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# Scaling of a random walk on a supercritical contact process 

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#### Abstract

A proof is provided of a strong law of large numbers for a one-dimensional random walk in a dynamic random environment given by a supercritical contact process in equilibrium. The proof is based on a coupling argument that traces the space-time cones containing the infection clusters generated by single infections and uses that the random walk eventually gets trapped inside the union of these cones. For the case where the local drifts of the random walk are smaller than the speed at which infection clusters grow, the random walk eventually gets trapped inside a single cone. This in turn leads to the existence of regeneration times at which the random walk forgets its past. The latter are used to prove a functional central limit theorem and a large deviation principle.

The qualitative dependence of the speed, the volatility and the rate function on the infection parameter is investigated, and some open problems are mentioned.


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Key words and phrases. Random walk, dynamic random environment, contact process, strong law of large numbers, functional central limit theorem, large deviation principle, space-time cones, clusters of infections, coupling, regeneration times.

## 1 Introduction

### 1.1 Background, motivation and outline

Background. A random walk in a dynamic random environment on $\mathbb{Z}^{d}, d \geq 1$, is a random process where a "particle" makes random jumps with transition rates that depend on its location and themselves evolve with time. A typical example is when the dynamic random environment is given by an interacting particle system

$$
\begin{equation*}
\xi=\left(\xi_{t}\right)_{t \geq 0} \text { with } \xi_{t}=\left\{\xi_{t}(x): x \in \mathbb{Z}^{d}\right\} \in \Omega, \tag{1.1}
\end{equation*}
$$

[^0]where $\Omega$ is the configuration space, and $\xi_{0}$ is typically drawn from equilibrium. In the case where $\Omega=\{0,1\}^{\mathbb{Z}^{d}}$, the configurations can be thought of as consisting of "particles" and "holes". Given $\xi$, run a random walk $W=\left(W_{t}\right)_{t \geq 0}$ on $\mathbb{Z}^{d}$ that jumps at a fixed rate, but uses different transition kernels on a particle and on a hole. The key question is: What are the scaling properties of $W$ and how do these properties depend on the law of $\xi$ ?

The literature on random walks in dynamic random environments is still modest (for a recent overview, see Avena [1], Chapter 1). In Avena, den Hollander and Redig [4] a strong law of large numbers (SLLN) was proved for a class of interacting particle systems satisfying a mild space-time mixing condition, called cone-mixing. Roughly speaking, this is the requirement that for every $m>0$ all states inside the space-time cone (see Fig. 1)

$$
\begin{equation*}
\operatorname{CONE}_{t}:=\left\{(x, s) \in \mathbb{Z}^{d} \times[t, \infty):\|x\| \leq m(s-t)\right\} \tag{1.2}
\end{equation*}
$$

are conditionally independent of the states at time zero in the limit as $t \rightarrow \infty$. The proof of the SLLN uses a regeneration-time argument. Under a cone-mixing condition involving multiple cones, a functional central limit theorem (FCLT) can be derived as well, and under monotonicity conditions also a large deviation principle (LDP).


Figure 1: The cone defined in (1.2).

Many interacting particle systems are cone-mixing, including spin-flip systems with spin-flip rates that are weakly dependent on the configuration, e.g. the stochastic Ising model above the critical temperature. However, also many interacting particle systems are not cone-mixing, including independent simple random walks, the exclusion process, the contact process and the voter model. Indeed, these systems have slowly decaying space-time correlations. For instance, in the exclusion process particles are conserved and cannot sit on top of each other. Therefore, if at time zero there are particles everywhere in the box $\left[-t^{2}, t^{2}\right] \cap \mathbb{Z}^{d}$, then these particles form a "large traffic jam around the origin". This traffic jam will survive up to time $t$ with a probability tending to 1 as $t \rightarrow \infty$, and will therefore affect the states near the tip of $\mathrm{CONE}_{t}$. Similarly,
in the contact process, if at time zero there are no infections in the box $\left[-t^{2}, t^{2}\right] \cap \mathbb{Z}^{d}$, then no infections will be seen near the tip of $\mathrm{CONE}_{t}$ as well.

Motivation. Several attempts have been made to extend the SLLN to interacting particle systems that are not cone-mixing, with partial success. Examples include: independent simple random walks (den Hollander, Kesten and Sidoravicius [11]) and the exclusion process (Avena, dos Santos and Völlering [5], Avena [2]). The present paper considers the supercritical contact process. We exploit the graphical representation, which allows us to simultaneously couple all realizations of the contact process starting from different initial configurations. This coupling in turn allows us to first prove the SLLN when the initial configuration is "all infected" (with the help of a subadditivity argument), and then show that the same result holds when the initial configuration is drawn from equilibrium. The main idea is to use the coupling to show that configurations agree in large space-time cones containing the infection clusters generated by single infections and that the random walk eventually gets trapped inside the union of these cones.

Under the assumption that the local drifts of the random walk are smaller than the speed at which infection clusters grow, the random walk eventually gets trapped inside a single cone. We show that this implies the existence of regeneration times at which the random walk "forgets its past". The latter in turn allow us to prove the FCLT and the LDP.

It is typically difficult to obtain information about the speed in the SLLN, the volatility in the FCLT and the rate function in the LDP. In general, these are non-trivial functions of the parameters in the model, a situation that is well known from the literature on random walks in static random environments (for overviews, see Sznitman [16] and Zeitouni [17]). The reason is that these quantities depend on the environment process (i.e., the process of environments as seen from the location of the walk), which is typically hard to analyze. For the supercritical contact process we are able to derive a few qualitative properties as a function of the infection parameter, but it remains a challenge to obtain a full quantitative description.

A model of a random walk on the infinite cluster of supercritical oriented percolation (the discrete-time analogue of the contact process) is treated in Birkner, Černý, Depperschmidt and Gantert [8], where a SLLN and a quenched and annealed CLT are obtained. This model can be viewed as a random walk in a dynamic random environment, but it has non-elliptic transition probabilities different from the ones we consider here, because the random walk is confined to the infinite cluster.

Outline. In Section 1.2 we define the model. In Section 1.3 we state our main results: two theorems claiming the SLLN, the FCLT and the LDP under appropriate conditions on the model parameters. In Section 1.4 we mention some open problems. The proofs of the theorems are given in Sections 3 and 5 , respectively, Section 6. Sections 2 and 4 contain preparatory work.

### 1.2 Model

In this paper we consider the case where the dynamic random environment is the onedimensional linear contact process $\xi=\left(\xi_{t}\right)_{t \geq 0}$, i.e., the spin-flip system on $\Omega:=\{0,1\}^{\mathbb{Z}}$ with local transition rates given by

$$
\eta \rightarrow \eta^{x} \text { with rate } \begin{cases}1 & \text { if } \eta(x)=1  \tag{1.3}\\ \lambda\{\eta(x-1)+\eta(x+1)\} & \text { if } \eta(x)=0\end{cases}
$$

where $\lambda \in(0, \infty)$ and $\eta^{x}$ is defined by $\eta^{x}(y):=\eta(y)$ for $y \neq x, \eta^{x}(x):=1-\eta(x)$. We call a site infected when its state is 1 , and healthy when its state is 0. See Liggett [12], Chapter VI, for proper definitions.

The empty configuration $\mathbf{0} \in \Omega$, given by $\mathbf{0}(x)=0$ for all $x \in \mathbb{Z}$, is an absorbing state for $\xi$, while the full configuration $\mathbf{1} \in \Omega$, given by $\mathbf{1}(x)=1$ for all $x \in \mathbb{Z}$, evolves towards an equilibrium measure $\nu_{\lambda}$ that is stationary and ergodic under space-shifts. There is a critical threshold $\lambda_{c} \in(0, \infty)$ such that: (1) for $\lambda \in\left(0, \lambda_{c}\right], \nu_{\lambda}=\delta_{0}$; (2) for $\lambda \in\left(\lambda_{c}, \infty\right), \rho_{\lambda}:=\nu_{\lambda}(\eta(0)=1)>0$. In the latter case, $\delta_{0}$ and $\nu_{\lambda}$ are the only equilibrium measures. It is known that $\nu_{\lambda}$ has exponentially decaying correlations, and that $\lambda \mapsto \rho_{\lambda}$ is continuous and non-decreasing with $\lim _{\lambda \rightarrow \infty} \rho_{\lambda}=1$.

For a fixed realization of $\xi$, we define the random walk $W:=\left(W_{t}\right)_{t \geq 0}$ as the timeinhomogeneous Markov process on $\mathbb{Z}$ that, given $W_{t}=x$, jumps to

$$
\begin{array}{lll}
x+1 & \text { at rate } & \alpha_{1} \xi_{t}(x)+\alpha_{0}\left[1-\xi_{t}(x)\right], \\
x-1 & \text { at rate } & \beta_{1} \xi_{t}(x)+\beta_{0}\left[1-\xi_{t}(x)\right], \tag{1.4}
\end{array}
$$

where $\alpha_{i}, \beta_{i} \in(0, \infty), i=0,1$. We assume that

$$
\begin{equation*}
\alpha_{0}+\beta_{0}=\alpha_{1}+\beta_{1}=: \gamma, \tag{1.5}
\end{equation*}
$$

and that

$$
\begin{equation*}
v_{1}>v_{0} \text { with } v_{1}:=\alpha_{1}-\beta_{1} \text { and } v_{0}:=\alpha_{0}-\beta_{0}, \tag{1.6}
\end{equation*}
$$

i.e., the jump rate is constant and equal to $\gamma$ everywhere, while the drift to the right is larger on infected sites than on healthy sites. Observe that the assumption in (1.6) is made without loss of generality: since the contact process is invariant under reflection in the origin, $-W$ has the same law as $W$ with inverted jump rates.

### 1.3 Theorems

Let $\mathbb{P}_{\nu_{\lambda}}$ denote the joint law of $W$ and $\xi$ when the latter is started from $\nu_{\lambda}$. Our SLLN reads as follows.

Theorem 1.1. Suppose that (1.5-1.6) hold.
(a) For every $\lambda \in\left(\lambda_{c}, \infty\right)$ there exists a $v(\lambda) \in\left[v_{0}, v_{1}\right]$ such that

$$
\begin{equation*}
\lim _{t \rightarrow \infty} t^{-1} W_{t}=v(\lambda) \quad \mathbb{P}_{\nu_{\lambda}} \text { a.s. and in } L^{p}, p \geq 1 . \tag{1.7}
\end{equation*}
$$

(b) The function $\lambda \mapsto v(\lambda)$ is non-decreasing and right-continuous on $\left(\lambda_{c}, \infty\right)$, with $v(\lambda) \in\left(v_{0}, v_{1}\right)$ for all $\lambda \in\left(\lambda_{c}, \infty\right)$ and $\lim _{\lambda \rightarrow \infty} v(\lambda)=v_{1}$.

We note in passing that if $\lambda \in\left(0, \lambda_{c}\right)$, then $\xi_{t}$ agrees with $\mathbf{0}$ on an interval that grows exponentially fast in $t$ (Liggett [12], Chapter VI), and so it is trivial to deduce that $W$ satisfies the SLLN with $v(\lambda)=v_{0}$.

A FCLT and an LDP hold under an additional restriction, namely, $\lambda \in\left(\lambda_{W}, \infty\right)$ with

$$
\begin{equation*}
\lambda_{W}:=\inf \left\{\lambda \in\left(\lambda_{c}, \infty\right):\left|v_{0}\right| \vee\left|v_{1}\right|<\iota(\lambda)\right\} . \tag{1.8}
\end{equation*}
$$

Here, $\lambda \mapsto \iota(\lambda)$ is the infection propagation speed (see (2.4) in Section 2.1), which is known to be continuous, strictly positive and strictly increasing on $\left(\lambda_{c}, \infty\right)$, with $\lim _{\lambda \downarrow \lambda_{c}} \iota(\lambda)=0$ and $\lim _{\lambda \rightarrow \infty} \iota(\lambda)=\infty$.

Theorem 1.2. Suppose that (1.5-1.6) hold.
(a) For every $\lambda \in\left(\lambda_{W}, \infty\right)$ there exists a $\sigma(\lambda) \in(0, \infty)$ such that, under $\mathbb{P}_{\nu_{\lambda}}$,

$$
\begin{equation*}
\left(\frac{W_{n t}-v(\lambda) n t}{\sigma(\lambda) \sqrt{n}}\right)_{t \geq 0} \Longrightarrow\left(B_{t}\right)_{t \geq 0} \quad \text { as } n \rightarrow \infty \tag{1.9}
\end{equation*}
$$

where $B$ is standard Brownian motion and $\Longrightarrow$ denotes weak convergence in path space.
(b) The functions $\lambda \mapsto v(\lambda)$ and $\lambda \mapsto \sigma(\lambda)$ are continuous on $\left(\lambda_{W}, \infty\right)$.
(c) For every $\lambda \in\left(\lambda_{W}, \infty\right)$, $\left(t^{-1} W_{t}\right)_{t>0}$ under $\mathbb{P}_{\nu_{\lambda}}$ satisfies the large deviation principle on $\mathbb{R}$ with a finite and convex rate function that has a unique zero at $v(\lambda)$.

The intuitive reason why the rate function has a unique zero is that deviations of the empirical speed in the: (i) upward direction require a density of infected sites larger than $\rho_{\lambda}$, which is costly because infections become healthy independently of the states at the other sites; (ii) downward direction require a density of infected sites smaller than $\rho_{\lambda}$, which is costly because infection clusters grow at a linear speed and rapidly fill up healthy intervals everywhere.

