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Abstract

This paper presents an analysis of the stochastic recursion $W_{i+1} = [V_iW_i + Y_i]^+$ that can be interpreted as an autoregressive process of order 1, reflected at 0. We start our exposition by a discussion of the model's stability condition. Writing $Y_i = B_i - A_i$, for independent sequences of non-negative i.i.d. random variables $\{A_i\}_{i \in \mathbb{N}_0}$ and $\{B_i\}_{i \in \mathbb{N}_0}$, and assuming $\{V_i\}_{i \in \mathbb{N}_0}$ is an i.i.d. sequence as well (independent of $\{A_i\}_{i \in \mathbb{N}_0}$ and $\{B_i\}_{i \in \mathbb{N}_0}$), we then consider three special cases: (i) V_i attains negative values only and B_i has a rational LST, (ii) V_i equals a positive value *a* with certain probability $p \in (0, 1)$ and is negative otherwise, and both A_i and B_i have a rational LST, (iii) V_i is uniformly distributed on [0, 1], and A_i is exponentially distributed. In all three cases we derive transient and stationary results, where the transient results are in terms of the transform at a geometrically distributed epoch.

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1. Introduction

This paper focuses on the Lindley type stochastic recursion

$$W_{i+1} = [V_i W_i + Y_i]^+, \quad i = 0, 1, \dots,$$
(1)

where $[x]^+ = \max\{x, 0\}$ for $x \in \mathbb{R}$ and $\{V_i\}_{i \in \mathbb{N}_0}$ and $\{Y_i\}_{i \in \mathbb{N}_0}$ are independent sequences of i.i.d. (independent, identically distributed) random variables. The analysis of stochastic recursions has received much attention in the applied probability literature. This holds in particular for stochastic recursions of the autoregressive type, owing to their wide applicability across various scientific domains including biology, finance, and engineering [6, 8, 13].

An important subclass of first order autoregressive models corresponds to the case in which the $\{V_i\}_{i \in \mathbb{N}_0}$ are constant, i.e., a stochastic process defined through the recursion

$$W_{i+1} = aW_i + Y_i, \quad i = 0, 1, \dots,$$
 (2)

for a sequence of i.i.d. random variables $\{Y_i\}_{i \in \mathbb{N}_0}$ and a scalar *a*, with W_0 being given. When the quantities W_i cannot attain negative values, then it becomes natural to study the truncated counterpart of (2), i.e., the recursion

$$W_{i+1} = [aW_i + Y_i]^+, \quad i = 0, 1, \dots$$
(3)

When a = 1 we recover the classical Lindley recursion describing the waiting time in the G/G/1 queue, with Y_i representing the difference between the *i*-th service time and the (i + 1)-st interarrival time. The case of $a \in (0, 1)$ was studied in detail in [6], whereas the case a = -1 is covered by [22]. It should be observed that, while from the analysis it is clear that a is assumed to be positive in [6], the introduction of that paper incorrectly states that |a| < 1.

By studying (1), we significantly extend the analysis of the Lindley recursion as well as the analysis of the stochastic recursion (3). Our results focus on three different choices of $\{V_i\}$. In Model I, the V_i are negative random variables. A detailed analysis is shown to be possible as long as the positive part of the Y_i has a rational Laplace-Stieltjes transform (in the sequel abbreviated to LST). In Model II, the V_i are either negative or equal to the positive constant a. Here we demand that both the positive and the negative part of the Y_i have a rational LST. While Model II effectively contains Model I, we prefer to give a separate analysis of both models, to make the reader familiar with the specific mathematical intricacies due to V_i being negative (Model I) and V_i being a positive constant (the model in [6]). Finally, in Model III, the V_i are uniformly distributed on [0, 1], and the negative part of the Y_i is exponentially distributed; this case requires an entirely different approach.

The rationality assumptions are natural in the light of the existing theory that has been developed for the G/G/1 queue. While in principle the waiting-time distribution in the general G/G/1 queue can be obtained via a Wiener-Hopf decomposition (cf. [10, Chapter II.5]), the solution is a rather implicit one, unless one makes rationality assumptions on either the interarrival or the service-time LST. In addition, it can be argued that the distribution of any non-negative random variable can be approximated arbitrarily closely by the distribution of a random variable with a rational LST [1, Ch. III], so that a restriction to random variables with rational LST leads to just a minor loss of generality.

Notable studies of stochastic recursions are Borovkov and Foss [4], Foss et al. [15], and Diaconis and Freedman [12]; see in addition [3]. For the non-reflected case, stochastic recursions of the form $W_{i+1} = V_i W_i + Y_i$ have been studied frequently, partly under the name 'Vervaat perpetuity'; we mention

[8, 11, 13, 16, 18, 21]. For the reflected case, [7] considers another generalization of the Lindley recursion, by replacing V_iW_i in (1) by $S(W_i)$, where $\{S(t)\}_{t\geq0}$ is a Lévy subordinator. A model that is similar to the present model has been discussed in Whitt [23]. It is noted, though, that [23] primarily focuses on stability questions, limit theorems and questions related to queuing applications, whereas our primary focus lies on the derivation of results for the transient and stationary distribution of the process under investigation. Also related is the model in [5]; there (1) is considered with $\mathbb{P}(V_i = 1) = p$, $\mathbb{P}(V_i = -1) = 1 - p$.

The main contributions of the present paper are the following. For each of the three models that we introduced above, we state and solve a Wiener-Hopf boundary value problem, which allows us to study the transient behavior of the $\{W_i\}_{i \in \mathbb{N}_0}$ process. In particular, we obtain an expression for the object

$$\sum_{i=0}^{\infty} r^i \mathbb{E}(\mathrm{e}^{-sW_i}),$$

which can be interpreted as the generating function of the LST of the W_i , but also (up to the multiplicative constant 1 - r) as the LST after a geometrically distributed time. The stability condition of each of the three models is discussed, and the steady-state distribution of the $\{W_i\}_{i \in \mathbb{N}_0}$ process is also determined.

The remainder of the paper is organized as follows. Section 2 presents the description of the three models and some preliminaries. Sections 3, 4 and 5 are devoted to the transient and steady-state analysis of, respectively, Models I, II and III. Section 6 contains a discussion and provides suggestions for further research.

2. Model description and preliminaries

The main object of study is the stochastic recursion

$$W_{i+1} = [V_i W_i + Y_i]^+, \quad i = 0, 1, \dots,$$
(4)

where $\{V_i\}_{i \in \mathbb{N}_0}$ and $\{Y_i\}_{i \in \mathbb{N}_0}$ are sequences of i.i.d. random variables, which are in addition independent of each other. The initial state of the process is assumed to be $W_0 = w \in \mathbb{R}^+$. We write V and Y for generic random variables distributed as V_0 and Y_0 respectively.

In this paper, we discuss the following three variants of the model:

Model I:
$$\mathbb{P}(V < 0) = 1$$
;
Model II: $\mathbb{P}(V = a) = p$, $\mathbb{P}(V < 0) = 1 - p$, $a > 0$, $p \in (0, 1)$;
Model III: $\mathbb{P}(V < x) = x$, $0 \le x \le 1$.

In each of the cases we will assume that the Y_i are decomposed as $B_i - A_i$, with sequences $\{A_i\}_{i \in \mathbb{N}_0}$ and $\{B_i\}_{i \in \mathbb{N}_0}$ of i.i.d. *non-negative* random variables. In addition, depending on the chosen model, the random variables A_i and/or B_i are assumed to have a rational LST.

