient of the proof: attractiveness \rightarrow FKG, subadditivity.

Theorem 3 (Quenched LDP) For P^μ -a.e. ξ , the family $P_0^\xi(X_t/t \in \cdot), t > 0$, satisfies the LDP with rate t and with deterministic convex rate function $I^{\text{que}}: \mathbb{R} \rightarrow [0, \infty)$ satisfying

$$I^{\text{que}}(\theta) \begin{cases} = 0, & \text{if } \theta \in [v^-, v^+], \\ > 0, & \text{if } \theta \in \mathbb{R} \setminus [v^-, v^+], \end{cases}$$

for some $-(\alpha - \beta) \leq v^- \leq v \leq v^+ \leq \alpha - \beta$.

Furthermore, $I^{\text{que}} \geq I^{\text{ann}}$, and the following symmetry relation holds:

$$I^{\text{que}}(-\theta) = I^{\text{que}}(\theta) + \theta(2\rho - 1) \log(\alpha/\beta), \quad \theta \geq 0.$$

Moreover, for **small drift** $\alpha - \beta$ and **exponentially mixing** environments:

$$v^- = v = v^+ \text{ (no slow-down phenomenon).}$$

Does the slow-down phenomenon hold only in the *static* random environment?

Does the slow-down phenomenon hold only in the *static* random environment?

No, if the dynamic random environment is not "fast mixing", then slow-down can survive.

Does the slow-down phenomenon hold only in the *static* random environment?

No, if the dynamic random environment is not "fast mixing", then slow-down can survive.

We can prove this when ξ is Simple Symmetric Exclusion (SSE).

RW on Simple Symmetric Exclusion (SSE)

Definition: an IPS ξ evolves according to a **SSE** dynamics if each particle performs an independent simple symmetric random walk at rate 1 such that jumps to occupied sites are suppressed.

Equilibrium measures: Bernoulli product measures ν_ρ with parameter $\rho \in (0, 1)$.

RW on Simple Symmetric Exclusion (SSE)

Definition: an IPS ξ evolves according to a **SSE** dynamics if each particle performs an independent simple symmetric random walk at rate 1 such that jumps to occupied sites are suppressed.

Equilibrium measures: Bernoulli product measures ν_ρ with parameter $\rho \in (0, 1)$.

Unfortunately: SSE is not cone-mixing and so we do not have LLN and LDP.

We performed simulations showing that LLN does hold with speed $\nu > 0$ as soon as $\rho > 1/2$.

We can prove subexponential decay for the probability that the displacement is sublinear.

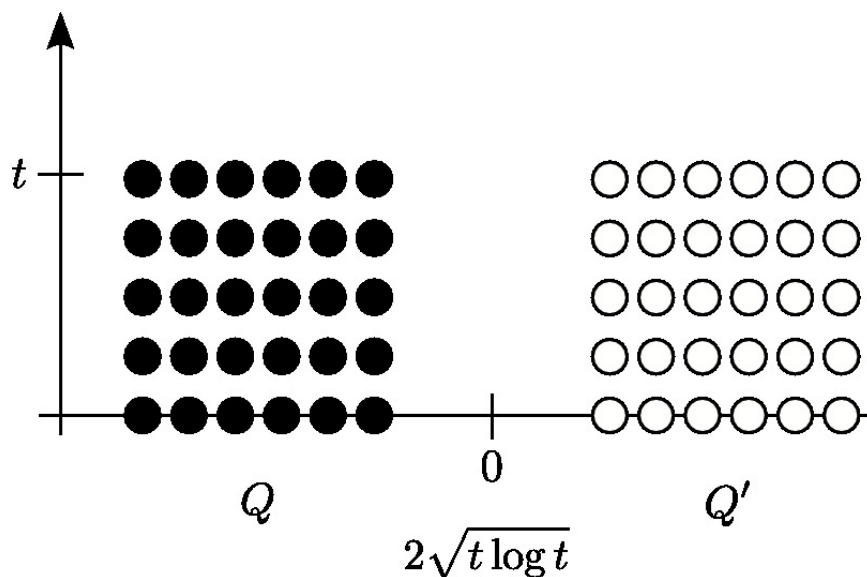
Proposition (Annealed slow-down for the SSE) For all $\rho \in (0, 1)$,

$$\lim_{t \rightarrow \infty} \frac{1}{t} \log \mathbb{P}_{\nu_\rho, 0}(|X_t| \leq 2\sqrt{t \log t}) = 0.$$