### 1.4 Discussion

1. It is natural to expect that $\lambda \mapsto v(\lambda)$ is continuous and strictly increasing on $\left(\lambda_{c}, \infty\right)$ with $\lim _{\lambda \downarrow \lambda_{c}} v(\lambda)=v_{0}$. Fig. 2 shows a qualitative plot of the speed in that setting. If $0 \in\left(v_{0}, v_{1}\right)$, then there is a critical threshold $\lambda^{*} \in\left(\lambda_{c}, \infty\right)$ at which the speed changes sign. It is natural to ask whether $\lambda \mapsto v(\lambda)$ is concave on $\left(\lambda_{c}, \infty\right)$ and Lipshitz at $\lambda_{c}$.
2. We know that $W$ is transient when $v(\lambda) \neq 0$. Is $W$ recurrent when $v(\lambda)=0$ ?
3. We expect (1.8) to be redundant. Moreover, we expect that for every $\lambda \in\left(\lambda_{c}, \infty\right)$ the environment process (i.e., the process of environments as seen from the location of the random walk) has a unique and non-trivial equilibrium measure that is absolutely continuous with respect to $\nu_{\lambda}$.
4. Theorems $1.1-1.2$ can presumably be extended to $\mathbb{Z}^{d}$ with $d \geq 2$. Also in higher dimensions single infections create infection clusters that grow at a linear speed (i.e., asymptotically form a ball with a linearly growing radius). The construction of the regeneration times when $\lambda \in\left(\lambda_{W}, \infty\right)$, with $\lambda_{W}$ the analogue of (1.8), is straightforward.


Figure 2: Qualitative plot of $\lambda \mapsto v(\lambda)$ when $0 \in\left(v_{0}, v_{1}\right)$.
5. It would be interesting to extend Theorems 1.1-1.2 to multi-type contact processes. On each type $i$ the random walk has transition rates $\alpha_{i}, \beta_{i}$ such that $\alpha_{i}+\beta_{i}=\gamma$ for all $i$. As long as the dynamics is monotone and $i \mapsto v_{i}$ is non-decreasing, many of the arguments in the present paper carry over.

## 2 Construction

In Section 2.1 we construct the contact process, in Section 2.2 the random walk on top of the contact process.

### 2.1 Contact process

A càdlàg version of the contact process can be constructed from a graphical representation in the following standard fashion. Let $:=(H(x))_{x \in \mathbb{Z}}$ and $I:=(I(x))_{x \in \mathbb{Z}}$ be two independent collections of i.i.d. Poisson processes with rates 1 and $\lambda$, respectively. On $\mathbb{Z} \times[0, \infty)$, draw the events of $H(x)$ as crosses over $x$ and the events of $I(x)$ as two-sided arrows between $x$ and $x+1$ (see Fig. 3).
(The standard graphical representation uses Poisson processes of one-sided arrows to the right and to the left on every time line, each with rate $\lambda$. This gives the same dynamics.)

For $x, y \in \mathbb{Z}$ and $0 \leq s \leq t$, we say that $(x, s)$ and $(y, t)$ are connected, written $(x, s) \leftrightarrow(y, t)$, if and only if there exists a nearest-neighbor path in $\mathbb{Z} \times[0, \infty)$ starting at $(x, s)$ and ending at $(y, t)$, going either upwards in time or sideways in space across arrows without hitting crosses. For $x \in \mathbb{Z}$, we define the cluster of $x$ at time $t$ by

$$
\begin{equation*}
C_{t}(x):=\{y \in \mathbb{Z}:(x, 0) \leftrightarrow(y, t)\} . \tag{2.1}
\end{equation*}
$$

For example, in Fig. 3, $C_{t}(0)=\{-2,-1,1,2\}$ and $C_{t}(2)=\emptyset$. Note that $C_{t}(x)$ is a function of $H$ and $I$.


Figure 3: Graphical representation. The crosses are events of $H$ and the arrows are events of $I$. The thick lines cover the region that is infected when the initial configuration has a single infection at the origin.

For a fixed initial configuration $\eta$, we declare $\xi_{t}(y)=1$ if there exists an $x$ such that $y \in C_{t}(x)$ and $\eta(x)=1$, and we declare $\xi_{t}(y)=0$ otherwise. Then $\xi$ is adapted to the filtration

$$
\begin{equation*}
\mathcal{F}_{t}:=\sigma\left(\xi_{0},\left(H_{s}, I_{s}\right)_{s \in[0, t]}\right) \tag{2.2}
\end{equation*}
$$

This construction allows us to simultaneously couple copies of the contact process starting from all configurations $\eta \in \Omega$. In the following we will write $\xi(\eta)$ and $\xi_{t}(\eta)(x)$ when we want to exhibit that the initial configuration is $\eta$.

We note two consequences of the graphical construction, stated in Lemmas 2.1-2.3 below. The first is the monotonicity of $\eta \mapsto \xi(\eta)$, the second concerns the state of the sites surrounded by the cluster of an infected site. The notation $\eta \leq \eta^{\prime}$ stands for $\eta(x) \leq \eta^{\prime}(x)$ for all $x \in \mathbb{Z}$.

Lemma 2.1. If $\eta \leq \eta^{\prime}$, then $\xi_{t}(\eta) \leq \xi_{t}\left(\eta^{\prime}\right)$ for all $t \geq 0$.
Proof. Immediate from the definition of $\xi_{t}$ in terms of $\eta$ and $\left(C_{t}(x)\right)_{x \in \mathbb{Z}}$.
For $x \in \mathbb{Z}$, define the left-most and the right-most site influenced by site $x$ at time $t$ as

$$
\begin{align*}
& L_{t}(x):=\inf C_{t}(x), \\
& R_{t}(x):=\sup C_{t}(x), \tag{2.3}
\end{align*}
$$

where $\inf \emptyset=\infty$ and $\sup \emptyset=-\infty$. By symmetry, for any $t \geq 0, R_{t}(x)-x$ and $x-L_{t}(x)$ have the same distribution, independently of $x$.

Lemma 2.2. Fix $x \in \mathbb{Z}$ and $t \geq 0$. If $C_{t}(x) \neq \emptyset$ and $y \in\left[L_{t}(x), R_{t}(x)\right] \cap \mathbb{Z}$, then $\eta \mapsto \xi_{t}(\eta)(y)$ is constant on $\{\eta \in \Omega: \eta(x)=1\}$.

Proof. It suffices to show that, under the conditions stated, $\xi_{t}(\eta)(y)=1$ if and only if $y \in C_{t}(x)$. The 'if' part is obvious. For the 'only if' part, note that if there is a $z \neq x$ such that $(z, 0) \leftrightarrow(y, t)$, then any path realizing the connection must cross a path connecting $(x, 0)$ to either $\left(R_{t}(x), t\right)$ or $\left(L_{t}(x), t\right)$, so that $(x, 0) \leftrightarrow(y, t)$ as well.

If $\xi_{0}=\mathbb{1}_{x}$, then $R_{t}(x)$ and $L_{t}(x)$ are, respectively, the right-most and the left-most infections present at time $t$. In particular, in this case the infection survives for all times if and only if $R_{t}(x)-L_{t}(x) \geq 0$ for all $t \geq 0$. For $\lambda \in\left(\lambda_{c}, \infty\right)$ it is well known that, given $\xi_{0}=\mathbb{1}_{0}$, the infection survives with positive probability and there exists a constant $\iota=\iota(\lambda)>0$ such that, conditionally on survival,

$$
\begin{equation*}
\lim _{t \rightarrow \infty} t^{-1} R_{t}(0)=\iota \quad \xi \text {-a.s. } \tag{2.4}
\end{equation*}
$$

### 2.2 Random walk on top of contact process

Under assumptions (1.5-1.6), the random walk $W$ can be constructed as follows. Let $N:=\left(N_{t}\right)_{t \geq 0}$ be a Poisson process with rate $\gamma$. Denote by $J:=\left(J_{k}\right)_{k \in \mathbb{N}_{0}}$ its generalized inverse, i.e., $J_{0}=0$ and $\left(J_{k+1}-J_{k}\right)_{k \in \mathbb{N}_{0}}$ are i.i.d. $\operatorname{EXP}(\gamma)$ random variables. Let $U:=\left(U_{k}\right)_{k \in \mathbb{N}}$ be an i.i.d. sequence of $\operatorname{UNIF}([0,1])$ random variables, independent of $N$. Set $S_{0}:=0$ and, recursively for $k \in \mathbb{N}_{0}$,

$$
\begin{equation*}
S_{k+1}:=S_{k}+2\left(\mathbb{1}_{\left\{0 \leq U_{k+1} \leq \alpha_{0} / \gamma\right\}}+\xi_{J_{k+1}}\left(S_{k}\right) \mathbb{1}_{\left\{\alpha_{0} / \gamma<U_{k+1} \leq \alpha_{1} / \gamma\right\}}\right)-1, \tag{2.5}
\end{equation*}
$$

i.e., $S_{k+1}=S_{k}+1$ with probability $\alpha_{i} / \gamma$ and $S_{k+1}=S_{k}-1$ with probability $\beta_{i} / \gamma=$ $1-\alpha_{i} / \gamma$ when $\xi_{J_{k+1}}\left(S_{k}\right)=i$, for $i=0,1$ (recall that $\alpha_{0}<\alpha_{1}$ by (1.5-1.6)). Setting

$$
\begin{equation*}
W_{t}:=S_{N_{t}}, \tag{2.6}
\end{equation*}
$$

we can use the right-continuity of $\xi$ to verify that $W$ indeed is a Markov process with the correct jump rates.

A useful property of the above construction is that it is monotone in the environment, in the following sense. For two dynamic random environments $\xi$ and $\xi^{\prime}$, we say that $\xi \leq \xi^{\prime}$ when $\xi_{t} \leq \xi_{t}^{\prime}$ for all $t \geq 0$. Writing $W=W(\xi)$ in the previous construction (i.e., exhibiting $W$ as a function of $\xi$ ), it is clear from (2.5) that

$$
\begin{equation*}
\xi \leq \xi^{\prime} \quad \Longrightarrow \quad W_{t}(\xi) \leq W_{t}\left(\xi^{\prime}\right) \quad \forall t \geq 0 . \tag{2.7}
\end{equation*}
$$

We denote by

$$
\begin{equation*}
\mathcal{G}_{t}:=\mathcal{F}_{t} \vee \sigma\left(\left(N_{s}\right)_{s \in[0, t]},\left(U_{k}\right)_{1 \leq k \leq N_{t}}\right) \tag{2.8}
\end{equation*}
$$

the filtration generated by all the random variables that are used to define the contact process $\xi$ and the random walk $W$.

## 3 SLLN

Theorem 1.1(a) is proved in two steps. In Section 3.1 we use subadditivity to prove the SLLN when $\xi$ starts from $\delta_{\mathbf{1}}$. In Section 3.2 we couple two copies of $\xi$ starting from $\nu_{\lambda}$ and $\delta_{\mathbf{1}}$, transfer the SLLN, and show that the speed is the same.

In the following, for a random process $X=\left(X_{t}\right)_{t \in \mathcal{I}}$ with $\mathcal{I}=\mathbb{R}$ or $\mathcal{I}=\mathbb{Z}$, we write

$$
\begin{equation*}
X_{[0, t]}:=\left(X_{s}\right)_{s \in[0, t] \cap \mathcal{I}} . \tag{3.1}
\end{equation*}
$$

### 3.1 Starting from the full configuration: subadditivity

Since $\eta \leq \mathbf{1}$ for all $\eta \in \Omega$, it follows from (2.7) and Lemma 2.1 that $W_{t}(\xi(\eta)) \leq W_{t}(\xi(\mathbf{1}))$ for all $t \geq 0$. Therefore, if in the graphical construction we replace $\xi_{s}$ by $\mathbf{1}$ at any given time $s$, then the new increments after time $s$ lie to the right of the old increments after time $s$, and are independent of the increments before time $s$. This leads us to a subadditivity argument, which we now formalize.

For $n \in \mathbb{N}_{0}$, let

$$
\begin{align*}
H^{(n)} & =\left(H_{t}^{(n)}(x)\right)_{t \geq 0, x \in \mathbb{Z}} \\
I^{(n)} & =\left(I_{t}^{(n)}(x)\right)_{t \geq 0, x \in \mathbb{Z}} \\
N^{(n)} & \left.:=\left(H_{t+n}\left(x+W_{n}\right)-H_{n}\left(x+W_{n}\right)\right)_{t \geq 0, x \in \mathbb{Z}}\left(x+W_{n}\right)-I_{n}\left(x+W_{n}\right)\right)_{t \geq 0, x \in \mathbb{Z}},  \tag{3.2}\\
U^{(n)} & =\left(U_{k}^{(n)}\right)_{k \geq 0} \\
U_{k \in \mathbb{N}} & :=\left(N_{t+n}-N_{n}\right)_{t \geq 0}, \\
& :=\left(U_{k+N_{n}}\right)_{k \in \mathbb{N}} .
\end{align*}
$$

Then, for any $n \in \mathbb{N}_{0},\left(H^{(n)}, I^{(n)}, N^{(n)}, U^{(n)}\right)$ has the same distribution as $(H, I, N, U)$ and is independent of

$$
\begin{equation*}
H_{[0, n-j]}^{(j)}, I_{[0, n-j]}^{(j)}, N_{[0, n-j]}^{(j)}, U_{\left[1, N_{n-j}^{(j)}\right]}^{(j)}, \quad 0 \leq j \leq n-1 . \tag{3.3}
\end{equation*}
$$

Abbreviate $\xi=\xi(\eta, H, I)$ and $W=W(\xi, N, U)$. For $n \in \mathbb{N}_{0}$, let

$$
\begin{align*}
& \xi^{(n)}:=\xi\left(\mathbf{1}, H^{(n)}, I^{(n)}\right), \\
& W^{(n)}:=W\left(\xi^{(n)}, N^{(n)}, U^{(n)}\right), \tag{3.4}
\end{align*}
$$

and define the double-indexed sequence

$$
\begin{equation*}
X_{m, n}:=W_{n-m}^{(m)}, \quad n, m \in \mathbb{N}_{0}, n \geq m . \tag{3.5}
\end{equation*}
$$

Lemma 3.1. The following properties hold:
(i) For all $n, m \in \mathbb{N}_{0}, n \geq m: X_{0, n} \leq X_{0, m}+X_{m, n}$.
(ii) For all $n \in \mathbb{N}_{0}:\left(X_{n, n+k}\right)_{k \in \mathbb{N}_{0}}$ has the same distribution as $\left(X_{0, k}\right)_{k \in \mathbb{N}_{0}}$.
(iii) For all $k \in \mathbb{N}$ : $\left(X_{n k,(n+1) k}\right)_{n \in \mathbb{N}_{0}}$ is i.i.d.
(iv) $\sup _{n \in \mathbb{N}} \mathbb{E}_{\delta_{1}}\left[n^{-1}\left|X_{0, n}\right|\right]<\infty$.