We start with investigating the stationary behavior of $\{W_i\}_{i \in \mathbb{N}_0}$. We always assume that both $\mathbb{E}|V|$ and $\mathbb{E}|Y|$ are finite. The first result was given in Whitt [23] and covers most cases of interest.

Theorem 1 (Whitt [23]). *If one of the following conditions holds, then* W_i *tends weakly to a proper limit* W *as* $i \rightarrow \infty$:

(C₁) $\mathbb{P}(V < 0) > 0$ and $\mathbb{P}(Y \le 0) > 0$,

(C₂) $V \ge 0$ *a.s.* and $\mathbb{P}(V = 0) > 0$,

(C₃) V > 0 a.s. and $\mathbb{E}(\log |V|) < 0$.

Moreover, W_i converges weakly to a possibly improper limit W as $i \to \infty$ if V > 0 a.s., $\mathbb{E}(\log |V|) = 0$, and $W_0 = 0$. If additionally V = 1 a.s. then W is proper for $\mathbb{E}(Y) < 0$ and improper for $\mathbb{E}(Y) > 0$.

It follows straightforwardly from the regenerative structure of W_i , and the proof of the above theorem in Whitt [23], that in cases (C₁) and (C₂) the limit *W* is unique. Obviously under the conditions of the theorem the limiting random variable *W* fulfils the associated distributional identity $W =_d [VW + Y]^+$. Regarding the above Condition (C₃), we add the following observation.

Theorem 2. In order to have convergence of W_i to a proper unique limit W as $i \to \infty$, it is sufficient to have $\mathbb{E}(\log |V|) < 0$, which in turn is implied by $\mathbb{E}|V| < 1$.

Proof. Recursion (4) can be written as a random iteration $W_{i+1} = f_{\theta_i}(W_i)$ with $f_{\theta_i}(x) = [xV_i + Y_i]^+$, $\theta_i = (V_i, Y_i)$. This means that $f_{\theta}(\cdot)$ enjoys the Lipschitz property

$$|f_{\theta}(x) - f_{\theta}(y)| \leq K_{\theta}|x - y|,$$

with random Lipschitz constant $K_{\theta} = |V|$. As a result, Diaconis and Freedman [12, Thm. 1.1] is applicable. By Jensen's inequality $\mathbb{E} \log |V| \le \log \mathbb{E} |V|$ and so $\mathbb{E} |V| < 1$ implies condition $\mathbb{E}(\log |V|) < 0$.

The case where $\mathbb{P}(V < 0) > 0$ and $Y \ge 0$ a.s., which was omitted in Whitt [23], is more involved due to the fact that the process might not be aperiodic, even if Y is not deterministic. As an example suppose that the distribution of Y is supported on [1,2] and that $V \le -2$ a.s. If $W_0 = 0$, then $W_1 \in [1,2], W_2 = 0, W_3 \in [1,2]$, entailing that the process alternates between the set {0} and a value in [1,2]. On the other hand, if $W_0 > 2$, then $W_1 = 0, W_2 \in [1,2]$ and so on. As a consequence, there is no convergence $W_i \Longrightarrow W$ as $i \to \infty$. However, regarding the existence of a stationary distribution we can show the following.

Theorem 3. If $\mathbb{P}(V \leq 0) > 0$ and $Y \geq 0$ a.s., then there is convergence of W_i to a stationary random variable W as $i \to \infty$.

Proof. We define a majorizing process by $M_0 := W_0$ and $M_{i+1} := [V_i]^+ M_i + Y_i$. Then $W_i \leq M_i$, i = 0, 1, 2, ... and for $\{M_i\}_{i \in \mathbb{N}_0}$ we have $\mathbb{E}(\log | [V]^+ |) = -\infty < 0$ since $\mathbb{P}([V]^+ = 0) > 0$. So by Thm. 2 it follows that $M_i \Rightarrow M$ as $i \to \infty$ for some limiting random variable M. This implies the existence of a stationary distribution for $\{W_i\}_{i \in \mathbb{N}_0}$. Indeed, let $\varepsilon > 0$ and let k > 0 be such that $\mathbb{P}(M > k) \leq \varepsilon/2$. Letting $b_i = \mathbb{P}(W_i \leq k)$, we obtain $b_i \geq \mathbb{P}(M_i \leq k)$. The right-hand side converges to $\mathbb{P}(M \leq k)$ as $i \to \infty$ which is at least $1 - \varepsilon/2$, implying that there is an i_0 such that for $i \geq i_0, b_i > 1 - \varepsilon$ for all $i \geq i_0$. Consequently the family $\mathbb{P}(W_i \leq \cdot)$ is tight, guaranteeing [14, Thm. 4] the existence of a stationary distribution for $\{W_i\}_{i\in\mathbb{N}_0}$.

We end this section with a lemma that forms the starting-point of the analysis of all three models. For this, we need to introduce some additional notation. For a given non-negative random variable X

we write $\Phi_X(s) = \mathbb{E}e^{-sX}$ for its LST, defined at least for Re $s \ge 0$. We say that $\Phi_X \in \mathbb{Q}[s_1, s_2, \dots, s_n]$ if *X* has a rational LST with poles at s_1, s_2, \dots, s_n , i.e., if $\Phi_X(s)$ is of the form

$$\Phi_X(s) = \frac{N_X(s)}{D_X(s)},\tag{5}$$

where $D_X(s) = \prod_{i=1}^n (s-s_i)$ and $N_X(s)$ is a polynomial of degree at most n-1 not sharing zeros with $D_X(s)$. Note that this implies that $\mathbb{P}(X = 0) = \lim_{s \to \infty} \Phi_X(s) = 0$. With this notation we then have e.g. $\Phi_Y(s) = \Phi_B(s) \Phi_A(-s)$. We also write $D_Y(s) = D_B(s) D_A(-s)$ and $N_Y(s) = N_B(s) N_A(-s)$ if *A* and *B* have rational LSTs of the form (5).

For a sequence $\{X_i\}_{i \in \mathbb{N}_0}$ of random variables we introduce the generating function, for $r \in (0, 1)$:

$$U_X(r,s) = \sum_{i=0}^{\infty} r^i \Phi_{X_i}(s).$$

Note that since $\Phi_{VX}(s) = \int \Phi_X(ys) \mathbb{P}(V \in dy)$, we have

$$U_{VX}(r,s) = \int U_X(r,ys) \mathbb{P}(V \in dy).$$
(6)

The following lemma plays a key role in our analysis. Define $W_i^* := [V_i W_i + Y_i]^-$, where $[x]^- := \min\{x, 0\}$.

Lemma 4. $U_W(r, s)$ and $U_{W^*}(r, s)$ are, for $r \in (0, 1)$ and Re s = 0, related via

$$U_W(r,s) = e^{-sw} + r \Big(\Phi_Y(s) U_{VW}(r,s) + \frac{1}{1-r} - U_{W^*}(r,s) \Big).$$
(7)

Proof. First observe that the basic identity $\exp([x]^+) = \exp(x) + 1 - \exp([x]^-)$ applies. It thus follows from (4) that

$$\Phi_{W_{i+1}}(s) = \Phi_{V_i W_i + Y_i}(s) + 1 - \Phi_{W_i^*}(s), \quad i = 0, 1, \dots$$
(8)

Multiplying both sides of (8) by r^{i+1} and summing yields the identity (7).