Key estimate for slow-down in the SSE case

There exist $C, \delta > 0$ such that, for all intervals $Q, Q' \subset \mathbb{Z}$ separated by a distance at least $2\sqrt{t \log t}$ and all $t \geq 1$,

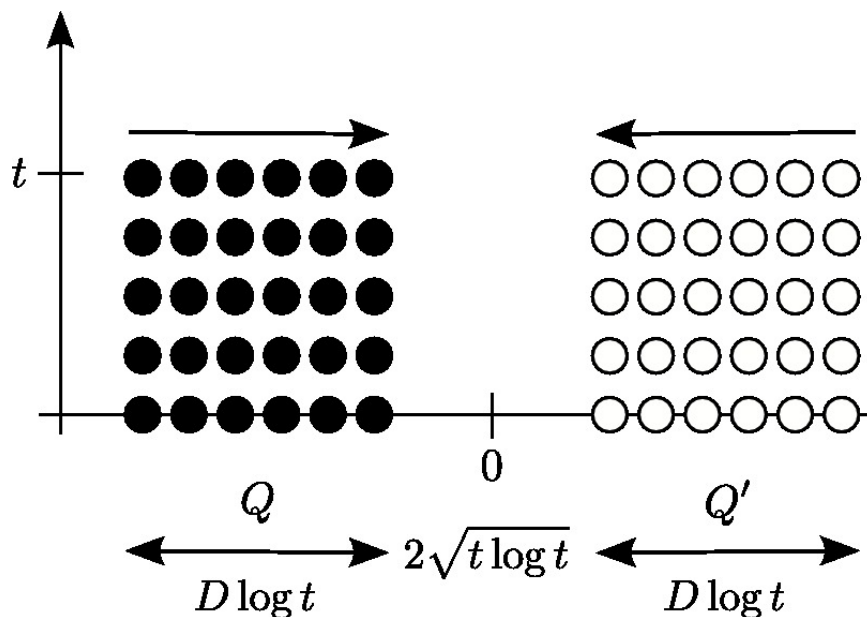
$$\mathbb{P}_{\mathbf{v}_\rho} \left(\xi(x, s) = 1, \xi(y, s) = 0 \forall x \in Q \forall y \in Q' \forall s \in [0, t] \right) \geq \delta e^{-C(|Q|+|Q'|)\sqrt{t}}.$$