Proof. (i) Fix $n, m \in \mathbb{N}_{0}, n \geq m$ and define $\hat{\xi}:=\xi\left(\hat{\eta}, H^{(m)}, I^{(m)}\right)$, where $\hat{\eta}(x)=\xi_{m}(x+$ $\left.W_{m}\right)$. This is the contact process after time $m$ as seen from $W_{m}$. Note that $X_{0, n}-X_{0, m}=$ $W_{n}-W_{m}=W_{n-m}\left(\hat{\xi}, N^{(m)}, U^{(m)}\right)$. Since $\hat{\eta} \leq \mathbf{1}$, it follows from (2.7) and Lemma 2.1 that the latter is $\leq W_{n-m}\left(\xi^{(m)}, N^{(m)}, U^{(m)}\right)=W_{n-m}^{(m)}$.
(ii) Immediate from the construction.
(iii) By definition, $X_{n k,(n+1) k}=W_{k}\left(\xi^{(n k)}, N^{(n k)}, U^{(n k)}\right)$. By construction, for each $t \geq 0$, $W_{t}(\xi, N, U)$ is a function of $N_{[0, t]}, U_{\left[1, N_{t}\right]}$ and $\xi_{[0, t]}$, which in turn is a function of $H_{[0, t]}, I_{[0, t]}$ and $\eta$. Therefore $X_{n k,(n+1) k}$ is equal to a (fixed) function of

$$
\begin{equation*}
H_{[0, k]}^{(n k)}, I_{[0, k]}^{(n k)}, N_{[0, k]}^{(n k)}, U_{\left[1, N_{(n+1) k}^{(n k)}\right]}^{(n k)}, \tag{3.6}
\end{equation*}
$$

which are jointly i.i.d. in $n$ (when $k$ is fixed).
(iv) This follows from the fact that $\left|W_{t}\right| \leq N_{t}$.

Lemma 3.1 allows us to prove the SLLN when $\xi$ starts from $\delta_{\mathbf{1}}$.
Proposition 3.2. Let

$$
\begin{equation*}
v(\lambda):=\inf _{n \in \mathbb{N}} \mathbb{E}_{\delta_{1}}\left[n^{-1} W_{n}\right] . \tag{3.7}
\end{equation*}
$$

Then

$$
\begin{equation*}
\lim _{t \rightarrow \infty} t^{-1} W_{t}=v(\lambda) \quad \mathbb{P}_{\delta_{1}} \text {-a.s. and in } L^{p}, p \geq 1 \tag{3.8}
\end{equation*}
$$

Proof. Conditions (i)-(iv) in Lemma 3.1 allow us to apply the subbaditive ergodic theorem of Liggett [13] (see also Liggett [12], Theorem VI.2.6) to the sequence $\left(X_{0, n}\right)_{n \in \mathbb{N}_{0}}=$ $\left(W_{n}\right)_{n \in \mathbb{N}_{0}}$, which gives $\lim _{n \rightarrow \infty} n^{-1} W_{n}=v \mathbb{P}_{\delta_{1}}$-a.s. Via a standard argument this can subsequently be extended to $\left(t^{-1} W_{t}\right)_{t \geq 0}$ by using that, for any $n \in \mathbb{N}_{0}$,

$$
\begin{equation*}
\sup _{s \in[0,1]}\left|W_{n+s}-W_{n}\right| \leq N_{n+1}-N_{n} \tag{3.9}
\end{equation*}
$$

which implies that $\lim _{t \rightarrow \infty} t^{-1}\left|W_{t}-W_{\langle t\rfloor}\right|=0 \mathbb{P}_{\delta_{1}}$-a.s. The convergence also holds in $L^{p}$, because $\left|W_{t}\right| \leq N_{t}$ and so $\left(t^{-p}\left|W_{t}\right|^{p}\right)_{t \geq 1}$ is uniformly integrable for any $p \geq 1$.

### 3.2 Starting from equilibrium: coupling

In this section we show that two copies of the contact process starting from $\nu_{\lambda}$ and $\delta_{1}$ and coupled via the graphical representation are with a large probability equal inside space-time cones with tips at large times. Since the random walk eventually gets trapped inside a dense union of such cones, this will be enough to transfer the result of Proposition 3.2 from $\mathbb{P}_{\delta_{1}}$ to $\mathbb{P}_{\nu_{\lambda}}$, with the same velocity $v(\lambda)$, and will complete the proof of Theorem 1.1(a).

For $m, r>0$ and $t \geq 0$, let

$$
\begin{equation*}
V_{m, r}(t):=\{(x, s) \in \mathbb{Z} \times[t, \infty):|x| \leq r \vee m(s-t)\}, \tag{3.10}
\end{equation*}
$$

i.e., $V_{m, r}(t)$ is the union of the cylinder $[-r, r] \cap \mathbb{Z} \times[t, \infty)$ and the cone with tip at $(0, t)$ opening upwards in space-time with inclination $m$ (recall (1.2)).

Let $\eta$ be distributed according to $\nu_{\lambda}$, and let $\xi^{(1)}:=\xi(\eta), \xi^{(2)}:=\xi(\mathbf{1})$, i.e., take $\xi^{(1)}$ and $\xi^{(2)}$ to be copies of the contact process constructed from the same graphical representation and initial configurations $\eta$ and $\mathbf{1}$, respectively. Denote by $\mathbb{P}$ the joint distribution of all random variables needed to define $\xi^{(1)}, \xi^{(2)}$ and $W$, i.e., $\mathbb{P}$ is the product of the distributions of $\eta, H, I, N$ and $U$.

Lemma 3.3. For any $m, r>0$,

$$
\begin{equation*}
\lim _{T \rightarrow \infty} \mathbb{P}\left(\exists(x, t) \in V_{m, r}(T): \xi_{t}^{(1)}(x) \neq \xi_{t}^{(2)}(x)\right)=0 \tag{3.11}
\end{equation*}
$$

Before proving Lemma 3.3, we show how it leads to Theorem 1.1(a).

Proof of Theorem 1.1(a). Fix $\epsilon>0$. Let $D_{T}^{1}(r):=\left\{N_{T+t}-N_{T} \leq r \vee 2 \gamma t \forall t \geq 0\right\}$. Since $\lim _{t \rightarrow \infty} t^{-1} N_{t}=\gamma$ a.s. and $\left(N_{T+t}-N_{T}\right)_{t \geq 0}$ is equal in distribution to $N$, there exists an $r_{0}>0$ such that

$$
\begin{equation*}
\mathbb{P}\left(D_{T}^{1}\left(r_{0}\right)\right) \geq 1-\frac{1}{2} \epsilon \quad \forall T>0 . \tag{3.12}
\end{equation*}
$$

Let $D_{T}^{2}:=\left\{\xi_{t}^{(1)}(x)=\xi_{t}^{(2)}(x) \forall(x, t) \in V_{2 \gamma, r_{0}}(T)\right\}$ and $D_{T}:=D_{T}^{1}\left(r_{0}\right) \cap D_{T}^{2}$. By (3.12) and Lemma 3.3, there exists a $T_{0}>0$ large enough such that

$$
\begin{equation*}
\mathbb{P}\left(D_{T_{0}}\right)>1-\epsilon . \tag{3.13}
\end{equation*}
$$

Let $\Gamma_{0}:=\left\{N_{T_{0}}=0\right\}$, which has positive probability and is independent of $\xi^{(i)}, i=$ 1,2. Let $W^{(i)}:=W\left(\xi^{(i)}\right), i=1,2$. Note that $W^{(1)}=W^{(2)}$ on $\Gamma_{0} \cap D_{T_{0}}$. Since $\lim _{t \rightarrow \infty} t^{-1} W_{t}^{(2)}=v(\lambda) \mathbb{P}$-a.s., we therefore get

$$
\begin{equation*}
\mathbb{P}\left(\lim _{t \rightarrow \infty} t^{-1}\left(W_{t+T_{0}}^{(1)}-W_{T_{0}}^{(1)}\right)=v \mid \Gamma_{0}\right) \geq 1-\epsilon . \tag{3.14}
\end{equation*}
$$

However, because $\nu_{\lambda}$ is an equilibrium and $W_{T_{0}}^{(1)}=0$ on $\Gamma_{0},\left(W_{t+T_{0}}^{(1)}-W_{T_{0}}^{(1)}\right)_{t \geq 0}$ has under $\mathbb{P}\left(\cdot \mid \Gamma_{0}\right)$ the same distribution as $W$ under $\mathbb{P}_{\nu_{\lambda}}$, so the SLLN is obtained by letting $\epsilon \downarrow 0$. Convergence in $L^{p}, p \geq 1$, follows as in the proof of Proposition 3.2.

Proof of Lemma 3.3. Denote by $P$ the joint law of $\eta, H$ and $I$. The law of $\left(\xi^{(1)}, \xi^{(2)}\right)$ is the same under $P$ or $\mathbb{P}$. We can regard $P$ as a law on the product space

$$
\begin{equation*}
\left(\{0,1\} \times D\left(\mathbb{N}_{0},[0, \infty)\right)^{2}\right)^{\mathbb{Z}}=\{0,1\}^{\mathbb{Z}} \times\left(D\left(\mathbb{N}_{0},[0, \infty)\right)^{2}\right)^{\mathbb{Z}} \tag{3.15}
\end{equation*}
$$

$P$ is shift-ergodic because it is the product of probability measures that are shift-ergodic, namely, $\nu_{\lambda}$ and the distributions of $H$ and $I$. Let

$$
\begin{equation*}
\Lambda_{x}:=\left\{\eta(x)=1,\left(x-L_{t}(x)\right) \wedge\left(R_{t}(x)-x\right) \geq\lfloor(\iota / 2) t\rfloor \forall t \geq 0\right\}, \tag{3.16}
\end{equation*}
$$

i.e., the event that $x$ generates a "wide-spread infection" (moving at speed at least half the typical asymptotic speed $\iota$ ). Since $\Lambda_{x}$ is a translation of $\Lambda_{0}$, we have

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{1}{n} \sum_{x=1}^{n} \mathbb{1}_{\Lambda_{x}}=P\left(\Lambda_{0}\right)=: \varrho>0 \quad P \text {-a.s. } \tag{3.17}
\end{equation*}
$$

where the last inequality is justified by (2.4) and local modifications of the graphical representation.

Next, for $n \in \mathbb{N}$, define $Z_{n}$ by the equation

$$
\begin{equation*}
\sum_{x=1}^{Z_{n}} \mathbb{1}_{\Lambda_{x}}=n \tag{3.18}
\end{equation*}
$$

Then we also have

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \frac{Z_{n}}{n}=\varrho^{-1} \quad P \text {-a.s. } \tag{3.19}
\end{equation*}
$$

$\left(Z_{n}\right)_{n \in \mathbb{N}}$ marks the positions of wide-spread infections to the right of the origin, i.e., $x>0$ such that $\Lambda_{x}$ occurs. Equation (3.19) means that these wide-spread infections are not too far apart. Extending the definition of $Z_{n}$ to the negative integers, we obtain analogously that $\lim _{n \rightarrow \infty} n^{-1}\left(-Z_{-n}\right)=\varrho^{-1} P$-a.s. Let $\mathcal{Z}:=\cup_{n \in \mathbb{N}}\left\{Z_{n}, Z_{-n}\right\}$ and

$$
\begin{equation*}
\mathcal{S}:=\{(y, t) \in \mathbb{Z} \times[2 / \iota, \infty): \exists x \in \mathcal{Z} \text { such that }|y-x| \leq(\iota / 2) t-1\} \tag{3.20}
\end{equation*}
$$

Then $\mathcal{S}$ is the union of cones of inclination angle $\iota / 2$ with tips at $(2 / \iota, z)$ with $z \in \mathcal{Z}$ (see Fig. 4). We call $\mathcal{S}$ the safe region. This is justified by the following fact, whose proof is a direct consequence of Lemma 2.2.
Lemma 3.4. If $(x, t) \in \mathcal{S}$, then $\xi_{t}^{(1)}(x)=\xi_{t}^{(2)}(x)$.