3. Model I: The negative case

The model we analyze in this section assumes that each V_i attains only negative values and that Y_i is the difference $B_i - A_i$ of two independent non-negative random variables, where B_i has a rational LST. In other words, we impose the conditions

(A) V < 0 a.s.,

(B) $\Phi_B \in \mathbb{Q}[s_1, \ldots, s_\ell]$ with $\operatorname{Re} s_j < 0$ for $j = 1, \ldots, \ell$.

Theorem 5. Suppose that the Conditions (A) and (B) hold. Then, for $r \in (0, 1)$,

$$U_W(r,s) = e^{-sw} + \frac{\sum_{k=0}^{\ell} a_k(r) s^k}{D_B(s)}, \quad \text{Re } s \ge 0,$$
(9)

$$U_{W^*}(r,s) = \frac{1}{1-r} - \frac{\sum_{k=0}^{\ell} a_k(r)s^k}{rD_B(s)} + \Phi_Y(s) \int_{-\infty}^0 U_W(r,sy) \mathbb{P}(V \in \mathrm{d}y), \quad \mathrm{Re}\, s \le 0, \tag{10}$$

where

$$a_0(r) = \frac{r}{1-r} (-1)^{\ell} \prod_{j=1}^{\ell} s_j,$$
(11)

and the remaining constants $a_1(r), \ldots, a_\ell(r)$ can be determined from the linear system (18) that will be given below.

Proof. Multiplying both sides of (7) by the denominator $D_B(s)$ gives

$$D_B(s) (U_W(r,s) - e^{-sw})$$

= $rN_B(s) \Phi_A(-s) U_{VW}(r,s) + rD_B(s) (\frac{1}{1-r} - U_{W^*}(r,s)).$ (12)

Now observe the following:

- (i) the left-hand side of (12) is analytic in Re s > 0 and continuous in $\text{Re } s \ge 0$,
- (ii) the right-hand side of (12) is analytic in Re s < 0 and continuous in Re $s \le 0$,
- (iii) for large s, both sides are $O(s^l)$ in their respective half-planes

At the boundary Re s = 0, both sides are well-defined. Determination of the unknown functions $U_W(r, s)$ and $U_{W^*}(r, s)$ from (12) and conditions (i), (ii) and (iii) is a Wiener-Hopf boundary value problem of a type that has been extensively studied in the queuing theory before, cf. the expository paper [9]. By introducing a function G(r, s) that is equal to the left-hand side of (12) for $\text{Re } s \ge 0$ and to the right-hand side of (12) for $\text{Re } s \le 0$, we have a function that is analytic in the whole *s*-plane, and that for large *s* is $O(s^{\ell})$. Liouville's theorem [20, p. 85] now states that both sides of (12), in their respective half-planes, are equal to the same ℓ -th degree polynomial in *s*. In other words,

$$D_B(s)(U_W(r,s) - e^{-sw}) = \sum_{k=0}^{\ell} a_k(r)s^k$$
(13)

for Re $s \ge 0$ and

$$rN_B(s) \Phi_A(-s) U_{VW}(r,s) + rD_B(s) \left(\frac{1}{1-r} - U_{W^*}(r,s)\right) = \sum_{k=0}^{\ell} a_k(r) s^k$$
(14)

for Re $s \le 0$. We still need to determine the $\ell + 1$ unknown functions $a_0(r), \ldots, a_\ell(r)$. Taking s = 0 in either (13) or (14) gives the expression in (11) for $a_0(r)$. Next we take $s = s_j, j = 1, \ldots, \ell$. We do this in (14), observing that Re $s_j < 0$. Using that $D_B(s_j) = 0$ we thus obtain

$$rN_B(s_j) \Phi_A(-s_j) U_{VW}(r, s_j) = \sum_{k=0}^{\ell} a_k(r) s_j^k, \quad j = 1, \dots, \ell.$$
(15)

Applying (6), this identity can be rewritten into

$$rN_B(s_j) \Phi_A(-s_j) \int_{-\infty}^0 U_W(r, s_j y) \mathbb{P}(V \in dy) = \sum_{k=0}^\ell a_k(r) s_j^k, \quad j = 1, \dots, \ell.$$
(16)

Using (13), Equation (16) becomes, for $j = 1, ..., \ell$,

$$rN_B(s_j) \Phi_A(-s_j) \int_{-\infty}^0 \left(e^{s_j y_W} + \frac{\sum_{k=0}^\ell a_k(r)(s_j y)^k}{D_B(s_j y)} \right) \ \mathbb{P}(V \in \mathrm{d}y) = \sum_{k=0}^\ell a_k(r) s_j^k.$$
(17)

We can rewrite this equation as follows: for $j = 1, ..., \ell$,

$$\sum_{k=0}^{\ell} a_k(r) s_j^k \left(1 - r N_B(s_j) \, \Phi_A(-s_j) \int_{-\infty}^0 \frac{y^k}{\prod_{m=1}^{\ell} (s_j y - s_m)} \, \mathbb{P}(V \in \mathrm{d}y) \right)$$

= $r \, N_B(s_j) \, \Phi_A(-s_j) \, \Phi_V(-s_j w).$ (18)

One can determine the remaining unknowns $a_1(r), \ldots, a_\ell(r)$ from this set of ℓ linear equations. Subsequently, from (13), (9) follows. Expression (10) then follows from (14).

We proceed by discussing the stationary behavior of $\{W_i\}_{i \in \mathbb{N}_0}$.

Theorem 6. Suppose that the Conditions (A) and (B) hold. If $\mathbb{P}(B \leq A) > 0$ then W_i converges weakly to a proper limit W as $i \to \infty$, and

$$\Phi_W(s) = \frac{\sum_{k=0}^{\ell} a_k s^k}{D_B(s)},$$
(19)

where

$$a_0 = (-1)^{\ell} \prod_{i=1}^{\ell} s_i,$$
(20)

and the remaining constants a_1, \ldots, a_ℓ can be determined from the linear system (22) that will be given below.

Proof. If $\mathbb{P}(B \le A) > 0$ holds then Condition (C₁) is fulfilled, so W_n weakly converges to a proper limit. We obtain the steady-state behavior (19) from its transient counterpart (9) in a standard manner, viz. by using an Abelian theorem for power series:

$$\Phi_W(s) = \lim_{r \uparrow 1} (1 - r) U_W(r, s) = \frac{\sum_{k=0}^{\ell} a_k s^k}{D_B(s)}, \quad \text{Re } s \ge 0,$$
(21)

where $a_k := \lim_{r \uparrow 1} (1 - r) a_k(r)$, for $k = 0, ..., \ell$. Using (11) and (18), we readily obtain the linear system

$$\sum_{k=0}^{\ell} a_k s_j^k \left(1 - N_B(s_j) \Phi_A(-s_j) \int_{-\infty}^0 \frac{y^k}{\prod_{m=1}^{\ell} (s_j y - s_m)} \mathbb{P}(V \in \mathrm{d}y) \right) = 0$$
(22)

for $a_j, j = 1, ..., \ell$.