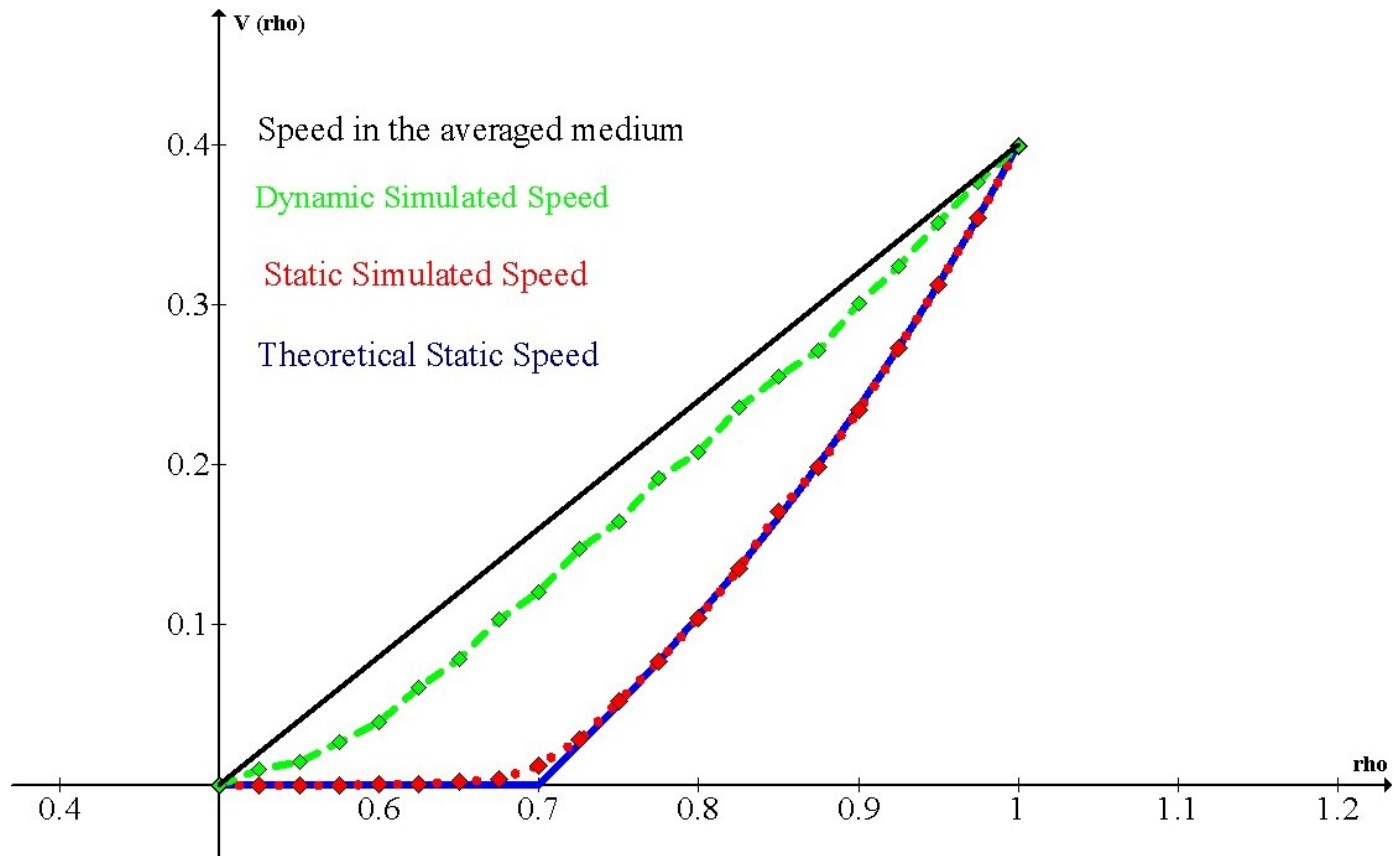
Key estimate for slow-down in the SSE case

There exist $C, \delta > 0$ such that, for all intervals $Q, Q' \subset \mathbb{Z}$ separated by a distance at least $2\sqrt{t \log t}$ and all $t \geq 1$,

$$\mathbb{P}_{\nu_\rho} \left(\xi(x, s) = 1, \xi(y, s) = 0 \forall x \in Q \forall y \in Q' \forall s \in [0, t] \right) \geq \delta e^{-C(|Q|+|Q'|)\sqrt{t}}.$$



Simulations: SSE



$$\alpha / (\alpha + \beta) = 0.7, \quad \alpha + \beta = 1.$$

Essential references

- ★ L. Avena, F. den Hollander and F. Redig, *Law of large numbers for a class of random walks in dynamic random environments*, EURANDOM Report 2009-032.
- ★ L. Avena, F. den Hollander and F. Redig, *Large deviation principle for one-dimensional random walk in dynamic random environment: attractive spin-flips and simple symmetric exclusion*, EURANDOM Report 2009-033.
- ★ F. Comets and O. Zeitouni, *A law of large numbers for random walks in random mixing environment*, Ann. Probab. 32 (2004) 880–914.
- ★ T.M. Liggett, *Interacting Particle Systems*, Grundlehren der Mathematischen Wissenschaften 276, Springer, New York, 1985.

Thanks



Leiden, Oude Rijn.