Figure 4: Cones have inclination angle $\iota / 2$. The safe region $\mathcal{S}$ lies above the thick lines.
By Lemma 3.4, it is enough to prove that $\mathcal{S}$ contains $V_{m, r}(t)$ with a large probability when $t$ is large. Instead, we will prove that, for any $m>0$,

$$
\begin{equation*}
V_{m, 0}(0) \cap \mathcal{S}^{c} \text { is a bounded subset of } \mathbb{Z} \times[0, \infty) \quad P \text {-a.s. } \tag{3.21}
\end{equation*}
$$

This will also be enough, because it implies that $V_{m, r}(t) \subset \mathcal{S}$ for large $t, P$-a.s. for any $r>0$,

Now, $\mathcal{S}^{c}$ is contained in the union of space-time "houses" (unions of triangles and rectangles) with base at time 0 . The tips of the houses to the right of 0 form a sequence with spatial coordinates $\frac{1}{2}\left(Z_{n+1}+Z_{n}\right)$ and temporal coordinates $\left(Z_{n+1}-Z_{n}+2\right) / \iota$, $n \in \mathbb{N}$. By (3.19), the ratio of temporal/spatial coordinates tends to 0 as $n \rightarrow \infty$, so that only finitely many tips can be inside $V_{m, 0}(0)$. The same is true for the tips of the houses to the left of 0 . Therefore $V_{m, 0}(0)$ touches only finitely many houses, which proves (3.21).

## 4 More on the contact process

In this section we collect some additional facts about the contact process on $\mathbb{Z}$ that will be needed in the remainder of the paper. The proofs rely on geometric observations that will also illuminate the proof strategies developed in Sections 5-6.

In the following we will use the notation

$$
\begin{equation*}
\mathbb{Z}_{\leq x}:=\mathbb{Z} \cap(\infty, x] \tag{4.1}
\end{equation*}
$$

and analogously for $\mathbb{Z}_{\geq x}$.

Stochastic domination. We start with a useful alternative construction of the equilibrium $\nu_{\lambda}$. Let $\eta(x):=\mathbb{1}_{\left\{C_{t}(x) \neq \emptyset \forall t \geq 0\right\}}$. Then, by the graphical representation, $\eta$ has distribution $\nu_{\lambda}$. This follows from duality (see Liggett [12], Chapter VI). We can also graphically construct the contact process starting from $\nu_{\lambda}$ : extend the graphical representation to negative times, and declare $\xi_{t}(x)=1$ if and only if for all $0 \leq s \leq t$ there exists a $y$ such that $(y, s) \leftrightarrow(x, t)$, i.e., if and only if there exists an infinite infection path going backwards in time from $(x, t)$.

Let $\bar{\nu}_{\lambda}$ denote the restriction of $\nu_{\lambda}$ to $\mathbb{Z}_{\leq-1}$. Abusing notation, we will write the same symbol to denote the measure on $\Omega$ that is the product of $\bar{\nu}_{\lambda}$ with the measure concentrated on all sites healthy to the right of -1 . Using the alternative construction above, we can prove that the restriction of $\nu_{\lambda}(\cdot \mid \eta(0)=1)$ to $\mathbb{Z}_{\leq-1}$ is stochastically larger than $\bar{\nu}_{\lambda}$. In the following, we will focus on a similar result for the distribution of $\xi_{t}$ to the left of certain infection paths.

For $\varpi_{[0, t]}$ a nearest-neighbor càdlàg path with values in $\mathbb{Z}$, let

$$
\begin{equation*}
\overline{\mathcal{R}}_{t}^{\varpi}:=\sigma\left(\left(\xi_{0}(x)\right)_{x \geq \varpi_{0}},\left(H_{s}(x), I_{s}(x)\right)_{s \in[0, t], x \geq \varpi_{s}}\right) . \tag{4.2}
\end{equation*}
$$

Suppose that $\pi_{[0, t]}$ is a random path of the same type, with the following properties:
(p1) $\xi_{0}\left(\pi_{0}\right)=1$ a.s. and $\left(\pi_{s}, s\right) \leftrightarrow\left(\pi_{u}, u\right)$ for all $s, u \in[0, t]$.
(p2) $\pi$ is $\mathcal{F}$-adapted and $\left\{\pi_{s} \geq \varpi_{s} \forall s \in[0, t]\right\} \in \overline{\mathcal{R}}_{t}^{\varpi}$ for all deterministic paths $\varpi$.
We call $\pi$ a random infection path (see Fig. 5), a name that is justified by (p1). Property (p2) means that $\pi$ is causal and that, when we discover it, we leave the graphical representation to its left untouched. For such $\pi$, let

$$
\begin{equation*}
\mathcal{R}_{t}^{\pi}:=\sigma\left(\pi,\left(\xi_{0}(x)\right)_{x \geq \pi_{0}},\left(H_{s}(x), I_{s}(x)\right)_{s \in[0, t], x \geq \pi_{s}}\right) . \tag{4.3}
\end{equation*}
$$

Note that, since $\pi$ is an infection path, also $\left(\xi_{s}(x)\right)_{x \geq \pi_{s}} \in \mathcal{R}_{t}^{\pi}$ for each $s \in[0, t]$ (see the proof of Lemma 2.2). We have the following stochastic domination result.

Lemma 4.1. For any random infection path $\pi_{[0, t]}$ as above, the law of $\xi_{t}\left(\cdot+\pi_{t}+1\right)$ under $\mathbb{P}_{\bar{\nu}_{\lambda}}\left(\cdot \mid \mathcal{R}_{t}^{\pi}\right)$ is stochastically larger than $\bar{\nu}_{\lambda}$.

Proof. Construct $\mathbb{P}_{\bar{\nu}_{\lambda}}$ from a graphical representation on $\mathbb{Z} \times \mathbb{R}$ as outlined above by adding healing events on $(x, 0)$ for each $x \in \mathbb{Z}_{\geq 0}$. Extend $\pi$ to negative times by making it equal to the right-most infinite infection path going backwards in time from $\left(\pi_{0}, 0\right)$. (Such a path exists because $\xi_{0}\left(\pi_{0}\right)=1$.) We may check that the resulting path still has properties (p1) and (p2). Extend also $\mathcal{R}_{t}^{\pi}$ to include negative times.

Next, regard $H$ and $I$ as Poisson point processes on subsets of $\mathbb{Z} \times \mathbb{R}$. Let (see Fig. 5)

$$
\begin{equation*}
D:=\left\{(x, s) \in \mathbb{Z} \times \mathbb{R}: s>t \text { or } \pi_{s}>x\right\} \tag{4.4}
\end{equation*}
$$



Figure 5: The thick line represents the random infection path $\pi$. The dashed lines represent other infection paths.

Given $\mathcal{R}_{t}^{\pi}$, by (p2) $H$ and $I$ are still Poisson point processes with the same densities on $D$. This can be justified first for $\pi$ taking values in a countable set and then for general $\pi$ using right-continuity.

With this observation we can couple $\mathbb{P}_{\nu_{\lambda}}$ to $\mathbb{P}_{\bar{\nu}_{\lambda}}\left(\cdot \mid \mathcal{R}_{t}^{\pi}\right)$ in the following way. Draw independent Poisson point processes $\hat{H}, \hat{I}$ on $D^{c}$. Take $\hat{\xi}$ to be the contact process obtained by using $H, I$ on $D$ and $\hat{H}, \hat{I}$ on $D^{c}$. Then $\hat{\xi}$ is distributed as the contact process under $\mathbb{P}_{\nu_{\lambda}}$, and is independent of $\mathcal{R}_{t}^{\pi}$. Furthermore, $\xi_{t}(x) \geq \hat{\xi}_{t}(x)$ for all $x<\pi_{t}$. Indeed, if $\xi_{t}(x)=1$, then infinite infection paths going backwards in time must either stay inside $D$ or cross $\pi$, so that, by ( p 1$), \xi_{t}(x)=1$ as well.

Remark 4.2. In Lemma 4.1, we may replace $t$ by a finite stopping time $\mathcal{T}$ w.r.t. the filtration $\mathcal{F}$, as long as the event in (p2) is replaced by $\left\{\mathcal{T} \leq t, \pi_{s} \geq \varpi_{s} \forall s \in[0, \mathcal{T}]\right\}$ and we add $\mathcal{T}$ to $\mathcal{R}_{\mathcal{T}}^{\pi}$. We may also enlarge all filtrations by adding information that is independent of $\xi_{0}, H, I$, in particular, $N_{[0, t]}$ and $U_{\left[1, N_{t}\right]}$ (recall Section 2.2).

Infection range. Lemma 4.3 below concerns the positions of wide-spread infections. For $\delta \in(0, \iota)$ and $x \in \mathbb{Z}$, let $\mathcal{W}_{x}^{\delta}:=\{(z, t) \in \mathbb{Z} \times[0, \infty):(\iota-\delta) t-1<z-x \leq(\iota+\delta) t\}$ be a wedge between two lines of inclination $\iota-\delta$ and $\iota+\delta$. Set $C_{t}^{\delta}(x):=\{y \in \mathbb{Z}:(y, t) \leftrightarrow$ $(x, 0)$ via a path contained in $\left.\mathcal{W}_{x}^{\delta}\right\}$, and

$$
\begin{equation*}
Z_{\delta}(x):=\sup \left\{z \in \mathbb{Z}_{<x}: \xi_{0}(z)=1, C_{t}^{\delta}(z) \neq \emptyset \forall t \geq 0\right\} \tag{4.5}
\end{equation*}
$$

i.e., the first infected site to the left of $x$ that spreads its infection forever inside a wedge.

Lemma 4.3. If $\lambda \in\left(\lambda_{c}, \infty\right)$ then $\left|Z_{\delta}(x)-x\right|$ has exponential moments under $\mathbb{P}_{\bar{\nu}_{\lambda}}$ for every $\delta \in(0, \iota)$, uniformly in $x \in \mathbb{Z}_{\leq 0}$.

Proof. We will use the fact that, for any $\lambda \in\left(\lambda_{c}, \infty\right), \nu_{\lambda}$ stochastically dominates a non-trivial Bernoulli product measure $\mu_{\lambda}$. This follows from Liggett and Steif [15], Theorem 1.2, Durrett and Schonmann [9], Theorem 1, and van den Berg, Häggström and Kahn [7], Theorem 3.5. Since $Z_{\delta}(x)$ is monotone in $\xi_{0}$, it is therefore enough to prove the statement under $\mathbb{P}_{\mu_{\lambda}}$. We may also assume $x=0$, as $Z_{\delta}(x)$ does not depend on $\left(\xi_{0}(z)\right)_{z \geq x}$.

Construct a sequence of pairs $\left(Z_{n}, T_{n}\right)_{n \in \mathbb{N}_{0}}$ as follows. Set $Z_{0}=T_{0}:=0$ and, recursively for $n \in \mathbb{N}_{0}$,

Conditionally on $T_{n}<\infty, \Delta_{n+1}:=Z_{n+1}-Z_{n}+\left\lceil(\iota+\delta) T_{n}\right\rceil$ and $T_{n+1}$ are independent of $\left(Z_{k}, T_{k}\right)_{k=1}^{n}$ and distributed as $\left(Z_{1}, T_{1}\right)$. This is because the region of the graphical representation plus initial configuration on which $T_{n+1}$ and $\Delta_{n+1}$ depend is disjoint from the region on which the previous random variables depend. Since $\mu_{\lambda}$ is a non-trivial product measure, $\left|Z_{1}\right|$ has exponential moments. Noting that $T_{1}$ is independent of $Z_{1}$ we conclude, using standard facts about the contact process (see Liggett [12], Chapter VI, Theorem 2.2, Corollary 3.22 and Theorem 3.23), that $\mathbb{P}_{\mu_{\lambda}}\left(T_{1}=\infty\right)>0$ and that, conditionally on $T_{1}<\infty, T_{1}$ has exponential moments. Defining the random index

$$
\begin{equation*}
K:=\inf \left\{n \in \mathbb{N}: T_{n}=\infty\right\} \tag{4.7}
\end{equation*}
$$

whose distribution is $\operatorname{GEO}\left(\mathbb{P}_{\mu_{\lambda}}\left(T_{1}=\infty\right)\right.$ ), we see that $\left|Z_{\delta}(0)\right| \leq\left|Z_{K}\right|$. Taking $a>$ 0 such that $\mathbb{E}_{\mu_{\lambda}}\left[e^{a\left(\left|Z_{1}\right|+\left\lceil(\iota+\delta) T_{1}\right\rceil\right)} \mid T_{1}<\infty\right]<1 / \mathbb{P}_{\mu}\left(T_{1}<\infty\right)$, we get after a short calculation that $\mathbb{E}_{\mu_{\lambda}}\left[\mathbb{1}_{\{K=n\}} e^{a\left|Z_{n}\right|}\right]$ decays exponentially in $n$.

## 5 Properties of the speed

In this section we prove Theorem 1.1(b).
For each $n \in \mathbb{N}, W_{n}$ depends on $\xi$ in a finite space-time region. Therefore $\lambda \mapsto$ $\mathbb{E}_{\delta_{1}}\left[n^{-1} W_{n}\right]$ is continuous (see Liggett [14], Part I). Since, by monotonicity, the latter is non-decreasing, it follows from (3.7) that $\lambda \mapsto v(\lambda)$ is right-continuous and nondecreasing.

It remains to show that $v(\lambda) \in\left(v_{0}, v_{1}\right)$ and $\lim _{\lambda \rightarrow \infty} v(\lambda)=v_{1}$. This will be done in Sections 5.1-5.2 below. These properties come from the fact that the random walk spends positive fractions of its time on top of infected sites and on top of healthy sites. To keep track of this, define $N_{t}^{i}:=\#\left\{n \in \mathbb{N}: \xi_{J_{n}}\left(W_{J_{n-1}}\right)=i\right\}, i \in\{0,1\}$. Recalling the construction of $W$ in Section 2.2, we may write

$$
\begin{equation*}
W_{t}=S_{N_{t}^{0}}^{0}+S_{N_{t}^{1}}^{1}, \tag{5.1}
\end{equation*}
$$

where $S_{n}^{i}, i=0,1$, are discrete-time homogeneous random walks that jump to the right with probability $\alpha_{i} / \gamma$ and to the left with probability $\beta_{i} / \gamma$. From this representation we immediately get the following.