The mean of *W* directly follows by differentiation of (19): $\Phi'_W(0) = (a_1 D_B(0) - a_0 D'_B(0))/D_B(0)^2$, with $D'_B(0) = 1$ if $\ell = 1$ and $D'_B(0) = -\sum_{i=1}^{\ell} s_i$ if $\ell = 2, 3, \ldots$; hence

$$\mathbb{E}(W) = \begin{cases} \frac{1-a_1}{a_0}, & \ell = 1; \\ -\frac{a_1 + \sum_{i=1}^{\ell} s_i}{a_0}, & \ell = 2, 3, \dots \end{cases}$$
(23)

Remark 1. It follows from (19) that *W* is a mixture of an atom at zero (with probability a_{ℓ}) and ℓ exponential terms. This is not surprising: as $V_i < 0$, the only way for W_{i+1} to be positive is to have $B_i > A_i - V_i W_i$. Now use the fact that B_i has a phase-type distribution with ℓ exponential phases, in combination with the memoryless property of the exponential distribution.

Remark 2. When $\ell = 1$, one obtains (using that $a_0 = -s_1$, cf. (20))

$$a_1 = \frac{1 - N_B(s_1)\Phi_A(-s_1)\int_{-\infty}^0 \frac{1}{y-1}\mathbb{P}(V \in dy)}{1 - N_B(s_1)\Phi_A(-s_1)\left(\int_{-\infty}^0 \frac{1}{y-1}\mathbb{P}(V \in dy) + 1\right)}.$$

For general ℓ , we have not been able to verify formally that the set of ℓ linear equations (22) in the ℓ unknowns a_1, \ldots, a_ℓ has a unique solution (as they involve the zeroes s_j and the distribution of V in an intricate way); similarly for the set of equations (18) for $a_1(r), \ldots, a_\ell(r)$. However, since W_i has a unique limiting distribution with LST $\Phi_W(s)$ as $i \to \infty$, there is no reason to suspect that anomalies in this set of equations will occur.

Example 1. Suppose that *B* has an exponential distribution with mean $1/\mu$. Then $\Phi_B(s) = \mu/(s+\mu)$, $\ell = 1$ and $s_1 = -\mu$. Suppose also that V = -a a.s. with a > 0. We then obtain

$$a_0(r) = \frac{r\mu}{1-r}, \quad a_1(r) = \frac{r}{1-r} \left(1 - \frac{(1+a)r\Phi_A(\mu)}{1+a+ar\Phi_A(\mu)} \right) - \frac{(1+a)r\Phi_A(\mu)}{1+a+ar\Phi_A(\mu)} e^{-a\mu w}.$$

Multiplying with (1 - r) and letting $r \uparrow 1$ yields the coefficients

$$a_0 = \mu,$$
 $a_1 = 1 - \frac{(1+a)\Phi_A(\mu)}{1+a+a\Phi_A(\mu)},$

so that the LST of W is given by

$$\Phi_W(s) = \frac{a_0 + a_1 s}{\mu + s} = \mathbb{P}(W > 0) \frac{\mu}{\mu + s} + \mathbb{P}(W = 0),$$
(24)

where the last equality follows from $\mathbb{P}(W = 0) = \lim_{s \to \infty} \Phi_W(s) = a_1$. We then obtain $\mathbb{E}(W) = (1 - a_1)/\mu$ in accordance with (23).

The case where a = 1, yielding the Lindley-type recursion $W_{i+1} = [B_i - A_i - W_i]^+$, has been extensively studied in [22]. We obtain for the stationary process

$$a_0 = \mu, \quad a_1 = \frac{2 - \Phi_A(\mu)}{2 + \Phi_A(\mu)},$$

which is in agreement with [22, Formula (4.12), p. 74]. It is easy to see that $\mathbb{P}(W = 0) = a_1$ is increasing in *a*.

For a = 0 we have $\mathbb{P}(W = 0) = 1 - \Phi_A(\mu)$. This relation is explained by observing that now $\mathbb{P}(W = 0) = \mathbb{P}(B < A)$, with $B \sim \exp(\mu)$. For $a \uparrow \infty$ we have $\mathbb{P}(W = 0) = 1/(1 + \Phi_A(\mu))$, which is explained by observing that a positive W is followed by a geometric(q) number of zeroes, with $q = \mathbb{P}(B < A) = 1 - \Phi_A(\mu)$.

4. Model II: The mixed case

In this section we consider the following variant of the model of Section 3. We again start from the recursion (4), but now assume that V = a, a > 0 with probability p and V < 0 with probability 1 - p. Let

$$V^- = (V \mid V < 0).$$

We keep the assumption that *B* has a rational LST, but add the requirement that *A* has a rational LST. Summarizing, we impose the conditions

- (A*) either V < 0 or V = a > 0 a.s.,
 - (B) $\Phi_B \in \mathbb{Q}[s_1, \ldots, s_\ell]$ with $\operatorname{Re} s_j < 0$ for $j = 1, \ldots, \ell$,
 - (C) $\Phi_A \in \mathbb{Q}[t_1, \dots, t_m]$ with $\operatorname{Re} t_i < 0$ for $i = 1, \dots, m$.

Theorem 7. Suppose that the Conditions (A*), (B), and (C) hold. Then, for $r \in (0, 1)$, *1. if* a = 1 *then*

$$U_W(r,s) = \frac{D_Y(s) e^{-sw} + \sum_{k=0}^{m+\ell} a_k(r) s^k}{D_Y(s) - rp N_Y(s)};$$
(25)

where

$$a_0(r) = \frac{r}{1-r} (1-p)(-1)^{\ell+m} \prod_{j=1}^{\ell} s_j \prod_{i=1}^{m} t_i,$$
(26)

while the remaining constants $a_1(r), \ldots, a_{m+\ell}(r)$ can be determined from the linear systems (32) and (34) that will be given below.

2. if $a \neq 1$ then

$$U_W(r,s) = \sum_{h=0}^{\infty} \left(e^{-a^h s_W} + \frac{\sum_{k=0}^{m+\ell} a_k(r)(a^h s)^k}{D_Y(a^h s)} \right) (rp)^h \prod_{j=0}^{h-1} \Phi_Y(a^j s),$$
(27)

where $a_0(r)$ is as in (26) and the remaining constants $a_1(r), \ldots, a_{m+\ell}(r)$ can be determined from the linear systems (42) and (43) that will be given below.

Proof. In this situation

$$U_{VW}(r,s) = p U_W(r,as) + (1-p) \int_{-\infty}^0 U_W(r,sy) \mathbb{P}(V^- \in dy).$$
(28)

Then (7) becomes, after multiplication by $D_Y(s)$,

$$D_{Y}(s) \left(U_{W}(r,s) - e^{-sw} \right) - rpN_{Y}(s) U_{W}(r,as)$$

= $r(1-p) N_{Y}(s) \int_{-\infty}^{0} U_{W}(r,sy) \mathbb{P}(V^{-} \in dy) + rD_{Y}(s) \left(\frac{1}{1-r} - U_{W^{*}}(r,s) \right).$ (29)

Now the following is true:

- (i) the left-hand side of (29) is analytic in Re s > 0 and continuous in $\text{Re } s \ge 0$,
- (ii) the right-hand side of (29) is analytic in Re s < 0 and continuous in Re $s \le 0$,

(iii) for large *s*, both sides are $O(s^{m+\ell})$ in their respective half-planes.