Lemma 5.1.

$$
\begin{align*}
& \liminf _{t \rightarrow \infty} t^{-1} W_{t}=v_{0}+\left(v_{1}-v_{0}\right) \liminf _{t \rightarrow \infty}(\gamma t)^{-1} N_{t}^{1} \\
& \limsup _{t \rightarrow \infty} t^{-1} W_{t}=v_{1}-\left(v_{1}-v_{0}\right) \liminf _{t \rightarrow \infty}(\gamma t)^{-1} N_{t}^{0} \tag{5.2}
\end{align*}
$$

Lemma 5.1 is valid for any dynamic random environment, even without a SLLN for $W$. But (5.2) shows that a SLLN for $W$ holds with speed $v$ if and only if a SLLN holds for $N^{1}$ with limit $\gamma \rho_{\text {eff }}$, where $\rho_{\text {eff }}:=\left(v-v_{0}\right) /\left(v_{1}-v_{0}\right)$ is the effective density of 1 's seen by $W$. Thus, $v>v_{0}$ and $v<v_{1}$ are equivalent to, respectively, $\rho_{\text {eff }}>0$ and $\rho_{\text {eff }}<1$.

### 5.1 Proof of $v(\lambda)<v_{1}$

In the contact process, infected sites heal spontaneously. Therefore it is easier to find 0 's than 1's. For this reason, it is easier to prove that $W$ often jumps from healthy sites than from infected sites.

Proof. For $k \in \mathbb{N}$, let $Y_{k}:=\xi_{J_{k}}\left(W_{J_{k-1}}\right)$, and note that $\left\{Y_{k+1}=0\right\}$ contains all configurations that between times $J_{k}$ and $J_{k+1}$ have a cross at site $W_{J_{k}}$ and no arrows between $W_{J_{k}}$ and its nearest-neighbors, i.e., such that the events $H_{J_{k+1}}\left(W_{J_{k}}\right)-H_{J_{k}}\left(W_{J_{k}}\right) \geq 1$ and $I_{J_{k+1}}\left(W_{J_{k}}\right)-I_{J_{k}}\left(W_{J_{k}}\right)=I_{J_{k+1}}\left(W_{J_{k}}-1\right)-I_{J_{k}}\left(W_{J_{k}}-1\right)=0$ occur. The probability of the latter events given $\sigma\left\{\left(J_{k}, \xi_{s}, W_{s}\right)_{0 \leq s \leq J_{k}}\right\}$ is constant in $k$ and equal to $p:=\gamma /(\gamma+2 \lambda)(1+\gamma+2 \lambda)$. Therefore the sequence $\left(Y_{k}\right)_{k \in \mathbb{N}}$ is stochastically dominated by a sequence of i.i.d. BERN $(1-p)$ random variables, which implies that $\liminf _{t \rightarrow \infty} t^{-1} N_{t}^{0} \geq \gamma p>0$, so that $v(\lambda)<v_{1}$ by Lemma 5.1.

### 5.2 Proof of $v(\lambda)>v_{0}$ and $\lim _{\lambda \rightarrow \infty} v(\lambda)=v_{1}$

This is the harder part of the proof. We will need results from Section 4. In the following we will assume that $v_{0} \leq 0$. The case $v_{0}>0$ can be treated analogously.

Let us start with an informal description of the argument. The idea is that there are "waves of infection" coming from $\pm \infty$ from which the random walk cannot escape. When $v_{0} \leq 0$, we can concentrate on the waves coming from the left, represented schematically in Fig. 6. Each time the random walk hits a new wave, there is an infection path starting from its current location and going backwards in time entirely to the left of the random walk path. By Lemma 4.1, at this time the law of $\xi$ to the left of the random walk has an appreciable density, which means that there are new waves coming in from locations not very far to the left. On the other hand, any infections to the right of the random walk can be ignored, since they only push it to the right. But doing so makes the random walk behave as a homogeneous random walk with a non-positive drift, meaning that it does not take the random walk long to hit the next infection wave. Since at each collision there is a fixed probability for the random walk
to jump while sitting on an infection, $v(\lambda)>v_{0}$ will follow from Lemma 5.1. With some care in the computations we also get the limit for large $\lambda$.


Figure 6: The dashed lines represent infection waves. The thick line represents the path of $W$.

Proof. Using the graphical representation, we will construct, on a larger probability space, a second random walk $\hat{W}$ coupled to $W$ in such a way that $\hat{W}_{t} \leq W_{t}$ for all $t \geq 0$ and that $\hat{W}$ has a speed with the desired properties. Let

$$
\begin{equation*}
V_{1}:=\inf \left\{t>0: \xi_{t}\left(W_{t}\right)=1\right\} . \tag{5.3}
\end{equation*}
$$

Note that $V_{1}$ has exponential moments under $\mathbb{P}_{\bar{\nu}_{\lambda}}$ by Lemma 4.3 and the fact that $v_{0} \leq 0$. Let

$$
\begin{equation*}
\tau_{1}:=\inf \left\{t>V_{1}: W_{t} \neq W_{V_{1}} \text { or } H_{t}\left(W_{V_{1}}\right)>H_{V_{1}}\left(W_{V_{1}}\right)\right\}, \tag{5.4}
\end{equation*}
$$

i.e., $\tau_{1}$ is the first time after time $V_{1}$ at which either $W$ jumps or there is a healing event at the position of the random walk. Note that $\tau_{1}$ is a stopping time w.r.t. the filtration $\mathcal{G}$ and that, given $\mathcal{G}_{V_{1}}, \tau_{1}-V_{1}$ has distribution $\operatorname{EXP}(1+\gamma)$.

We will construct a sequence $\left(W^{(n)}, \tau_{n}\right)_{n \in \mathbb{N}}$ with the following properties:
(A1) $W_{t}^{(n+1)} \leq W_{\tau_{n}+t}^{(n)}-W_{\tau_{n}}^{(n)}$ for all $t \geq 0$;
(A2) $\left(W^{(n)}, \tau_{n}\right)$ is distributed as $\left(W, \tau_{1}\right)$ under $\mathbb{P}_{\bar{\nu}_{\lambda}}$;
(A3) $\left(W_{\left[0, \tau_{n}\right]}^{(n)}, \tau_{n}\right)_{n \in \mathbb{N}}$ is i.i.d.;
(A4) If $\hat{v}(\lambda):=\mathbb{E}_{\bar{\nu}_{\lambda}}\left[W_{\tau_{1}}\right] / \mathbb{E}_{\bar{\nu}_{\lambda}}\left[\tau_{1}\right]$, then $\hat{v}(\lambda)>v_{0}$ and $\lim _{\lambda \rightarrow \infty} \hat{v}(\lambda)=v_{1}$.
Once we have this sequence, we can put $T_{0}:=0, T_{n}:=\sum_{k=1}^{n} \tau_{k}$ for $n \in \mathbb{N}$, and

$$
\begin{equation*}
\hat{W}_{t}:=\sum_{k=1}^{n} W_{\tau_{k}}^{(k)}+W_{t-T_{n}}^{(n+1)} \quad \text { for } \quad T_{n} \leq t<T_{n+1} . \tag{5.5}
\end{equation*}
$$

By (A1), $\hat{W}_{t} \leq W_{t}^{(1)}$ for all $t \geq 0$. By (A2), the latter is distributed as $W$ under $\mathbb{P}_{\bar{\nu}_{\lambda}}$, which by monotonicity is stochastically smaller than $W$ under $\mathbb{P}_{\nu_{\lambda}}$. By (A3), $\lim _{n \rightarrow \infty} T_{n}^{-1} \hat{W}_{T_{n}}=\hat{v}(\lambda)$, and so the claim follows from (A4). Thus, it remains to construct the sequence $\left(W^{(n)}, \tau_{n}\right)_{n \in \mathbb{N}}$ with properties (A1)-(A4).

To do so, we draw $\xi_{0}$ from $\bar{\nu}_{\lambda}$, let $\xi^{(1)}:=\xi, W^{(1)}:=W$, define $\tau_{1}$ as above, and note the following.

Lemma 5.2. Under $\mathbb{P}_{\bar{\nu}_{\lambda}}\left(\cdot \mid \tau_{1}, W_{\left[0, \tau_{1}\right]}\right)$, the law of $\xi_{\tau_{1}}\left(\cdot+W_{\tau_{1}}\right)$ is stochastically larger than $\bar{\nu}_{\lambda}$.

Proof. Since $\xi_{V_{1}}\left(W_{V_{1}}\right)=1$, there exists a right-most path $\pi_{\left[0, V_{1}\right]}$ connecting $\left(W_{V_{1}}, V_{1}\right)$ to $\mathbb{Z}_{\leq-1} \times\{0\}$. Extend $\pi$ to $\left[V_{1}, \tau_{1}\right]$ by making it constant and equal to $W_{V_{1}}$ on this time interval. Since $\pi_{s} \leq W_{s}$ for all $0 \leq s<\tau_{1}$, we have $\left(\tau_{1}, W_{\left[0, \tau_{1}\right]}\right) \in \mathcal{R}_{\tau_{1}}^{\pi} \vee \sigma\left(N_{\left[0, \tau_{1}\right]}, U_{\left[1, N_{\tau_{1}}\right]}\right)$. Note that $\pi$ is not an infection path, but only because of a possible healing event at time $\tau_{1}$, which does not affect $\left(\xi_{\tau_{1}}\left(x+W_{V_{1}}\right)\right)_{x \leq-1}$. Therefore, by Lemma 4.1, the distribution of the latter given $\left(\tau_{1}, W_{\left[0, \tau_{1}\right]}\right)$ is stochastically larger than $\bar{\nu}_{\lambda}$. Using this observation and noting that $W_{\tau_{1}} \neq W_{V_{1}}$ if and only if $\xi_{\tau_{1}}\left(W_{V_{1}}\right)=1$, we can verify that the claim holds for each possible outcome of $W_{\tau_{1}}-W_{V_{1}} \in\{0, \pm 1\}$.

By Lemma 5.2, there exists a configuration $\xi_{0}^{(2)}$ distributed as $\bar{\nu}_{\lambda}$, independent of $\left(\tau_{1}, W_{\left[0, \tau_{1}\right]}\right)$ and stochastically smaller than $\xi_{\tau_{1}}^{(1)}\left(\cdot+W_{\tau_{1}}\right)$. We may now define $\xi^{(2)}$ by using the events of the graphical representation that lie above time $\tau_{1}$ with the origin shifted to $W_{\tau_{1}}$, using $\xi_{0}^{(2)}$ as starting configuration. We may then define $W^{(2)}$ and $\tau_{2}$ from $\xi^{(2)},\left(N_{t+\tau_{1}}-N_{\tau_{1}}\right)_{t \geq 0}$ and $\left(U_{k}\right)_{k>N_{\tau_{1}}}$. With this coupling, clearly $W_{t}^{(2)} \leq W_{\tau_{1}+t}^{(1)}-W_{\tau_{1}}^{(1)}$ for all $t \geq 0$. Furthermore, since $\xi_{0}^{(2)}$ is independent of $\left(\tau_{1}, W_{\left[0, \tau_{1}\right]}\right)$, the distribution of $\xi_{\tau_{2}}^{(2)}\left(\cdot+W_{\tau_{2}}^{(2)}\right)$ given $\left(W_{\left[0, \tau_{i}\right]}^{(i)}, \tau_{i}\right)_{i=1,2}$ depends only on the random variables with $i=2$ and hence, by Lemma 5.2 , is again stochastically larger than $\bar{\nu}_{\lambda}$.

We may therefore repeat the argument. More precisely, suppose by induction that we have defined $\xi^{(k)}$, $W^{(k)}$ and $\tau_{k}$ for $k=1, \ldots, n$ and $n \geq 2$, in such a way that:
(B1) $W_{t}^{(k+1)} \leq W_{\tau_{n}+t}^{(k)}-W_{\tau_{n}}^{(n)}$ for all $t \geq 0$ and $k=1 \ldots n-1$;
(B2) $\left(W^{(k)}, \tau_{k}\right)$ is distributed as $\left(W, \tau_{1}\right)$ under $\mathbb{P}_{\bar{\nu}_{\lambda}}$ for all $k=1, \ldots, n$;
(B3) $\left(W_{\left[0, \tau_{k}\right]}^{(k)}, \tau_{k}\right)_{k=1}^{n}$ is i.i.d.;
(B4) The law of $\xi^{(n)}\left(\cdot+W_{\tau_{n}}^{(n)}\right)$ given $\left(W_{\left[0, \tau_{k}\right]}^{(k)}, \tau_{k}\right)_{k=1}^{n}$ is stochastically larger than $\bar{\nu}_{\lambda}$.
Then we proceed as before: there exists a configuration $\xi_{0}^{(n+1)}$ distributed as $\bar{\nu}_{\lambda}$, stochastically smaller than $\xi^{(n)}\left(\cdot+W_{\tau_{n}}^{(n)}\right)$ and independent of $\left(W_{\left[0, \tau_{k}\right]}^{(k)}, \tau_{k}\right)_{k=1}^{n}$, from which we obtain $\xi^{(n+1)}, W^{(n+1)}$ and $\tau_{n+1}$, and we prove (B1)-(B4) like in the case $n=2$. This settles the existence of the sequence $\left(W^{(n)}, \tau_{n}\right)_{n \in \mathbb{N}}$. All that is left to show is that $\hat{v}(\lambda)>v_{0}$ and $\lim _{\lambda \rightarrow \infty} \hat{v}(\lambda)=v_{1}$.

Note that Lemma 5.1 is valid also for $\hat{W}$, and write $\hat{N}_{t}^{1}$ to denote the number of jumps that $\hat{W}$ takes on infected sites. Then $\hat{N}_{T_{n}}^{1}$ has distribution $\operatorname{BINOM}(n, \gamma /(1+\gamma))$, and by standard arguments we obtain

$$
\begin{equation*}
\lim _{t \rightarrow \infty} t^{-1} \hat{N}_{t}^{1}=\frac{\gamma}{(1+\gamma) \mathbb{E}_{\bar{\nu}_{\lambda}}\left[\tau_{1}\right]}>0 \tag{5.6}
\end{equation*}
$$

which proves $\hat{v}(\lambda)>v_{0}$. Furthermore, we claim that $\lim _{\lambda \rightarrow \infty} \mathbb{E}_{\bar{\nu}_{\lambda}}\left[V_{1}\right]=0$. Indeed, $V_{1}$ is nonincreasing in $\lambda$ and, since $\lim _{\lambda \rightarrow \infty} \rho_{\lambda}=1$ (recall Section 1.2), it is not hard to see that $V_{1}$ converges in probability to zero as $\lambda \rightarrow \infty$. Therefore $\lim _{\lambda \rightarrow \infty} \mathbb{E}_{\bar{\nu}_{\lambda}}\left[\tau_{1}\right]=1 /(1+\gamma)$, and so $\lim _{\lambda \rightarrow \infty} \hat{v}(\lambda)=v_{1}$.

## 6 FCLT and LDP

The proof of Theorem 1.2 depends on the construction of regeneration times, i.e., times at which the random walk forgets its past. This construction will be carried out in Section 6.1 and is based on two propositions (Propositions 6.1-6.2 below), which are proved in Sections 6.2-6.3. At the end of Section 6.1 we will see that these propositions imply Theorem $1.2(\mathrm{a}, \mathrm{c})$. The proof of Theorem 1.2(b) is deferred to Section 6.4.

### 6.1 Regeneration times

If the infection propagation speed $\iota=\iota(\lambda)$ is larger than $\left|v_{0}\right| \vee\left|v_{1}\right|$, the maximum absolute speed at which the random walk can move, then each time $W$ finds itself on an infected site it can become "trapped" forever in an infection cluster generated by this site alone. In that case, by Lemma 2.2, the future increments of $W$ become independent of its past. The issue is therefore to find enough moments when $W$ sits on an infection. This can be dealt with in a way similar to what was done in the proof of $v(\lambda)>v_{0}$ in Section 5.2.

Hitting, failure and trial times. In order to build the regeneration structure, we first need to extend some definitions related to clusters and right-most infections. For $s \geq t$ and $x \in \mathbb{Z}$, let

$$
\begin{equation*}
C_{t, s}(x):=\{y \in \mathbb{Z}:(x, t) \leftrightarrow(y, s)\} \tag{6.1}
\end{equation*}
$$

and

$$
\begin{equation*}
R_{t, s}(x):=\sup C_{t, s}(x), \quad L_{t, s}(x):=\inf C_{t, s}(x) \tag{6.2}
\end{equation*}
$$

Furthermore, let

$$
\begin{equation*}
r_{t, s}(x):=\sup _{\substack{y<x \\ \xi_{t}(y)=1}} R_{t, s}(y), \tag{6.3}
\end{equation*}
$$

i.e., the right-most infection at time $s$ that comes from $\mathbb{Z}_{\leq x-1} \times\{t\}$.

For $t \geq 0$ and $z \in \mathbb{Z}$, let

$$
\begin{equation*}
V_{t}(z):=\inf \left\{s>t: W_{s}=r_{t, s}(z)\right\} \tag{6.4}
\end{equation*}
$$

be the first time after time $t$ at which $W$ meets the right-most infection coming from $\mathbb{Z}_{\leq z-1}$. We will call this the $z$-wave hitting time after $t$. It is not hard to see that $V_{t}(z)<\infty \mathbb{P}_{\nu_{\lambda}}$-a.s. for any $t$ and $z \leq W_{t}$. Indeed, at any time $t$ there is an infected site $x<z$ whose infection survives forever, and in this case $\lim _{s \rightarrow \infty} s^{-1} R_{t, s}(x)=\iota>$ $\left|v_{0}\right| \vee\left|v_{1}\right|$. Therefore there must be an $s>t$ for which $R_{t, s}(x)=W_{s}$. By right-continuity, $\mathbb{P}_{\nu_{\lambda}}\left(V_{t}(z)<\infty \forall z \leq W_{t}, t \geq 0\right)=1$ as well.

Now define the first failure time after time $t$ by (see Fig. 8)

$$
\begin{equation*}
F_{t}:=\inf \left\{s>t: W_{s} \notin\left[L_{t, s}\left(W_{t}\right), R_{t, s}\left(W_{t}\right)\right]\right\} \tag{6.5}
\end{equation*}
$$

i.e., the first time after time $t$ when $W$ exits the region surrounded by the cluster of $\left(W_{t}, t\right)$. To keep track of the space-time region on which the failure time depends, define, for $t \geq 0$ and $x \in \mathbb{Z}$,

$$
\begin{equation*}
\left(Y_{t, s}(x)\right)_{s \geq t} \tag{6.6}
\end{equation*}
$$

as the process with values in $\mathbb{Z}$ that starts at time $t$ at site $x$ and jumps down by following the infection arrows to the left in the graphical representation (see Fig. 7). Then, given $\mathcal{G}_{t},\left(x-Y_{t, t+s}(x)\right)_{s \geq 0}$ is a Poisson process with rate $\lambda$.


Figure 7: $Y_{t, s}(x)$ starts at $x$ and goes upwards and to the left across the arrows of the graphical representation.

With the above observations we can define the trial time after a failure time (see Fig. 8):

$$
T_{t}:= \begin{cases}\infty & \text { if } F_{t}=\infty  \tag{6.7}\\ V_{F_{t}}\left(Y_{t, F_{t}}\left(W_{t}\right)\right) & \text { otherwise }\end{cases}
$$

i.e., $T_{t}$ is the $Y_{t, F_{t}}\left(W_{t}\right)$-wave time after time $F_{t}$ when the latter is finite. This wave ensures "good conditions" at the trial time, meaning an appreciable density of infections to the left of $W$.


Figure 8: A failure time $F_{t}$ and a trial time $T_{t}$ after time $t$. The dashed lines represent infection paths. The thick line represents the path of $W$.

Regeneration times. We can now define our regeneration time $\tau$. First let

$$
\begin{equation*}
\mathcal{T}_{1}:=V_{0}(0) \tag{6.8}
\end{equation*}
$$

and, under the assumption that $\mathcal{T}_{1}, \ldots, \mathcal{T}_{k}, k \in \mathbb{N}$, are all defined, let

$$
\mathcal{T}_{k+1}:= \begin{cases}\infty & \text { if } \mathcal{T}_{k}=\infty  \tag{6.9}\\ T_{\mathcal{T}_{k}} & \text { otherwise }\end{cases}
$$

Note that the $\mathcal{T}_{k}$ 's are stopping times w.r.t. the filtration $\mathcal{G}$. Finally, put

$$
\begin{equation*}
K:=\inf \left\{k \in \mathbb{N}: \mathcal{T}_{k}<\infty, \mathcal{T}_{k+1}=\infty\right\} \tag{6.10}
\end{equation*}
$$

and let

$$
\begin{equation*}
\tau:=\mathcal{T}_{K} \tag{6.11}
\end{equation*}
$$

Note that $K<\infty$ a.s. since, at any trial time, the probability for the next failure time to be infinite is uniformly bounded from below. We will prove in Sections 6.2-6.3 that $\tau$ is a regeneration time and has exponential moments. This is stated in the following two propositions.

Proposition 6.1. The distribution of $\left(W_{t+\tau}-W_{\tau}\right)_{t \geq 0}$ under both $\mathbb{P}_{\nu_{\lambda}}\left(\cdot \mid \tau, W_{[0, \tau]}\right)$ and $\mathbb{P}_{\nu_{\lambda}}\left(\cdot \mid \Gamma, \tau, W_{[0, \tau]}\right)$ is the same as that of $W$ under $\mathbb{P}_{\nu_{\lambda}}(\cdot \mid \Gamma)$, where

$$
\begin{equation*}
\Gamma:=\left\{\xi_{0}(0)=1, F_{0}=\infty\right\} . \tag{6.12}
\end{equation*}
$$

Proposition 6.2. $\tau$ and $\left|W_{\tau}\right|$ have exponential moments under both $\mathbb{P}_{\nu_{\lambda}}$ and $\mathbb{P}_{\nu_{\lambda}}(\cdot \mid \Gamma)$, uniformly in $\lambda \in\left[\lambda_{-}, \lambda_{+}\right]$for any fixed $\lambda_{-}, \lambda_{+} \in\left(\lambda_{W}, \infty\right)$.

These two propositions imply the LLN and Theorem 1.2(a), with

$$
\begin{equation*}
v(\lambda)=\frac{\mathbb{E}_{\nu_{\lambda}}\left[W_{\tau} \mid \Gamma\right]}{\mathbb{E}_{\nu_{\lambda}}[\tau \mid \Gamma]}, \quad \sigma(\lambda)^{2}=\frac{\mathbb{E}_{\nu_{\lambda}}\left[\left(W_{\tau}\right)^{2} \mid \Gamma\right]-\mathbb{E}_{\nu_{\lambda}}\left[W_{\tau} \mid \Gamma\right]^{2}}{\mathbb{E}_{\nu_{\lambda}}[\tau \mid \Gamma]} \tag{6.13}
\end{equation*}
$$

They also imply that

$$
\begin{equation*}
\limsup _{t \rightarrow \infty} \frac{1}{t} \log \mathbb{P}_{\nu_{\lambda}}\left(t^{-1} W_{t} \notin(v-\epsilon, v+\epsilon)\right)<0 \quad \forall \epsilon>0 \tag{6.14}
\end{equation*}
$$

For a proof of these facts, the reader can follow word-by-word the arguments given in Avena, dos Santos and Völlering [5], Theorem 3.8 and Section 4.1 (which do not require (1.5)-(1.6)).

Theorem 1.2(c) follows from (6.14) and the partial LDP proven in Avena, den Hollander and Redig [3] for attractive spin-flip systems (including the contact process). Here, partial means that the LDP is shown to hold outside a possible interval where the rate function is zero. However, (6.14) precisely precludes the presence of such an interval. (See Glynn and Whitt [10], Theorem 3, for more details.)

The proof of Theorem 1.2(b) is deferred to Section 6.4.

### 6.2 Proof of Proposition 6.1

We first show that the regeneration strategy indeed makes sense.
Lemma 6.3. For all $t \geq 0$,

$$
\begin{equation*}
\mathbb{P}_{\nu_{\lambda}}\left(F_{t}=\infty,\left(W_{s+t}-W_{t}\right)_{s \geq 0} \in \cdot \mid \mathcal{G}_{t}\right)=\mathbb{P}_{\mathbb{1}_{0}}\left(\Gamma_{0}, W \in \cdot\right) \text { a.s. on }\left\{\xi_{t}\left(W_{t}\right)=1\right\} \text {, } \tag{6.15}
\end{equation*}
$$

where $\Gamma_{0}:=\left\{F_{0}=\infty\right\}$. The same is true for a finite stopping time w.r.t. $\mathcal{G}$ replacing $t$.

Proof. First note that $\mathbb{P}_{\eta}\left(\Gamma_{0}, W \in \cdot\right)=\mathbb{P}_{\mathbb{1}_{0}}\left(\Gamma_{0}, W \in \cdot\right)$ for any $\eta$ with $\eta(0)=1$. This follows from Lemma 2.2 because, on $\Gamma_{0}, W$ depends on $\xi$ only through $\left\{\xi_{t}(x): t \geq\right.$ $\left.0, x \in\left[L_{t}(0), R_{t}(0)\right]\right\}$, and $\Gamma_{0}$ does not depend on $\xi_{0}$. Now, letting $\hat{\xi}_{t}(\cdot):=\xi_{t}\left(\cdot+W_{t}\right)$, we can write (recall (6.5))

$$
\begin{align*}
\mathbb{P}_{\nu_{\lambda}}\left(\xi_{t}\left(W_{t}\right)=1, F_{t}\right. & \left.=\infty,\left(W_{s+t}-W_{t}\right)_{s \geq 0} \in \cdot \mid \mathcal{G}_{t}\right) \\
= & \mathbb{E}_{\nu_{\lambda}}\left[\xi_{t}\left(W_{t}\right) \mathbb{P}_{\hat{\xi}_{t}}\left(\Gamma_{0}, W \in \cdot\right) \mid \mathcal{G}_{t}\right]=\xi_{t}\left(W_{t}\right) \mathbb{P}_{\mathbb{1}_{0}}\left(\Gamma_{0}, W \in \cdot\right), \tag{6.16}
\end{align*}
$$

where the first equality is justified by the Markov property and the translation invariance of the graphical representation. To extend the result to stopping times we can use the strong Markov property of $(\xi, W)$.

With the help of Lemma 6.3 we are ready to prove Proposition 6.1.
Proof. We will closely follow the proof of Theorem 3.4 in [5]. Let $\mathcal{G}_{\tau}$ be the $\sigma$-algebra of all events $B$ such that, for all $n \in \mathbb{N}_{0}$, there exists a $B_{n} \in \mathcal{G}_{\mathcal{T}_{n}}$ such that $B \cap\{K=$ $n\}=B_{n} \cap\{K=n\}$. Note that $\tau$ and $W_{[0, \tau]}$ are in $\mathcal{G}_{\tau}$.