Again, both sides are well-defined at the boundary Re s = 0, so that we have a Wiener-Hopf boundary value problem. As before, the G(r, s) that is equal to the left-hand side of (29) for Re $s \ge 0$ and to the right-hand side of (29) for Re $s \le 0$ is analytic in the whole *s*-plane, and $G(r, s) = O(s^{m+\ell})$ for large *s*. According to Liouville's theorem both sides of (29), in their respective half-plane, are equal to the same $(m + \ell)$ -th degree polynomial in *s*, i.e., for Re $s \ge 0$

$$D_Y(s)(U_W(r,s) - e^{-sw}) - rpN_Y(s)U_W(r,as) = \sum_{k=0}^{m+\ell} a_k(r)s^k,$$
(30)

and for $\operatorname{Re} s \leq 0$

$$r(1-p) N_Y(s) \int_{-\infty}^0 U_W(r,sy) \mathbb{P}(V^- \in \mathrm{d}y) + r D_Y(s) \Big(\frac{1}{1-r} - U_{W^*}(r,s)\Big) = \sum_{k=0}^{m+\ell} a_k(r) s^k.$$
(31)

Taking s = 0 in either (30) or (31) yields, after a straightforward calculation, the expression for $a_0(r)$ in (26). Next we set $s = s_j$, $j = 1, ..., \ell$ in (31). Since $D_B(s_j) = 0$ it follows that

$$r(1-p)N_Y(s_j)\int_{-\infty}^0 U_W(r,s_jy) \mathbb{P}(V^- \in \mathrm{d}y) = \sum_{k=0}^{m+\ell} a_k(r)s_j^k, \quad j = 1,\dots,\ell.$$
(32)

We thus have obtained ℓ linear equations in the remaining $m + \ell$ unknown $a_k(r)$; however, they are expressed in the yet unknown function $U_W(r, \cdot)$.

We turn to (30), which provides a relation between $U_W(r, s)$ and $U_W(r, as)$. As it turns out, we have to distinguish between the two cases a = 1 and $a \neq 1$.

• *Case i:* For a = 1, after division by the denominators, Relation (30) can be rewritten as

$$U_W(r,s)(1 - rp\Phi_Y(s)) = e^{-sw} + \frac{\sum_{k=0}^{m+\ell} a_k(r)s^k}{D_Y(s)}.$$
(33)

Cohen [10], in his study of the $K_m/G/1$ queue, proves that the term between brackets in the left-hand side of (33) has *m* zeroes $\delta_1(r), \ldots, \delta_m(r)$ in the right half plane Re s > 0. The analyticity of $U_W(r, s)$ for Re $s \ge 0$ now implies that the right-hand side of (33) must be zero for all these *m* zeroes. This results in the *m* linear equations

$$\sum_{k=0}^{m+\ell} \delta_i^k(r) a_k(r) = -e^{-\delta_i(r)w} D_Y(\delta_i(r)), \quad i = 1, \dots, m.$$
(34)

Formula (32) contains ℓ more equations in the $a_k(r)$. Relying on (33), we can rewrite it into

$$r(1-p)N_Y(s_j) \int_{-\infty}^0 \frac{e^{-s_j y_W} + \sum_{k=0}^{m+\ell} a_k(r) \frac{(s_j y)^k}{D_Y(s_j y)}}{1 - rp \Phi_Y(s_j y)} \mathbb{P}(V^- \in \mathrm{d}y) = \sum_{k=0}^{m+\ell} a_k(r) s_j^k,$$
(35)

for $j = 1, \ldots, \ell$. From this we obtain

$$\sum_{k=0}^{m+\ell} c_k(r, s_j) a_k(r) = \int_{-\infty}^0 \frac{\mathrm{e}^{-s_j y w}}{1 - r p \Phi_Y(s_j y)} \mathbb{P}(V^- \in \mathrm{d}y), \quad j = 1, 2, \dots, \ell,$$
(36)

where, for $j = 1, \ldots, \ell$,

$$c_k(r, s_j) = \frac{s_j^k}{r(1-p) N_Y(s_j)} - \int_{-\infty}^0 \frac{(s_j y)^k}{D_Y(s_j y) - rp N_Y(s_j y)} \mathbb{P}(V^- \in dy).$$

• *Case ii:* For a < 1, Relation (30) has the same structure as [6, Formula (2.3)]. Proceeding in a similar way as in [6], we write

$$U_W(r,s) = K(r,s) U_W(r,as) + L(r,s),$$
(37)

with

$$K(r,s) := rp \,\Phi_Y(s), \quad L(r,s) := e^{-sw} + \frac{\sum_{k=0}^{m+\ell} a_k(r) \, s^k}{D_Y(s)}, \tag{38}$$

and iteration of (37) now yields

$$U_W(r,s) = \sum_{h=0}^{\infty} L(r,a^h s) \prod_{j=0}^{h-1} K(r,a^j s),$$
(39)

where convergence of the infinite sum can be proven using the d'Alembert test. Indeed, for a < 1 the limit as $h \to \infty$ of the ratio of two successive terms is

$$\lim_{h \to \infty} \left| \frac{L(r, a^h s)}{L(r, a^{h+1} s) K(r, a^h s)} \right| = \frac{1}{rp} > 1,$$
(40)

while for a > 1, $K(r, a^h s) \to 0$ and $|L(r, a^h s)| \to a_{m+\ell}(r)$, causing divergence of the lefthand side in (40) to infinity. Insertion of (38) in (39) gives (27). The only unknowns are $a_1(r), \ldots, a_{m+\ell}(r)$. We obtain *m* linear equations in the unknown $a_k(r)$ by observing that substitution of $s = -t_i$, $i = 1, \ldots, m$, in (30) results in the following identity:

$$-rpN_Y(-t_i)U_W(r, -at_i) = \sum_{k=0}^{m+\ell} a_k(r)(-t_i)^k, \quad i = 1, \dots, m.$$
(41)

Substituting the right-hand side of (27), with $s = -at_i$, into (41) now gives the *m* linear equations

$$\sum_{k=0}^{m+\ell} (-t_i)^k a_k(r) \left(1 + rp \, N_Y(-t_i) \sum_{h=0}^{\infty} \frac{a^{k(h+1)}(rp)^h \prod_{j=0}^{h-1} \Phi_Y(-a^{j+1}t_i)}{D_Y(-a^{h+1}t_i)} \right)$$
$$= -rp N_Y(-t_i) \sum_{h=0}^{\infty} e^{a^{h+1}t_i w} (rp)^h \prod_{j=0}^{h-1} \Phi_Y(-a^{j+1}t_i)$$
(42)

for i = 1, ..., m. The remaining ℓ equations are provided by substituting (27) into (32), yielding

$$\sum_{k=0}^{m+l} d_k(r, s_j) a_k(r)$$

= $r(1-p) N_Y(s_j) \sum_{h=0}^{\infty} (rp)^h \int_{-\infty}^0 e^{-a^h s_j y w} \prod_{i=0}^{h-1} \Phi_Y(a^i s_j y) \mathbb{P}(V^- \in dy)$ (43)

for $j = 1, \ldots, \ell$, where

$$d_k(r, s_j) = s_j^k \left(1 - r(1 - p) N_Y(s_j) \sum_{h=0}^{\infty} (rp)^h \int_{-\infty}^0 \frac{(a^h y)^k \prod_{i=0}^{h-1} \Phi_Y(a^i s_j y)}{D_Y(a^h s_j y)} \mathbb{P}(V^- \in dy) \right).$$

is hes the proof.

This finishes the proof.