In the following, we abreviate $W^{(t)}:=\left(W_{s+t}-W_{t}\right)_{s \geq 0}$. Pick $f$ bounded and measurable, $B \in \mathcal{G}_{\tau}$, and write (recall (6.9))

$$
\begin{align*}
& \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{B} f\left(W^{(\tau)}\right)\right]=\sum_{n \in \mathbb{N}_{0}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{B_{n}} \mathbb{1}_{\{K=n\}} f\left(W^{\left(\mathcal{T}_{n}\right)}\right)\right] \\
& =\sum_{n \in \mathbb{N}_{0}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{B_{n}} \mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\left\{F_{\mathcal{F}_{n}}=\infty\right\}} f\left(W^{\left(\mathcal{T}_{n}\right)}\right) \mid \mathcal{G}_{\mathcal{T}_{n}}\right]\right] . \tag{6.17}
\end{align*}
$$

Since $\xi_{\mathcal{T}_{n}}\left(W_{\mathcal{T}_{n}}\right)=1$ on $\left\{\mathcal{T}_{n}<\infty\right\}$, by Lemma 6.3 the last line of (6.17) equals

$$
\begin{align*}
& \mathbb{E}_{\mathbb{1}_{0}}\left[f(W) \mathbb{1}_{\Gamma_{0}}\right] \sum_{n \in \mathbb{N}_{0}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{B_{n}} \mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}}\right]  \tag{6.18}\\
& =\mathbb{E}_{\mathbb{1}_{0}}\left[f(W) \mid \Gamma_{0}\right] \sum_{n \in \mathbb{N}_{0}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{B_{n}} \mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}}\right] \mathbb{P}_{\mathbb{1}_{0}}\left(\Gamma_{0}\right),
\end{align*}
$$

which, again by Lemma 6.3, equals

$$
\begin{align*}
& \mathbb{E}_{\mathbb{1}_{0}}\left[f(W) \mid \Gamma_{0}\right] \sum_{n \in \mathbb{N}_{0}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{B_{n}} \mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}} \mathbb{P}_{\nu_{\lambda}}\left(F_{\mathcal{T}_{n}}=\infty \mid \mathcal{G}_{\mathcal{T}_{n}}\right)\right] \\
& =\mathbb{E}_{\mathbb{1}_{0}}\left[f(W) \mid \Gamma_{0}\right] \sum_{n \in \mathbb{N}_{0}} \mathbb{P}_{\nu_{\lambda}}\left(B_{n}, K=n\right)  \tag{6.19}\\
& =\mathbb{E}_{\mathbb{1}_{0}}\left[f(W) \mid \Gamma_{0}\right] \mathbb{P}_{\nu_{\lambda}}(B) \\
& =\mathbb{E}_{\nu_{\lambda}}[f(W) \mid \Gamma] \mathbb{P}_{\nu_{\lambda}}(B),
\end{align*}
$$

where the last equality is, one more time, justified by Lemma 6.3. This proves the claim under $\mathbb{P}_{\nu_{\lambda}}$.

To extend the claim to $\mathbb{P}_{\nu_{\lambda}}(\cdot \mid \Gamma)$, note that $\Gamma \in \mathcal{G}_{\tau}$ since

$$
\begin{equation*}
\Gamma \cap\{K=n\}=\left\{\xi_{0}(0)=1, W_{s} \in\left[L_{s}(0), R_{s}(0)\right] \forall s \in\left[0, \mathcal{T}_{n}\right]\right\} \cap\{K=n\} \tag{6.20}
\end{equation*}
$$

and apply (6.19) to $B \cap \Gamma$ instead of $B$.

### 6.3 Proof of Proposition 6.2

Exponential moments. We first show that $T_{0}$ has exponential moments when it is finite, uniformly for $\lambda$ in compact sets. Fix $\lambda_{-}, \lambda_{+} \in\left(\lambda_{W}, \infty\right)$.

Lemma 6.4. For every $\lambda \in\left[\lambda_{-}, \lambda_{+}\right]$and $\epsilon>0$ there exists an $a=a\left(\lambda_{-}, \lambda_{+}, \epsilon\right)>0$ such that, for any probability measure $\mu$ stochastically larger than $\bar{\nu}_{\lambda}$,

$$
\begin{align*}
& \text { (a) } \mathbb{E}_{\mu}\left[\mathbb{1}_{\left\{T_{0}<\infty\right\}} e^{a T_{0}}\right] \leq 1+\epsilon .  \tag{6.21}\\
& \text { (b) } \mathbb{E}_{\mu}\left[e^{a V_{0}(0)}\right] \leq 1+\epsilon
\end{align*}
$$

Proof. We couple systems with infection rates $\lambda_{-}, \lambda$ and $\lambda_{+}$starting, respectively, from $\bar{\nu}_{\lambda_{-}}, \mu$ and $\mathbf{1}$, by coupling their initial configurations and their infection events monotonically. Denote their joint law by $\mathbb{P}$. In what follows, we will refer to these systems by their rates and we will use a superscript to indicate on which system a random variable depends.

We will bound $T_{0} \mathbb{1}_{\left\{T_{0}<\infty\right\}}=T_{0} \mathbb{1}_{\left\{F_{0}<\infty\right\}}$ by a time $D_{0}$ that depends only on systems $\lambda_{ \pm}$and has exponential moments under $\mathbb{P}$. We start by bounding $F_{0} \mathbb{1}_{\left\{F_{0}<\infty\right\}}$ by a variable $D_{1}$ depending only on system $\lambda_{-}$. Let

$$
\begin{equation*}
r_{t}:=\sup _{x \in \mathbb{Z} \leq 0} R_{t}(x), \quad l_{t}:=\inf _{x \in \mathbb{Z} \geq 0} L_{t}(x) . \tag{6.22}
\end{equation*}
$$

Then $r_{t}$ is the same as $r_{0, t}(0)$ in (6.3) when all sites in $\mathbb{Z}_{\leq 0}$ are infected, and analogously for $l_{t}$. Furthermore, $R_{t}(0), L_{t}(0)$ are equal to $r_{t}, l_{t}$ while $C_{t}(0) \neq \emptyset$ : this can be seen by using the graphical representation (see e.g. Liggett [12] Chapter VI, Theorem 2.2). Therefore

$$
\begin{equation*}
F_{0}=\inf \left\{t \geq 0: r_{t}<W_{t} \text { or } l_{t}>W_{t}\right\} . \tag{6.23}
\end{equation*}
$$

Let $m:=\frac{1}{2}\left(\iota\left(\lambda_{-}\right)+\left|v_{0}\right| \vee\left|v_{1}\right|\right)$. Take homogeneous random walks $X^{i}$ jumping at rates $\alpha_{i}, \beta_{i}, i \in\{0,1\}$, independent of $\xi$ and coupled to $W$ in such a way that $X_{t}^{0} \leq W_{t} \leq X_{t}^{1}$ for all $t \geq 0$. Set

$$
\begin{align*}
& D_{1 a}:=\sup \left\{t \geq 0: l_{t}^{\lambda-} \geq-m t \text { or } r_{t}^{\lambda-} \leq m t\right\},  \tag{6.24}\\
& D_{1 b}:=\sup \left\{t \geq 0:\left|X_{t}^{0}\right| \vee\left|X_{t}^{1}\right|>m t\right\} .
\end{align*}
$$

Then $D_{1 a}$ depends only on system $\lambda_{-}$and has exponential moments by known large deviation bounds for $r_{t}$ (see Liggett [12] Chapter VI, Corollary 3.22), while $D_{1 b}$ is independent of $\xi$ and has exponential moments by standard large deviation bounds for $X^{0}$ and $X^{1}$. Noting that $r_{t}$ and $l_{t}$ are monotone, we can take $D_{1}:=D_{1 a} \vee D_{1 b}$, which does not depend on the initial configuration.

Set $\delta:=\frac{1}{2}\left(\iota\left(\lambda_{-}\right)-m\right), x_{0}:=Y_{0, D_{1}}^{\lambda_{+}}(0)-\left\lceil\left(\iota\left(\lambda_{+}\right)+\delta\right) D_{1}\right\rceil$ and note, using the graphical representation, that $\Delta_{0}:=x_{0}-Z_{\delta}^{\lambda-}\left(x_{0}\right)$ is independent of $x_{0}$, where $Z_{\delta}(x)$ is as in (4.5). Then

$$
\begin{equation*}
D_{0}:=\frac{\Delta_{0}+\left|x_{0}\right|+1}{\iota\left(\lambda_{-}\right)-\delta-m}=4 \frac{\Delta_{0}+\left|x_{0}\right|+1}{\iota\left(\lambda_{-}\right)-\left|v_{0}\right| \vee\left|v_{1}\right|} \tag{6.25}
\end{equation*}
$$

depends only on $\lambda_{-}, \lambda_{+}$and has exponential moments under $\mathbb{P}$ by Lemma 4.3. It is easy to check that $D_{0}$ is the intersection time of the line of inclination $\iota\left(\lambda_{-}\right)-\delta$ passing through $\left(Z_{\delta}^{\lambda-}\left(x_{0}\right)-1,0\right)$ and the line of inclination $m$ passing through the origin. Since system $\lambda$ has more infections than system $\lambda_{-}$and $D_{0} \geq D_{1}$, we have $T_{0} \mathbb{1}_{\left\{T_{0}<\infty\right\}} \leq D_{0}$, which proves (a). For (b), we can bound $V_{0}(0)$ analogously, taking $x_{0}=0$ instead.

Infections at trial times. We next show that at trial times there are more infections to the left of the random walk than under $\bar{\nu}_{\lambda}$.

Lemma 6.5. For all $n \in \mathbb{N}$, on the event $\left\{\mathcal{T}_{n}<\infty\right\}$ the law of $\xi_{\mathcal{T}_{n}}\left(\cdot+W_{\mathcal{T}_{n}}\right)$ under $\mathbb{P}_{\nu_{\lambda}}\left(\cdot \mid \mathcal{T}_{[1, n]}, W_{\left[0, \mathcal{T}_{n}\right]}\right)$ a.s. is stochastically larger than $\bar{\nu}_{\lambda}$.

Proof. Suppose that $n \geq 2$ (the case $n=1$ is simpler). Using the definition of $\mathcal{T}_{n}$, we can show by induction that, if $\mathcal{T}_{n}<\infty$, then there exist infection paths connecting $\left(W_{\mathcal{T}_{n}}, \mathcal{T}_{n}\right)$ to $\mathbb{Z}_{\leq-1} \times\{0\}$ and never touching the paths $Y^{\mathcal{T}_{k}}\left(W_{\mathcal{T}_{k}}\right), k=1, \ldots, n-1$, or the region to the right of $W$. Take $\pi$ to be the right-most of these infection paths. Then $\pi$ is a random infection path with properties (p1) and (p2), and

$$
\begin{equation*}
\left(\mathcal{T}_{[1, n]}, W_{\left[0, \mathcal{T}_{n}\right]}\right) \in \mathcal{R}_{\mathcal{T}_{n}}^{\pi} \vee \sigma\left(N_{\left[0, \mathcal{T}_{n}\right]}, U_{\left[1, N_{\mathcal{T}_{n}}\right]}\right) . \tag{6.26}
\end{equation*}
$$

Therefore the result follows from Lemma 4.1.

Conclusion. We are now ready to prove Proposition 6.2.
Proof. Let

$$
\begin{equation*}
\kappa:=\mathbb{P}_{\mathbb{1}_{0}}\left(\Gamma_{0}\right) . \tag{6.27}
\end{equation*}
$$

By Lemma 6.3, $\mathbb{P}_{\nu_{\lambda}}(\Gamma)=\kappa \rho_{\lambda} \geq \kappa \rho_{\lambda_{-}}$by monotonicity (recall the definition of $\rho_{\lambda}$ from Section 1.2). Also, there exists a $\kappa_{-}>0$ such that $\kappa \geq \kappa_{-}$for any $\lambda \geq \lambda_{-}$: we can
take $\kappa_{-}$to be the probability that $X^{0}$ and $X^{1}$ in the proof of Lemma 6.4 never cross $L(0)$ or $R(0)$ in system $\lambda_{-}$. Therefore it is enough to prove the claim for $\mathbb{P}_{\nu_{\lambda}}$. Since $|W|$ is dominated by $N$, which is Poisson process independent of $\xi$, we only need to worry about $\tau$.