The steady-state LST of W exists if $\mathbb{P}(B \leq A) > 0$, cf. Theorem 1. It can again be obtained by applying an Abelian theorem. For example in case (ii), in which a < 1, one gets

$$U_W(s) = \sum_{h=0}^{\infty} L(a^h s) \prod_{j=0}^{h-1} K(a^j s),$$
(44)

with

$$K(s) := p\Phi_Y(s), \quad L(s) := \frac{\sum_{k=0}^{m+l} a_k s^k}{D_Y(s)},$$
(45)

where $a_k := \lim_{r \neq 1} (1 - r) a_k(r)$. As this argumentation mimics the line of reasoning presented in the previous section, we omit details here.

5. Model III: the uniform proportional case

In this section we once more consider the stochastic recursion $W_{i+1} = [V_i W_i + Y_i]^+$, where $Y_i =$ $B_i - A_i$. Again we impose the usual independence assumptions on the sequences $\{V_i\}_{i \in \mathbb{N}_0}, \{A_i\}_{i \in \mathbb{N}_0}, \{A_$ and $\{B_i\}_{i\in\mathbb{N}_0}$. In addition, we assume that the A_i are $\exp(\lambda)$ distributed. The 'multiplicative adjustments' $\{V_i\}_{i \in \mathbb{N}_0}$ are assumed to form a sequence of unit uniformly distributed random variables on [0, 1]. By Thm. 2, since $\mathbb{E}(\log |V|) < 0$, a steady-state distribution of $\{W_i\}_{i \in \mathbb{N}_0}$ always exists. We shall first study its transient distribution, and then obtain the steady-state distribution.

We start with (8), i.e.,

$$\Phi_{W_{i+1}}(s) = \Phi_{V_i W_i + B_i - A_i}(s) + 1 - \Phi_{W_i^*}(s), \quad i = 0, 1, \dots,$$

where as before $W_i^* = [V_i W_i + B_i - A_i]^-$. This time the distribution of W_i^* is almost trivial: either $V_iW_i + B_i - A_i \ge 0$, in which case $W_i^* = 0$, or W_i^* has the same exponential distribution as A_i , due to the lack of memory property of the exponential distribution.

Using the independence between $\{A_i\}_{i \in \mathbb{N}_0}$, $\{B_i\}_{i \in \mathbb{N}_0}$, and $\{V_i\}_{i \in \mathbb{N}_0}$ and the exponentiality of the A_i , we obtain

$$\Phi_{W_{i+1}}(s) = \Phi_{V_i W_i}(s) \Phi_B(s) \frac{\lambda}{\lambda - s} + 1 - \mathbb{P}(V_i W_i + B_i - A_i < 0) \frac{\lambda}{\lambda - s} - \mathbb{P}(V_i W_i + B_i - A_i \ge 0)$$

$$= \Phi_{V_i W_i}(s) \Phi_B(s) \frac{\lambda}{\lambda - s} - p_{i+1} \frac{s}{\lambda - s},$$
(46)

where we set $p_i := \mathbb{P}(W_i = 0)$. Our goal is to write (46) fully in terms of the functions $\Phi_{W_i}(s)$. To this end, performing the change of variable v := su, we obtain

$$\Phi_{V_i W_i}(s) = \int_0^1 \mathbb{E}(e^{-suW_i}) \, \mathrm{d}u = \frac{1}{s} \int_0^s \Phi_{W_i}(v) \, \mathrm{d}v.$$
(47)

By multiplying with $\lambda - s$, we thus obtain the following recursive integral equation.

Lemma 8. For $i \in \mathbb{N}$,

$$\Phi_{W_{i+1}}(s) = \frac{\lambda \Phi_B(s)}{s(\lambda - s)} \int_0^s \Phi_{W_i}(v) \,\mathrm{d}v - \frac{s}{\lambda - s} p_{i+1}.$$
(48)

Since the LST $\Phi_{W_0}(s) = e^{-sw}$ of W_0 is known, Relation (48) in principle allows us to recursively determine all the transforms $\Phi_{W_i}(\cdot)$, $i \in \mathbb{N}$. Observe that, when $s = \lambda$, the right-hand side should become zero; using (47) we obtain

$$p_{i+1} = \frac{\Phi_B(\lambda)}{\lambda} \int_0^\lambda \Phi_{W_i}(v) \, \mathrm{d}v.$$
(49)

This formula can easily be interpreted probabilistically, using the memoryless property of the exponential distribution for A_i :

$$\mathbb{P}(W_{i+1}=0) = P(A_i \ge B_i + V_i W_i) = \mathbb{P}(A_i \ge B_i) P(A_i \ge V_i W_i) = \Phi_B(\lambda) \Phi_{V_i W_i}(\lambda).$$

It is not possible to obtain explicit expressions for Φ_{W_i} . However, as so often, one can utilize the method of generating functions to turn the recursion (48) into some sort of differential or integral equation. Therefore we multiply Equation (48) by r^{i+1} and sum over *i* to obtain

$$U_W(r,s) = \frac{\lambda r \,\Phi_B(s)}{s(\lambda - s)} \int_0^s U_W(r,v) \,\mathrm{d}v + K(s), \tag{50}$$

where, to simplify the notation,

$$K(s) = \Phi_{W_0}(s) - \frac{s}{\lambda - s} \left(U_W(r, \infty) - p_0 \right).$$
(51)

With $I(s) = \int_0^s U_W(r, v) dv$ we obtain the linear first order differential equation

$$I'(s) = \frac{\lambda r \Phi_B(s)}{s(\lambda - s)} I(s) + K(s).$$
(52)

As follows by standard techniques, this inhomogeneous differential equation is solved by

$$I(s) = \exp\left(\lambda r \int_{c}^{s} \frac{\Phi_{B}(t)}{t(\lambda - t)} dt\right) \left(\theta(c) + \int_{c}^{s} K(u) \exp\left(-\lambda r \int_{c}^{u} \frac{\Phi_{B}(t)}{t(\lambda - t)} dt\right) du\right),$$
(53)

where necessarily $\theta(c) = I(c)$ and we assume for the time being that $c \in (s, \lambda)$ if $s < \lambda$ and $c \in (\lambda, s)$ if $\lambda < s$. Since Φ_B is bounded and bounded away from zero, we have as $s \downarrow \lambda$ (and likewise if $s \uparrow \lambda$ in the case where $s < \lambda$)

$$\exp\left(\lambda r \int_c^s \frac{\Phi_B(t)}{t(\lambda-t)} \,\mathrm{d}t\right) \to \infty.$$

As a consequence, to make sure $U_W(r, \lambda)$ remains bounded,

$$\theta(c) = -\int_{c}^{\lambda} K(u) \exp\left(-\lambda r \int_{c}^{u} \frac{\Phi_{B}(t)}{t(\lambda - t)} dt\right) du,$$

and therefore

$$I(s) = \exp\left(\lambda r \int_{c}^{s} \frac{\Phi_{B}(t)}{t(\lambda - t)} dt\right) \int_{\lambda}^{s} K(u) \exp\left(-\lambda r \int_{c}^{u} \frac{\Phi_{B}(t)}{t(\lambda - t)} dt\right) du$$
$$= \int_{\lambda}^{s} K(u) \exp\left(\lambda r \int_{u}^{s} \frac{\Phi_{B}(t)}{t(\lambda - t)} dt\right) du.$$
(54)

It is to be noted that the integrand in (54) tends to ∞ as $s \to \lambda$, which follows from the finiteness of $I(\lambda)$. Inserting (54) into (52), recalling hat $U_W(r, s) = I'(s)$, yields the following expression for the generating function U_W :

$$U_W(r,s) = K(s) + \frac{\lambda r \Phi_B(s)}{s(\lambda - s)} \int_{\lambda}^{s} K(u) \exp\left(\lambda r \int_{u}^{s} \frac{\Phi_B(t)}{t(\lambda - t)} dt\right) du.$$
(55)

Plugging (51) into (55) we obtain

$$U_W(r,s) = K(s) + \frac{\lambda r \Phi_B(s)}{s(\lambda - s)} \int_{\lambda}^{s} \left(\Phi_{W_0}(u) - \frac{u \left(U_W(r, \infty) - p_0 \right)}{\lambda - u} \right) \exp\left(\lambda r \int_{u}^{s} \frac{\Phi_B(t)}{t(\lambda - t)} dt \right) du,$$

It remains to determine $U_W(r, \infty)$. Keeping in mind that $U_W(r, s) - K(s) \rightarrow 0$ by (51) we obtain after some rearrangements,

$$U_W(r,\infty) = p_0 - \frac{\int_{\lambda}^{\infty} \Phi_{W_0}(u) \exp\left(-\lambda r \int_{u}^{\infty} \frac{\Phi_B(t)}{t(t-\lambda)} dt\right) du}{\int_{\lambda}^{\infty} \frac{u}{u-\lambda} \exp\left(-\lambda r \int_{u}^{\infty} \frac{\Phi_B(t)}{t(t-\lambda)} dt\right) du}.$$

We summarize our findings in the following theorem.

Theorem 9. *For* $r \in (0, 1)$ *,*

$$U_W(r,s) = K(s) + \frac{\lambda r \Phi_B(s)}{s(\lambda - s)} \int_{\lambda}^{s} K(u) \exp\left(\lambda r \int_{u}^{s} \frac{\Phi_B(t)}{t(\lambda - t)} dt\right) du.$$
(56)

where

$$K(s) = \Phi_{W_0}(s) + \frac{\frac{s}{\lambda - s} \int_{\lambda}^{\infty} \Phi_{W_0}(u) \exp\left(-\lambda r \int_{u}^{\infty} \frac{\Phi_B(t)}{t(t-\lambda)} dt\right) du}{\int_{\lambda}^{\infty} \frac{u}{u-\lambda} \exp\left(-\lambda r \int_{u}^{\infty} \frac{\Phi_B(t)}{t(t-\lambda)} dt\right) du}.$$

We already noted that since the V_i are uniformly distributed on [0, 1] we have $\mathbb{E}(\log |V|) < 0$. Hence we always have $W_i \Rightarrow W$ as $i \to \infty$ for some proper random variable W. Its LST is given in the following theorem.

Theorem 10. W_i converges weakly to a proper limit W as $i \to \infty$, and

$$\Phi_W(s) = \frac{p_{\infty}}{s - \lambda} \left(s - \frac{\lambda \Phi_B(s)}{s} \int_{\lambda}^{s} \frac{u}{u - \lambda} \exp\left(\lambda \int_{u}^{s} \frac{\Phi_B(t)}{t(\lambda - t)} dt\right) du \right).$$
(57)

where

$$p_{\infty} = \left[\int_0^{\lambda} \frac{1}{\lambda - u} \exp\left(-\int_0^u \left(\frac{\lambda \Phi_B(t)}{t(\lambda - t)} - \frac{1}{t} \right) dt \right) du \right]^{-1}.$$
 (58)

Proof. We apply an Abelian theorem and obtain (57) after multiplying both sides of (56) by 1 - r and letting *r* tend to one. We thereby use the fact that

$$\lim_{r\uparrow 1}(1-r)K(s)=\frac{sp_{\infty}}{s-\lambda}.$$

The relation for p_{∞} follows by noting that $\Phi_W(0) = 1$, so that

$$\frac{1}{p_{\infty}} = \lim_{s \downarrow 0} \frac{1}{s - \lambda} \left(s - \frac{\lambda \Phi_B(s)}{s} \int_{\lambda}^{s} \frac{u}{u - \lambda} \exp\left(\lambda \int_{u}^{s} \frac{\Phi_B(t)}{t(\lambda - t)} dt\right) du \right)$$
$$= \lim_{s \downarrow 0} \frac{1}{s} \int_{s}^{\lambda} \frac{u}{\lambda - u} \exp\left(-\lambda \int_{s}^{u} \frac{\Phi_B(t)}{t(\lambda - t)} dt\right) du$$
$$= \lim_{s \downarrow 0} \int_{s}^{\lambda} \frac{1}{\lambda - u} \exp\left(-\int_{s}^{u} \left(\frac{\lambda \Phi_B(t)}{t(\lambda - t)} - \frac{1}{t}\right) dt\right) du.$$

Now

$$\frac{\lambda \Phi_B(t)}{t(\lambda - t)} - \frac{1}{t} \to \frac{1}{\lambda} - \mathbb{E}(B),$$

as $t \downarrow 0$, so that we can safely let $s \downarrow 0$ and obtain the finite and non-zero limit (58).

Remark 3. The expected value $\mathbb{E}(W)$ can be expressed in terms of the parameters λ , $\mathbb{E}(B)$ and p_{∞} as follows. Letting $i \to \infty$ in (46) yields

$$\Phi_W(s) = \Phi_{VW}(s)\Phi_B(s)\frac{\lambda}{\lambda-s} - \frac{s}{\lambda-s}p_{\infty}.$$
(59)

After a rearrangement of terms this becomes

$$p_{\infty} = \Phi_W(s) + \lambda \frac{1 - \Phi_W(s) + \Phi_{VW+B}(s) - 1}{s}.$$

As $s \downarrow 0$ the right-hand side tends to $1 + \lambda (\mathbb{E}(W) - \mathbb{E}(VW + B))$ and since $\mathbb{E}(VW + B) = \mathbb{E}(V)\mathbb{E}(W) + \mathbb{E}(B) = \frac{1}{2}\mathbb{E}(W) + \mathbb{E}(B)$, we obtain

$$\mathbb{E}(W) = 2\left(\mathbb{E}(B) - \frac{1 - p_{\infty}}{\lambda}\right).$$
(60)

This yields the inequality $\mathbb{P}(W = 0) \ge 1 - \lambda \mathbb{E}(B)$, the value that one would get if $V_i \equiv 1$.