For $\epsilon>0$ such that $(1+\epsilon)\left(1-\kappa_{-}\right)<1$, take $a>0$ as in Lemma 6.4. On the event $\left\{\mathcal{T}_{n}<\infty\right\}$, let $\hat{\xi}_{n}:=\xi_{\mathcal{T}_{n}}\left(\cdot+W_{\mathcal{T}_{n}}\right)$ and note that, given $\mathcal{G}_{\mathcal{T}_{n}}, \mathcal{T}_{n+1}-\mathcal{T}_{n}$ is distributed as $T_{0}$ under $\mathbb{P}_{\hat{\xi}_{n}}$. By Lemma 6.5, the law of $\hat{\xi}_{n}$ under $\mathbb{P}_{\nu_{\lambda}}\left(\cdot \mid \mathcal{T}_{[1, n]}, W_{\left[0, \mathcal{T}_{n}\right]}\right)$ is stochastically larger than $\bar{\nu}_{\lambda}$, and we get from Lemma 6.4 that

$$
\begin{align*}
\mathbb{E}_{\nu_{\lambda}} & {\left[\mathbb{1}_{\left\{\mathcal{T}_{n+1}<\infty\right\}} e^{a\left(\mathcal{T}_{n+1}-\mathcal{T}_{n}\right)} \mid \mathcal{T}_{[1, n]}, W_{\left[0, \mathcal{T}_{n}\right]}\right] } \\
& =\mathbb{E}_{\nu_{\lambda}}\left[\mathbb{E}_{\hat{\xi}_{n}}\left[\mathbb{1}_{\left\{T_{0}<\infty\right\}} e^{a T_{0}}\right] \mid \mathcal{T}_{[1, n]}, W_{\left[0, \mathcal{T}_{n}\right]}\right] \leq 1+\epsilon \tag{6.28}
\end{align*}
$$

Using this bound, estimate

$$
\begin{align*}
\mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\left\{\mathcal{T}_{n+1}<\infty\right\}} e^{a \mathcal{T}_{n+1}}\right] & =\mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}} e^{a \mathcal{T}_{n}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\left\{\mathcal{T}_{n+1}<\infty\right\}} e^{a\left(\mathcal{T}_{n+1}-\mathcal{T}_{n}\right)} \mid \mathcal{T}_{n}\right]\right]  \tag{6.29}\\
& \leq(1+\epsilon) \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}} e^{a \mathcal{T}_{n}}\right],
\end{align*}
$$

so that, by induction,

$$
\begin{equation*}
\mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}} e^{a \mathcal{T}_{n}}\right] \leq(1+\epsilon)^{n} . \tag{6.30}
\end{equation*}
$$

Using Lemma 6.3, write, for $n \in \mathbb{N}$,

$$
\begin{equation*}
\mathbb{P}_{\nu_{\lambda}}(K \geq n+1)=\mathbb{P}_{\nu_{\lambda}}\left(\mathcal{T}_{n}<\infty, F_{\mathcal{T}_{n}}<\infty\right)=(1-\kappa) \mathbb{P}_{\nu_{\lambda}}(K \geq n) \tag{6.31}
\end{equation*}
$$

to note that $K$ has distribution $\operatorname{GEO}(\kappa)$. To conclude, use (6.30)-(6.31) to write

$$
\begin{align*}
\mathbb{E}_{\nu_{\lambda}}\left[e^{\frac{a}{2} \tau}\right] & =\sum_{n \in \mathbb{N}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\{K=n\}} e^{\frac{a}{2} \mathcal{T}_{n}}\right]=\sum_{n \in \mathbb{N}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\{K=n\}} \mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}} e^{\frac{a}{2} \mathcal{T}_{n}}\right] \\
& \leq \sum_{n \in \mathbb{N}} \mathbb{P}_{\nu_{\lambda}}(K=n)^{\frac{1}{2}} \mathbb{E}_{\nu_{\lambda}}\left[\mathbb{1}_{\left\{\mathcal{T}_{n}<\infty\right\}} e^{a \mathcal{T}_{n}}\right]^{\frac{1}{2}}  \tag{6.32}\\
& \leq\left(1-\kappa_{-}\right)^{-\frac{1}{2}} \sum_{n \in \mathbb{N}}\left(\sqrt{\left(1-\kappa_{-}\right)(1+\epsilon)}\right)^{n}<\infty,
\end{align*}
$$

where in the second line we use the Cauchy-Schwarz inequality.

### 6.4 Continuity of the speed and the volatility

Given $\lambda_{-} \leq \lambda_{+}$in $\left(\lambda_{W}, \infty\right)$ and $\left(\lambda_{n}\right)_{n \in \mathbb{N}}, \lambda_{*} \in\left[\lambda_{-}, \lambda_{+}\right]$such that either $\lambda_{n} \uparrow \lambda_{*}$ or $\lambda_{n} \downarrow \lambda_{*}$ as $n \rightarrow \infty$, we can simultaneously construct systems with infection rates $\left(\lambda_{n}\right)_{n \in \mathbb{N}}, \lambda_{*}$ and $\lambda_{ \pm}$, starting from equilibrium, with a single graphical representation in the standard fashion, taking a monotone sequence of Poisson processes for infection events and coupling the initial configurations monotonically. For $n \in \mathbb{N} \cup\{*,+,-\}$, denote by $\Lambda^{n}:=\left(\xi_{0}^{n}, H, I^{n}, N, U\right)$ the system with infection rate $\lambda_{n}$, and by $\mathbb{P}$ their joint law. In the following, we will use a superscript $n$ to indicate functionals of $\Lambda^{n}$.

In view of (6.13) and Proposition 6.2, in order to prove convergence of $v\left(\lambda_{n}\right)$ and $\sigma\left(\lambda_{n}\right)$ it is enough to prove convergence in distribution of $\Gamma^{n}$ and of $\left(W_{\tau^{n}}^{n}, \tau^{n}\right) \mathbb{1}_{\Gamma^{n}}$.

The main step to achieve this will be to approximate relevant random variables with uniformly large probability by random variables depending on bounded regions of the graphical representation.

Note that, by monotonicity and continuity of $\lambda \mapsto \rho_{\lambda}$ (see Liggett [12] Chapter VI, Theorem 1.6),

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \xi_{0}^{n}(x)=\xi_{0}^{*}(x) \quad \forall x \in \mathbb{Z} \quad \mathbb{P} \text {-a.s. } \tag{6.33}
\end{equation*}
$$

Recall the definitions of $F_{0}, \mathcal{T}_{k}$ and $K$ in (6.5), (6.8)-(6.9) and (6.10), respectively. For $n \in \mathbb{N} \cup\{*\}$ and $k \in \mathbb{N}$, let

$$
\begin{equation*}
\Gamma_{k}^{n}:=\left\{\xi_{0}^{n}(0)=1, W_{s}^{n} \in\left[L_{s}^{n}(0), R_{s}^{n}(0)\right] \forall s \in\left[0, \mathcal{T}_{k}^{n}\right] \cap \mathbb{R}\right\} \tag{6.34}
\end{equation*}
$$

so that $\Gamma^{n}=\Gamma_{k}^{n}$ on $\left\{K^{n}=k\right\}$ as in (6.20).
Proposition 6.6. For every $k \in \mathbb{N}$, $\left(W_{\mathcal{T}_{k}^{n}}^{n}, \mathcal{T}_{k}^{n}, \mathbb{1}_{\Gamma_{k}^{n}}\right) \mathbb{1}_{\left\{\mathcal{T}_{k}^{n}<\infty\right\}}, \mathbb{1}_{\left\{\mathcal{T}_{k}^{n}<\infty\right\}}$ and $\mathbb{1}_{\left\{F_{0}^{n}<\infty\right\}}$ converge in distribution as $n \rightarrow \infty$ to the corresponding functionals of $\Lambda^{*}$.

Proof. We first show that, for every fixed $T \in(0, \infty)$,

$$
\begin{equation*}
\left(W_{\mathcal{T}_{k}^{n}}^{n}, \mathcal{T}_{k}^{n}, \mathbb{1}_{\Gamma_{k}^{n}}\right) \mathbb{1}_{\left\{\mathcal{T}_{k}^{n} \leq T\right\}}, \quad \mathbb{1}_{\left\{\mathcal{T}_{k}^{n} \leq T\right\}}, \quad \mathbb{1}_{\left\{F_{0}^{n} \leq T\right\}}, \tag{6.35}
\end{equation*}
$$

converge a.s. as $n \rightarrow \infty$ to the corresponding functionals of $\Lambda^{*}$. To that end, let $\bar{Y}_{t, s}(x)$ be the increasing analogue of $Y_{t, s}(x)$ in (6.6), starting from $x$ but jumping across the arrows of $I$ to the right. Let $\bar{Z}_{\delta}(x)$, analogously to $Z_{\delta}(x)$ in (4.5), be the first infected site to the right of $x$ whose infection spreads inside a wedge between lines of inclination $-(\iota+\delta)$ and $-(\iota-\delta)$. Take $\delta:=\iota\left(\lambda_{-}\right) / 2$, set $\underline{y}:=Y_{0, T}^{+}\left(-N_{T}\right)$ and $\underline{z}:=$ $Z_{\delta}^{-}\left(\underline{y}-\left\lceil\left(\iota\left(\lambda_{-}\right)+\delta\right) T\right\rceil\right)$. Analogously, put $\bar{y}:=\bar{Y}_{0, T}^{+}\left(N_{T}\right)$ and $\bar{z}:=\bar{Z}_{\delta}^{-}\left(\bar{y}+\left\lceil\left(\iota\left(\lambda_{-}\right)+\delta\right) T\right\rceil\right)$.

Now observe that, for any $n \in \mathbb{N} \cup\{*\}$, all random variables in (6.35) depend on $\Lambda^{n}$ only in the space-time box $\mathcal{B}:=[\underline{z}, \bar{z}] \times[0, T]$. Indeed, for any $0 \leq t \leq s \leq T$, we have $L_{t, s}^{n}\left(W_{t}^{n}\right) \geq Y_{t, s}^{n}\left(W_{t}^{n}\right) \geq y^{-}$and $R_{t, s}^{n}\left(W_{t}^{n}\right) \leq y^{+}$, so that $\left\{F_{t}^{n} \leq s\right\}$ depends on $\Lambda^{n}$ only inside $[y, \bar{y}] \times[0, T]$. Also, there are infection paths from time 0 to time $T$ inside $[\underline{z}, \underline{y})$ and $(\bar{y}, \bar{z}]$. Therefore the states of $\xi^{n}$ inside $[\underline{y}, \bar{y}] \times[0, T]$ depend on $\Lambda^{n}$ only in $\mathcal{B}$ (see the proof of Lemma 2.2). The same is true for $\left\{T_{t}^{n} \leq s\right\}$, since any infection path needed to discover $T_{t}^{n}$ can be taken inside $\mathcal{B}$. Therefore, by (6.33) (and since the graphical representation is a.s. eventually constant inside bounded space-time regions), the claim after (6.35) follows.

To conclude note that, because $\mathcal{T}_{k} \mathbb{1}_{\left\{\mathcal{T}_{k}<\infty\right\}} \leq \tau$ and $F_{0} \mathbb{1}_{\left\{F_{0}<\infty\right\}} \leq T_{0} \mathbb{1}_{\left\{T_{0}<\infty\right\}}$,

$$
\begin{equation*}
\lim _{T \rightarrow \infty} \sup _{n \in \mathbb{N} \cup\{*\}} \mathbb{P}\left(T<\mathcal{T}_{k}^{n}<\infty \text { or } T<F_{0}^{n}<\infty\right)=0 \tag{6.36}
\end{equation*}
$$

by Proposition 6.2 and Lemma 6.4, which implies that, for large $T$, the random variables in the statement are equal to the ones in (6.35) with uniformly large probability.

Corollary 6.7. Let $\kappa^{n}$ be as in (6.27). Then $\lim _{n \rightarrow \infty} \kappa^{n}=\kappa^{*}$ and $K^{n}$ converges in distribution to $K^{*}$.

Proof. This follows directly from Proposition 6.6 and the definition of $\kappa$ since, by (6.31), $K^{n}$ is a geometric random variable with parameter $\kappa^{n}$.

With these results we can conclude the proof of Theorem 1.2(c).
Proof. Let $f$ be a bounded measurable function. For $k \in \mathbb{N}$, write

$$
\begin{align*}
\mathbb{E}\left[f\left(W_{\tau^{n}}^{n}, \tau^{n}\right) \mathbb{1}_{\Gamma^{n}} \mathbb{1}_{\left\{K^{n}=k\right\}}\right] & =\mathbb{E}\left[f\left(W_{\mathcal{T}_{k}^{n}}^{n}, \mathcal{T}_{k}^{n}\right) \mathbb{1}_{\Gamma_{k}^{n}} \mathbb{1}_{\left\{\mathcal{T}_{k}^{n}<\infty, F_{\mathcal{T}_{k}^{n}=\infty}\right\}}\right] \\
& =\kappa^{n} \mathbb{E}\left[f\left(W_{\mathcal{T}_{k}^{n}}^{n}, \mathcal{T}_{k}^{n}\right) \mathbb{1}_{\Gamma_{k}^{n}} \mathbb{1}_{\left\{\mathcal{T}_{k}^{n}<\infty\right\}}\right]  \tag{6.37}\\
& \xrightarrow{n \rightarrow \infty} \kappa^{*} \mathbb{E}\left[f\left(W_{\mathcal{T}_{k}^{*}}^{*}, \mathcal{T}_{k}^{*}\right) \mathbb{1}_{\Gamma_{k}^{*}} \mathbb{1}_{\left\{\mathcal{T}_{k}^{*}<\infty\right\}}\right] \\
& =\mathbb{E}\left[f\left(W_{\tau^{*}}^{*}, \tau^{*}\right) \mathbb{1}_{\Gamma^{*}} \mathbb{1}_{\left\{K^{*}=k\right\}}\right],
\end{align*}
$$

where for the second and the third equality we use Lemma 6.3 and the strong Markov property, and for the convergence we use Proposition 6.6 and Corollary 6.7. Therefore

$$
\begin{align*}
& \left|\mathbb{E}\left[f\left(W_{\tau^{n}}^{n}, \tau^{n}\right) \mathbb{1}_{\Gamma^{n}}\right]-\mathbb{E}\left[f\left(W_{\tau^{*}}^{*}, \tau^{*}\right) \mathbb{1}_{\Gamma^{*}}\right]\right| \\
& \quad \leq\|f\|_{\infty}\left\{\mathbb{P}\left(K^{n}>M\right)+\mathbb{P}\left(K^{*}>M\right)\right\} \\
& \quad+\sum_{k=1}^{M}\left|\mathbb{E}\left[f\left(W_{\tau^{n}}^{n}, \tau^{n}\right) \mathbb{1}_{\Gamma^{n}} \mathbb{1}_{\left\{K^{n}=k\right\}}\right]-\mathbb{E}\left[f\left(W_{\tau^{*}}^{*}, \tau^{*}\right) \mathbb{1}_{\Gamma^{*}} \mathbb{1}_{\left\{K^{*}=k\right\}}\right]\right| \tag{6.38}
\end{align*}
$$

and we conclude by taking $n \rightarrow \infty$, using Corollary 6.7 and (6.37), and taking $M \rightarrow \infty$.

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