Example 2. Even if the B_i are $\exp(\mu)$ distributed the computation of $U_W(r, s)$ becomes quite involved. By means of a partial fraction expansion, after considerable calculus we obtain

$$\frac{u}{\lambda - u} \exp\left(\lambda \int_{u}^{s} \frac{\Phi_{B}(t)}{t(\lambda - t)} dt\right) = \frac{s}{\lambda - u} \left(\frac{\lambda - u}{\lambda - s}\right)^{\frac{\mu}{\lambda + \mu}} \left(\frac{\mu + u}{\mu + s}\right)^{\frac{\lambda}{\lambda + \mu}}$$

For $s < \lambda$ the integral can be expressed in terms of the incomplete beta function $B(x, a, b) = \int_0^x u^{a-1} (1-u)^{b-1} du$:

$$\int_{\lambda}^{s} \frac{s}{u-\lambda} \left(\frac{\lambda-u}{\lambda-s}\right)^{\frac{\mu}{\lambda+\mu}} \left(\frac{\mu+u}{\mu+s}\right)^{\frac{\lambda}{\lambda+\mu}} du = \frac{s(\lambda+\mu)}{(\mu+s)^{\frac{\lambda}{\lambda+\mu}}(\lambda-s)^{\frac{\mu}{\lambda+\mu}}} B\left(\frac{\lambda-s}{\lambda+\mu}, \frac{\mu}{\lambda+\mu}, 1+\frac{\lambda}{\lambda+\mu}\right).$$

We then obtain

$$\Phi_W(s) = p_{\infty} \cdot \left(\frac{\lambda \mu (\lambda + \mu) B\left(\frac{\lambda - s}{\lambda + \mu}, \frac{\mu}{\lambda + \mu}, 1 + \frac{\lambda}{\lambda + \mu}\right)}{(\mu + s)^{1 + \frac{\lambda}{\lambda + \mu}} (\lambda - s)^{1 + \frac{\mu}{\lambda + \mu}}} - \frac{s}{\lambda - s} \right).$$

This leads to

$$p_{\infty} = \frac{\mu^{\frac{\lambda}{\lambda+\mu}} \lambda^{\frac{\mu}{\lambda+\mu}}}{(\lambda+\mu)B\left(\frac{\lambda}{\lambda+\mu}, \frac{\mu}{\lambda+\mu}, 1+\frac{\lambda}{\lambda+\mu}\right)},$$

so that, at least for $s < \lambda$,

$$\Phi_{W}(s) = \frac{\mu^{\frac{\lambda}{\lambda+\mu}} \lambda^{\frac{\mu}{\lambda+\mu}}}{(\lambda+\mu)B\left(\frac{\lambda}{\lambda+\mu}, \frac{\mu}{\lambda+\mu}, 1+\frac{\lambda}{\lambda+\mu}\right)} \cdot \left(\frac{\lambda\mu(\lambda+\mu)B\left(\frac{\lambda-s}{\lambda+\mu}, \frac{\mu}{\lambda+\mu}, 1+\frac{\lambda}{\lambda+\mu}\right)}{(\mu+s)^{1+\frac{\lambda}{\lambda+\mu}}(\lambda-s)^{1+\frac{\mu}{\lambda+\mu}}} - \frac{s}{\lambda-s}\right).$$
(61)

Unfortunately a similar expression for the $s > \lambda$ case is not available. Instead one obtains expressions that involve hypergeometric functions. Also it seems very hard to obtain higher moments from (61) by means of differentiation. It is possible, however, to derive a recursion formula for the moments $\omega_k := \mathbb{E}(W^k)$ (where we assume their existence for k = 1, 2, ..., j, say) if we start with (59), which in our example becomes

$$(\lambda - s)\Phi_W(s) = \frac{\lambda\mu}{s(\mu + s)} \int_0^s \Phi_W(v) \,\mathrm{d}v - sp_\infty.$$
(62)

For s < 0 the expansion $\Phi_W(-s) = \sum_{k=0}^j \frac{\omega_k}{k!} s^k + o(s^j)$ holds. Inserting this into (62) yields

$$(\mu - s)(\lambda + s) \sum_{k=0}^{j} \frac{\omega_k}{k!} s^k + o(s^j) = \lambda \mu \sum_{k=0}^{j} \frac{\omega_k}{(k+1)!} s^k - s(\mu + s) p_{\infty} + o(s^j), \quad s \downarrow 0.$$

Equating the coefficients on both sides leads to

$$\begin{split} \lambda \mu \frac{\omega_{k}}{k!} &+ (\mu - \lambda) \frac{\omega_{k-1}}{(k-1)!} - \frac{\omega_{k-2}}{(k-2)!} \mathbb{1}_{\{k \ge 2\}} \\ &= \left(\lambda \mu \frac{\omega_{1}}{2} - \mu p_{\infty} \right) \mathbb{1}_{\{k=1\}} + \left(\lambda \mu \frac{\omega_{2}}{6} - p_{\infty} \right) \mathbb{1}_{\{k=2\}} + \frac{\lambda \mu \omega_{k}}{(k+1)!} \mathbb{1}_{\{k \ge 3\}}, \quad k \in \{1, \dots, j\}. \end{split}$$

Then, in accordance with the general result (60),

$$\omega_1 = 2\left(\frac{1}{\mu} - \frac{1 - p_{\infty}}{\lambda}\right).$$

Moreover,

$$\omega_2 = 3 \frac{1 - p_{\infty} - (\mu - \lambda)\omega_1}{\mu \lambda},$$

and

$$\omega_k = \frac{(k^2 - 1)\omega_{k-2} - (k+1)(\mu - \lambda)\omega_{k-1}}{\lambda\mu}, \quad k \in \{3, \dots, j\}.$$

6. Discussion and concluding remarks

This paper has analyzed three reflected autoregressive processes specified by the stochastic recursion $W_{i+1} = [V_iW_i + B_i - A_i]^+$. While the classical case of $V \equiv 1$ has been widely studied in the queueing literature, our more general setting allows explicit analysis only in special cases. The three special cases we have considered are: (i) V attains negative values only and B has a rational LST, (ii) V equals a positive value a with certain probability $p \in (0, 1)$ and is negative otherwise, and both A and B have a rational LST, (iii) V is uniformly distributed on [0, 1], and A is exponentially distributed. In all three cases we present transient and stationary results, where the transient results are in terms of the transform at a geometrically distributed epoch.

Cases which might allow explicit analysis are, for example:

1. A generalization of Model III to the case in which $V = U^{1/\alpha}$, where U has a uniform distribution on [0, 1] and $\alpha > 0$. In this case (48) becomes

$$s^{\alpha-1}\Phi_W(s) = \frac{\alpha\lambda\Phi_B(s)}{s(\lambda-s)}\int_0^s v^{\alpha-1}\Phi_W(v)\,\mathrm{d}v - \frac{s^\alpha}{(\lambda-s)}p_{i+1}.$$

Letting $U_W^{(\alpha)}(r,s) = s^{\alpha-1}U_W(r,s)$ this yields

$$U_W^{(\alpha)}(r,s) = \frac{\lambda r \alpha \, \Phi_B(s)}{s(\lambda-s)} \int_0^s U_W^{(\alpha)}(r,v) \, \mathrm{d}v + K^{(\alpha)}(s), \tag{63}$$

where

$$K^{(\alpha)}(s) = s^{\alpha - 1} \Phi_{W_0}(s) - \frac{s^{\alpha}}{\lambda - s} (U_W^{(\alpha)}(r, \infty) - p_0)$$

Equation (63) is of the exact same type as (50), only with *r* replaced by $r\alpha$. This allows one to derive $U_W^{(\alpha)}(r, s)$ in the same way as before.

- 2. A combination of Models I and III, allowing V to be either negative or having a distribution as in 1.
- 3. One might consider more general recursions, such as the high-order Lindley equations analyzed in, e.g., [2, 17, 19].

Another possible line of research concerns scaling limits and asymptotics. In particular, tail asymptotics seem to be within reach; in heavy-tailed cases these may be identified relying on a Tauberian approach. One also anticipates that, under particular scalings, an explicit analysis is possible. Specifically, one would expect that a diffusion analysis similar to the one presented in [6] can be performed.

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